

# FYS3410 - Vår 2010 (Kondenserte fasers fysikk)

<http://www.uio.no/studier/emner/matnat/fys/FYS3410/index-eng.xml>

Based on Introduction to Solid State Physics by Kittel

## Course content

- Periodic structures, understanding of diffraction experiment and reciprocal lattice
- Crystal binding, elastic strain and waves
- Imperfections in crystals: diffusion, point defects, dislocations
- Crystal vibrations: phonon heat capacity and thermal conductivity
- Free electron Fermi gas: density of states, Fermi level, and electrical conductivity
- Electrons in periodic potential: energy bands theory classification of metals, semiconductors and insulators
- Semiconductors: band gap, effective masses, charge carrier distributions, doping, pn-junctions
- Metals: Fermi surfaces, temperature dependence of electrical conductivity

Andrej Kuznetsov, Dept of Physics and Centre for Material Science and Nanotechnology  
Postboks 1048 Blindern, 0316 OSLO  
Tel: +47-22852870, e-post: [andrej.kuznetsov@fys.uio.no](mailto:andrej.kuznetsov@fys.uio.no)  
visiting address: MiNaLab, Gaustadaleen 23b

## FYS3410 lecture schedule and exams: Spring 2010

M/18/1/2010:	Introduction and motivation. Periodicity and lattices	2h
W/20/1/2010:	Index system for crystal planes. Crystal structures	1h
M/25/1/2010:	Reciprocal space, Laue condition and Ewald construction	2h
W/27/1/2010:	Brillouin Zones. Interpretation of a diffraction experiment	1h
M/01/2/2010:	Crystal binding, elastic strain and waves	2h
W/03/2/2010:	Point defects, case study – vacancies	1h
M/08/2/2010:	Point defects and atomic diffusion	2h
W/10/2/2010:	Diffusion (continuation); dislocations	1h
M/15/2/2010:	Crystal vibrations and phonons	2h
W/17/2/2010:	Crystal vibrations and phonons	1h
M/22/2/2010:	Planck distribution and density of states	2h
W/24/2/2010:	Debye model	1h
M/01/3/2010:	Einstein model and general result for density of states	2h
W/03/3/2010:	Thermal conductivity	1h
M/08/3/2010:	Free electron Fermi gas in 1D and 3D – ground state	2h
W/10/3/2010:	Density of states, effect of temperature – FD distribution	1h
M/15/3/2010:	Heat capacity of FEFG	2h
W/17/3/2010:	Repetition	1h
22/3/2010:	Mid-term exam	

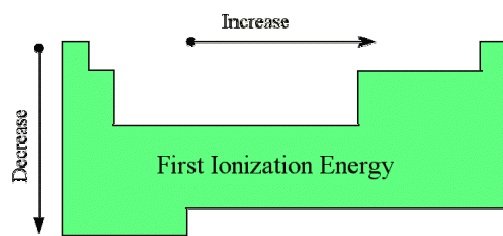
<b>M/14/4/2010:</b>	<b>Electrical and thermal conductivity in metals</b>	<b>2h</b>
<b>W/12/4/2010:</b>	<b>Bragg reflection of electron waves at the boundary of BZ</b>	<b>1h</b>
<b>M/19/4/2010:</b>	<b>Energy bands, Kronig - Penny model</b>	<b>2h</b>
<b>W/21/4/2010:</b>	<b>Empty lattice approximation; number of orbitals in a band</b>	<b>1h</b>
<b>M/26/4/2010:</b>	<b>Semiconductors, effective mass method, intrinsic carriers</b>	<b>2h</b>
<b>W/28/4/2010:</b>	<b>Impurity states in semiconductors and carrier statistics</b>	<b>1h</b>
<b>M/03/5/2010:</b>	<b>p-n junctions and heterojunctions</b>	<b>2h</b>
<b>W/05/5/2010:</b>	<b>surface structure, surface states, Schottky contacts</b>	<b>2h</b>
<b>M/10/5/2010:</b>	<b>no lectures</b>	
<b>W/12/5/2010:</b>	<b>no lectures</b>	
<b>W/19/5/2010:</b>	<b>Repetition</b>	<b>2h</b>
<b>W26/5/2010:</b>	<b>Repetition</b>	<b>2h</b>
<b>28/5/2010:</b>	<b>Final Exam (sensor: Prof. Arne Nylandsted Larsen at the Århus University, Denmark, <a href="http://person.au.dk/en/arl@phys.au.dk">http://person.au.dk/en/arl@phys.au.dk</a>)</b>	

## **Lecture 5: Crystal binding, elastic strain and waves**

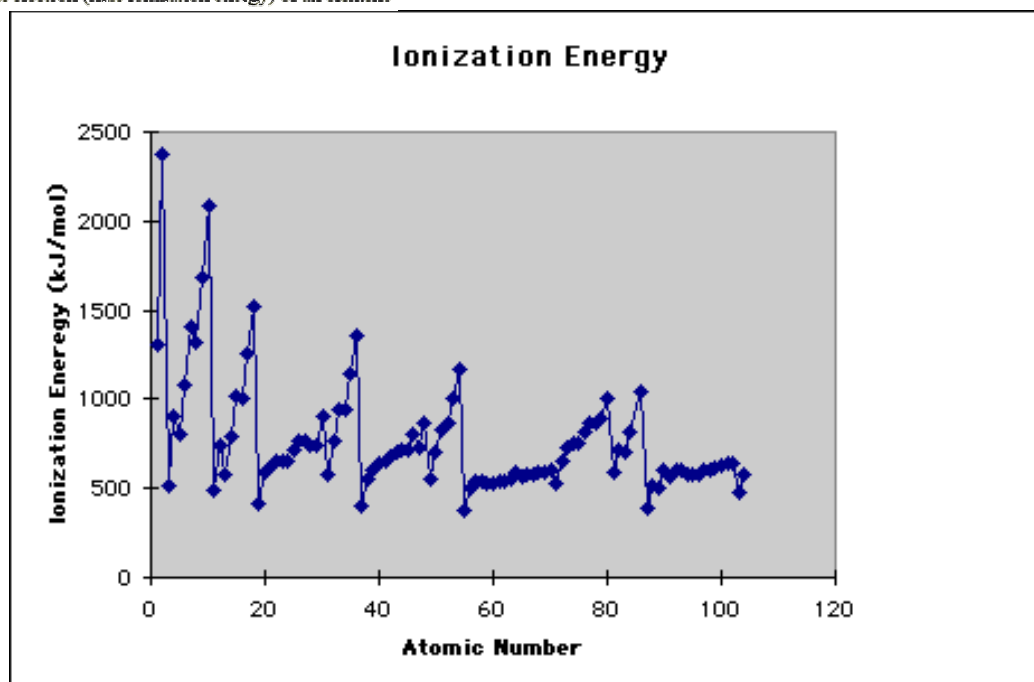
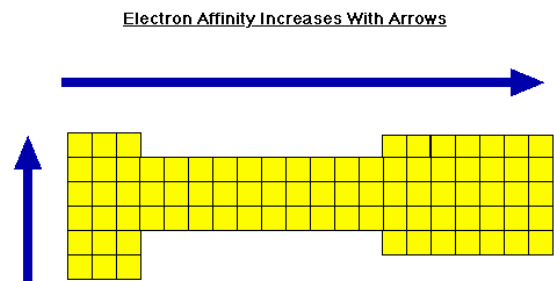
- **ionization and cohesive energies in the context of periodic table;**
- **interaction between two atoms in terms of attraction and repulsion forces;**
- **analysis of elastic strain**
- **elastic waves in cubic crystals**

## Lecture 5: Crystal binding, elastic strain and waves

- ionization and cohesive energies in the context of periodic table;
- interaction between two atoms in terms of attraction and repulsion forces;
- analysis of elastic strain
- elastic waves in cubic crystals



General trends for the energy required to remove the first electron (first ionization energy) of an element



**Table 1 Cohesive energies**

Li 158. 1.63 37.7		Be 320. 3.32 76.5		Energy required to form separated neutral atoms in their ground electronic state from the solid at 0 K at 1 atm. The data were supplied by Prof. Leo Brewer.												B 561 5.81 134	C 711 7.37 170.	N 474. 4.92 113.4	O 251 2.60 60.03	F 81.0 0.84 19.37	Ne 1.92 0.020 0.46												
Na 107. 1.113 25.67		Mg 145. 1.51 34.7														Al 327. 3.39 78.1	Si 446. 4.63 106.7	P 331. 3.43 79.16	S 275. 2.85 65.75	Cl 135. 1.40 32.2	Ar 7.74 0.080 1.85												
																←----- kJ/mol -----→						←----- eV/atom -----→						←----- kcal/mol -----→					
K 90.1 0.934 21.54	Ca 178. 1.84 42.5	Sc 376 3.90 89.9	Ti 468. 4.85 111.8	V 512. 5.31 122.4	Cr 395. 4.10 94.5	Mn 282. 2.92 67.4	Fe 413. 4.28 98.7	Co 424. 4.39 101.3	Ni 428. 4.44 102.4	Cu 336. 3.49 80.4	Zn 130 1.35 31.04	Ga 271. 2.81 64.8	Ge 372. 3.85 88.8	As 285.3 2.96 68.2	Se 237 2.46 56.7	Br 118. 1.22 28.18	Kr 11.2 0.116 2.68																
Rb 82.2 0.852 19.64	Sr 166. 1.72 39.7	Y 422. 4.37 100.8	Zr 603. 6.25 144.2	Nb 730. 7.57 174.5	Mo 658 6.82 157.2	Tc 661. 6.85 158.	Ru 650. 6.74 155.4	Rh 554. 5.75 132.5	Pd 376. 3.89 89.8	Ag 284. 2.95 68.0	Cd 112. 1.16 26.73	In 243. 2.52 58.1	Sn 303. 3.14 72.4	Sb 265. 2.75 63.4	Te 211 2.19 50.34	I 107. 1.11 25.62	Xe 15.9 0.16 3.80																
Cs 77.6 0.804 18.54	Ba 183. 1.90 43.7	La 431 4.47 103.1	Hf 621 6.44 148.4	Ta 782. 8.10 186.9	W 859. 8.90 205.2	Re 775. 8.03 185.2	Os 788. 8.17 188.4	Ir 670. 6.94 160.1	Pt 564. 5.84 134.7	Au 368. 3.81 87.96	Hg 65. 0.67 15.5	Tl 182. 1.88 43.4	Pb 196. 2.03 46.78	Bi 210. 2.18 50.2	Po 144. 1.50 34.5	At   	Rn 19.5 0.202 4.66																
Fr   	Ra 160. 1.66 38.2	Ac 410. 4.25 98.	Ce 417. 4.32 99.7	Pr 357. 3.70 85.3	Nd 328. 3.40 78.5	Pm   	Sm 206. 2.14 49.3	Eu 179. 1.86 42.8	Gd 400. 4.14 95.5	Tb 391. 4.05 93.4	Dy 294. 3.04 70.2	Ho 302. 3.14 72.3	Er 317. 3.29 75.8	Tm 233. 2.42 55.8	Yb 154. 1.60 37.1	Lu 428. 4.43 102.2																	
			Th 598. 6.20 142.9	Pa   	U 536. 5.55 128.	Np 456 4.73 109.	Pu 347 3.60 83.0	Am 264. 2.73 63.	Cm 385 3.99 92.1	Bk   	Cf   	Es   	Fm   	Md   	No   	Lr   																	

Kcal/mol = 0.0434 eV/molecule

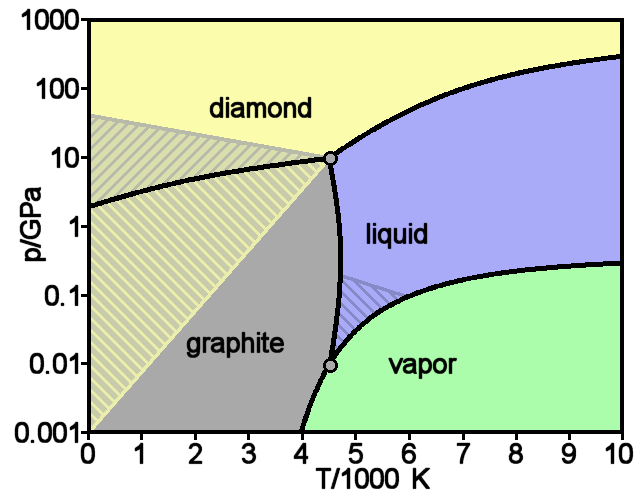
KJ/mol = 0.0104 eV/molecule

**Table 2 Melting points, in K.**

(After R. H. Lamoreaux)

Li 453.7		Be 1562														B 2365	C	N 63.15	O 54.36	F 53.48	Ne 24.56
Na 371.0		Mg 922														Al 933.5	Si 1687	P 317 863	S 388.4	Cl 172.2	Ar 83.81
K 336.3	Ca 1113	Sc 1814	Ti 1946	V 2202	Cr 2133	Mn 1520	Fe 1811	Co 1770	Ni 1728	Cu 1358	Zn 692.7	Ga 302.9	Ge 1211	As 1089	Se 494	Br 265.9	Kr 115.8				
Rb 312.6	Sr 1042	Y 1801	Zr 2128	Nb 2750	Mo 2895	Tc 2477	Ru 2527	Rh 2236	Pd 1827	Ag 1235	Cd 594.3	In 429.8	Sn 505.1	Sb 903.9	Te 722.7	I 386.7	Xe 161.4				
Cs 301.6	Ba 1002	La 1194	Hf 2504	Ta 3293	W 3695	Re 3459	Os 3306	Ir 2720	Pt 2045	Au 1338	Hg 234.3	Tl 577	Pb 600.7	Bi 544.6	Po 527	At  	Rn  				
Fr   	Ra 973	Ac 1324	Ce 1072	Pr 1205	Nd 1290	Pm   	Sm 1346	Eu 1091	Gd 1587	Tb 1632	Dy 1684	Ho 1745	Er 1797	Tm 1820	Yb 1098	Lu 1938					
			Th 2031	Pa 1848	U 1406	Np 910	Pu 913	Am 1449	Cm 1613	Bk 1562	Cf   	Es   	Fm   	Md   	No   	Lw   					

At atmospheric pressure carbon has no melting point as its triple point is at  $10.8 \pm 0.2$  MPa and  $4600 \pm 300$  K, so it sublimates at about 3900 K



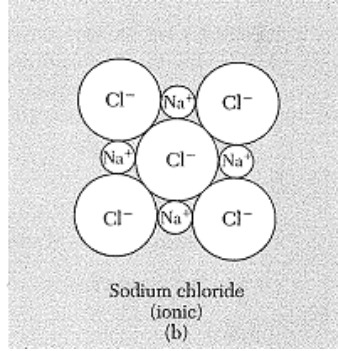
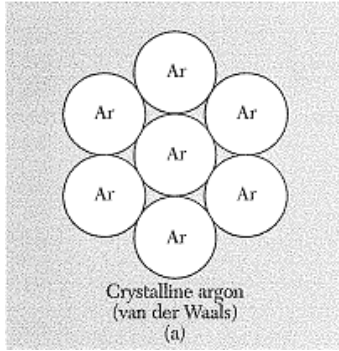
101.325 kPa is “one standard atmosphere” and  $1 \text{ Pa} = 1 \text{ N/m}^2$

$$1 \text{ J} = 1 \text{ N} \cdot \text{m} = \left( \frac{\text{kg} \cdot \text{m}}{\text{s}^2} \right) \cdot \text{m} = \frac{\text{kg} \cdot \text{m}^2}{\text{s}^2} = \text{Pa} \cdot \text{m}^3 = 1 \text{ W} \cdot \text{s}$$

$$1 \text{ eV} = 1.602176487 \times 10^{-19} \text{ Joule}$$

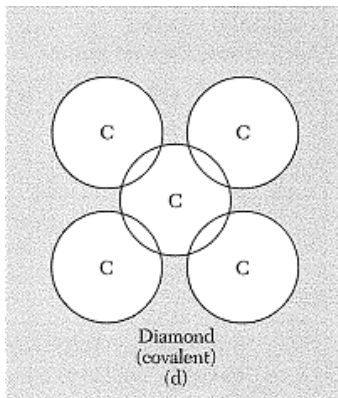
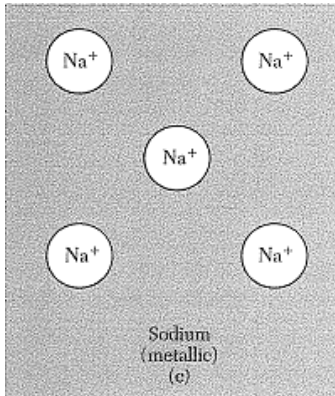
## Lecture 5: Crystal binding, elastic strain and waves

- ionization and cohesive energies in the context of periodic table;
- interaction between two atoms in terms of attraction and repulsion forces;
- analysis of elastic strain
- elastic waves in cubic crystals



Cohesive energy = energy required to break up crystal into neutral free atoms.

Lattice energy (ionic crystals) = energy required to break up crystal into free ions.

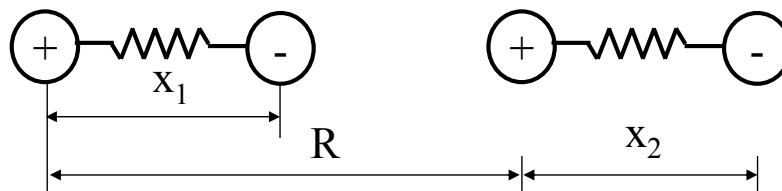


- ⊙ The atoms of the inert gasses attract each other via the so-called Van Der Waals interaction (dipole-dipole interaction)
- ⊙ The unperturbed Hamiltonian of the system is the one of two independent oscillators and is given by:

$$H_0 = \frac{p_1^2}{2m} + \frac{1}{2}cx_1^2 + \frac{p_2^2}{2m} + \frac{1}{2}cx_2^2$$

- ⊙ The interaction energy of the oscillators is given by:

$$H_1 = \frac{1}{4\pi\epsilon} \left[ \frac{e^2}{R} + \frac{e^2}{R+x_2-x_1} - \frac{e^2}{R+x_2} - \frac{e^2}{R-x_1} \right] \approx -2 \frac{e^2 x_1 x_2}{4\pi\epsilon R^3}$$



- ⊙ Besides the Van der Waals interaction, when two atoms are brought together, their charge distribution begins to overlap, giving rise to the exchange interaction due to the Pauli exclusion principle.
- ⊙ Since this interaction is difficult to be evaluated from first principles, it is usually parametrized with a term that goes as  $1/R^{12}$ , which then when added to the Van der Waals term gives rise to the famous Lennard-Jones potential that is of the following form:

$$U(R) = 4\varepsilon \left[ \left( \frac{\sigma}{R} \right)^{12} - \left( \frac{\sigma}{R} \right)^6 \right]$$

1. One first re-writes the Lennard-Jones potential as a sum of all pairs of atoms in the crystal as:

$$U_{TOT}(R) = \frac{1}{2} N(4\varepsilon) \left[ \sum_j \left( \frac{\sigma}{\rho_{ij}R} \right)^{12} - \sum_j \left( \frac{\sigma}{\rho_{ij}R} \right)^6 \right]$$

where  $\rho_{ij}R$  is the distance between the reference atom  $i$  and any other atom  $j$ . For FCC structure with 12 nearest neighbors, we have:

$$\sum_j \left( \frac{1}{\rho_{ij}} \right)^{12} = 12.13188, \quad \sum_j \left( \frac{1}{\rho_{ij}} \right)^6 = 14.45392$$

2. The condition that the net force on the atom is zero then gives  $R_0/\sigma = 1.09$ , and the cohesive energy is:

$$U_{TOT}(R) = -2.15(4N\varepsilon)$$



# Crystals of Inert Gases

**Table 4 Properties of inert gas crystals**  
(Extrapolated to 0 K and zero pressure)

	Nearest-neighbor distance, in Å	Experimental cohesive energy		Melting point, K	Ionization potential of free atom, eV	Parameters in Lennard-Jones potential, Eq. 10	
		kJ/mol	eV/atom			$\epsilon$ , in $10^{-16}$ erg	$\sigma$ , in Å
He	(liquid at zero pressure)				24.58	14	2.56
Ne	3.13	1.88	0.02	24.56	21.56	50	2.74
Ar	3.76	7.74	0.080	83.81	15.76	167	3.40
Kr	4.01	11.2	0.116	115.8	14.00	225	3.65
Xe	4.35	16.0	0.17	161.4	12.13	320	3.98

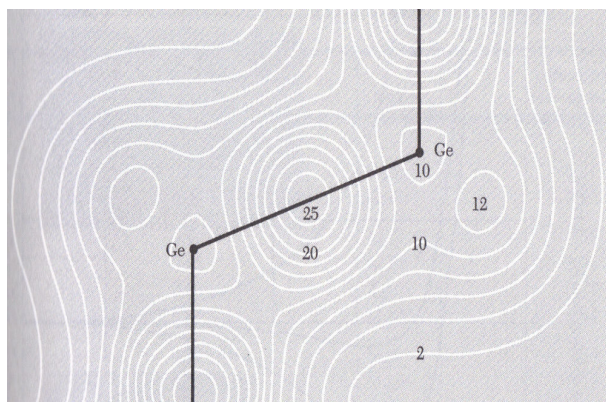
Atoms:

- high ionization energy
- outermost shell filled
- charge distribution spherical

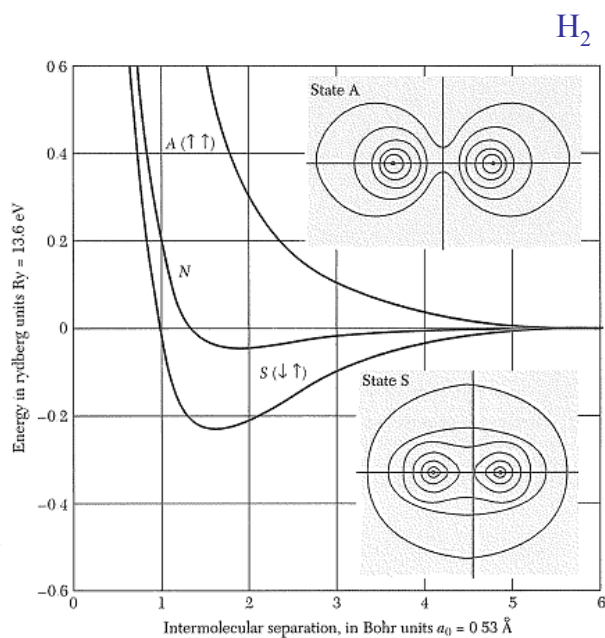
Crystal:

- transparent insulators
- weakly bonded
- low melting point
- closed packed (fcc, except He<sup>3</sup> & He<sup>4</sup>).

# Covalent Crystals



- Electron pair localized midway of bond.
- Tetrahedral: diamond, zinc-blende structures.
- Low filling: 0.34 vs 0.74 for closed-packed.



Pauli exclusion →  
exchange interaction



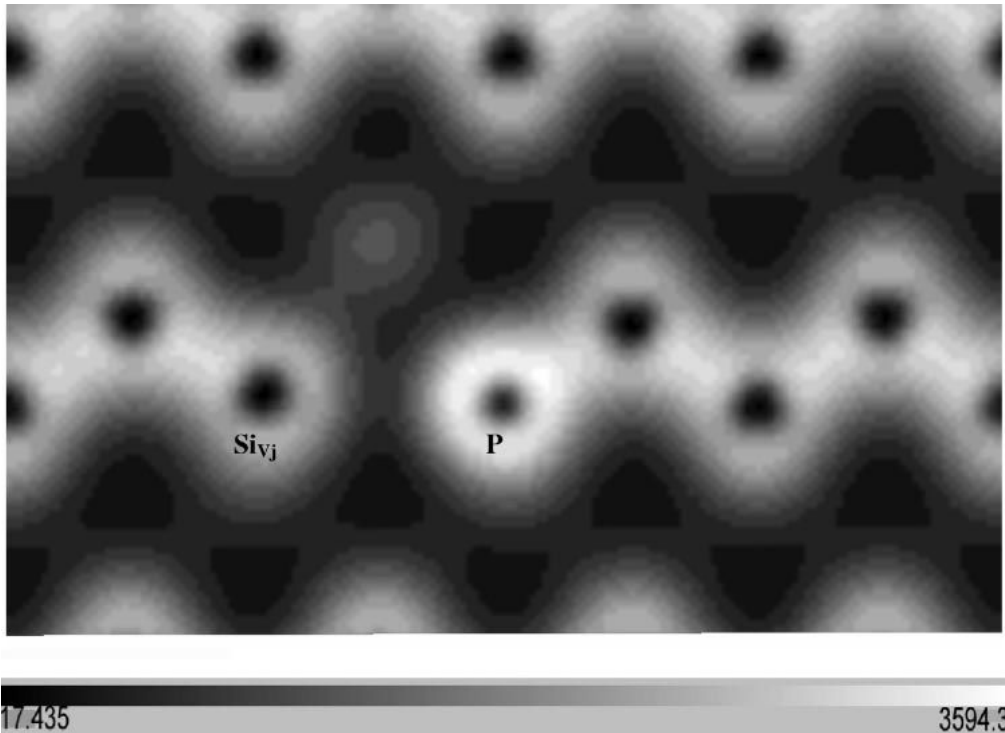


FIG. 4. Electron localization function map for the (110) plane, containing the phosphorus atom and vacancy for the case of the negatively charged PV complex (PHYSICAL REVIEW B 70, 115204 (2004))

## Lecture 5: Crystal binding, elastic strain and waves

- ionization and cohesive energies in the context of periodic table;
- interaction between two atoms in terms of attraction and repulsion forces;
- analysis of elastic strain
- elastic waves in cubic crystals

*stress*  $\sigma = \frac{\text{load } W}{\text{area } A}$

$$\epsilon_{ij} = S_{ij} \cdot \sigma_{ij}$$

**Hooke's law**

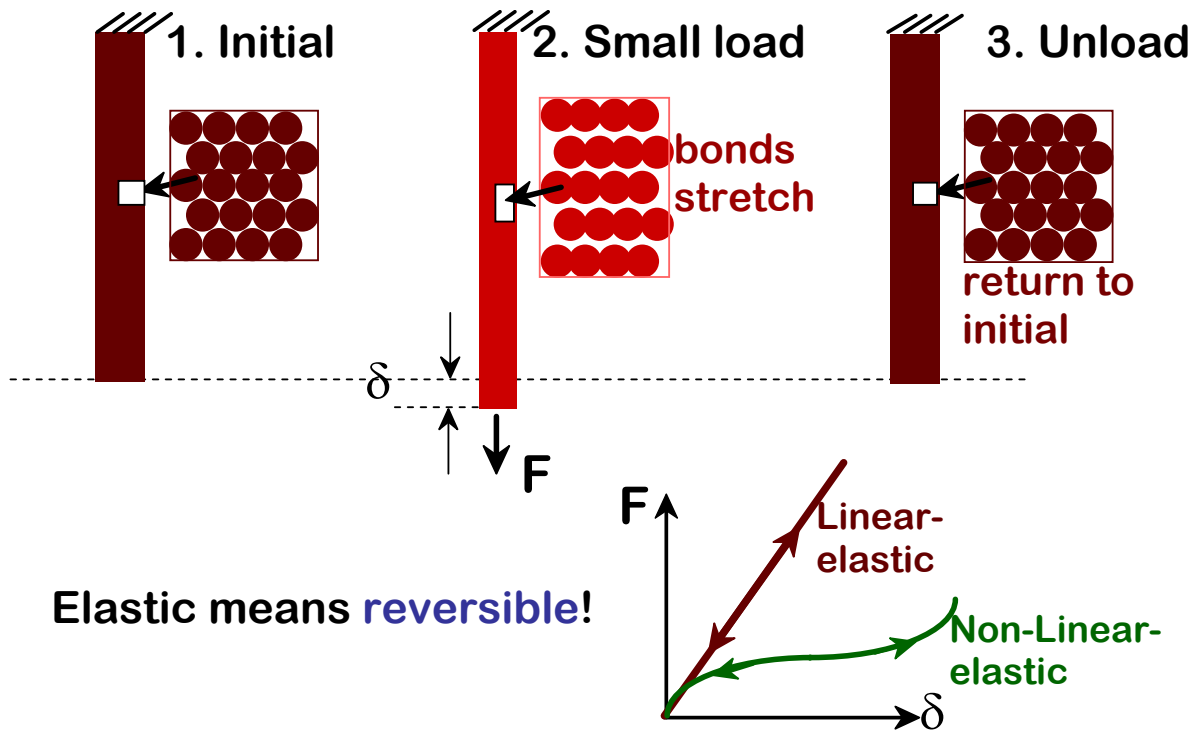
*strain*  $\epsilon = \frac{\text{increase in length } x}{\text{original length } L}$

$$\sigma_{ij} = C_{ij} \cdot \epsilon_{ij}$$

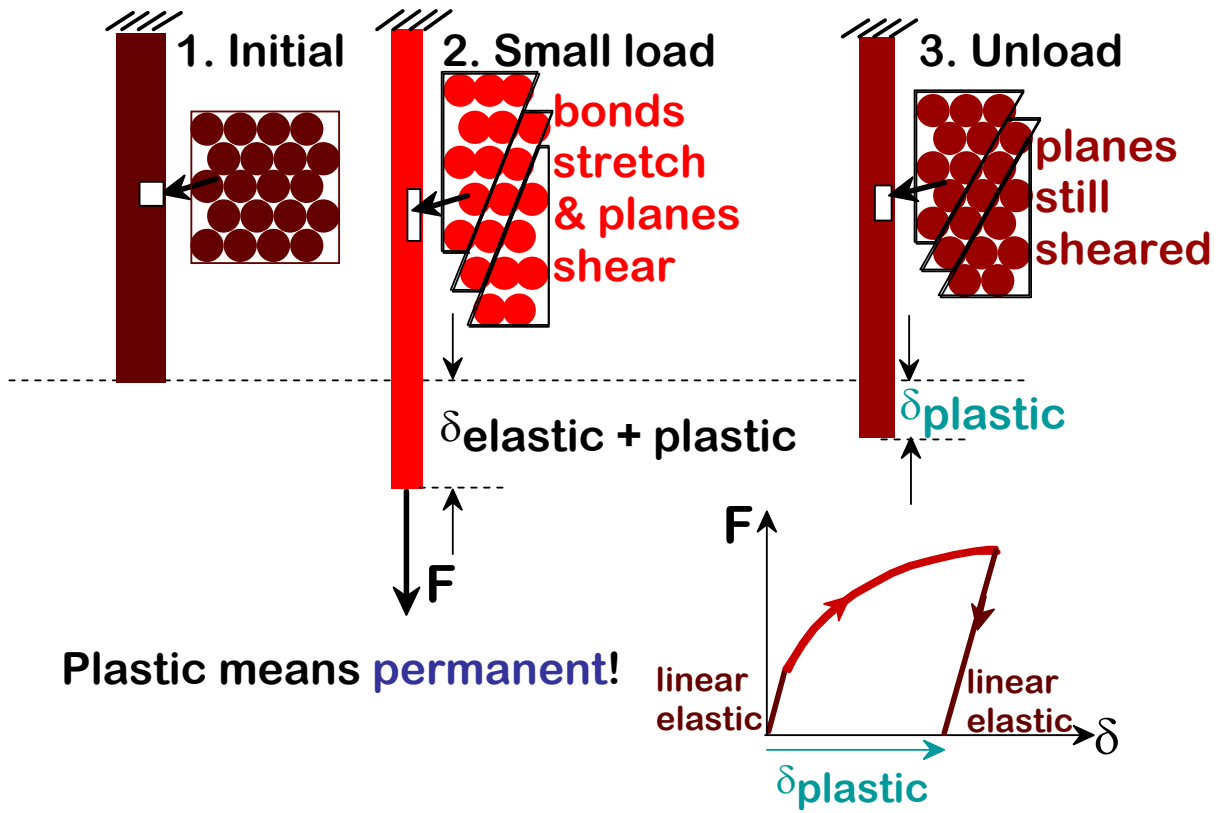
See Eqs 37 and 38 in Kittel, p.77

Is there a good reason to introduce complications with so many different indexes as in p.73-80? Yes it is, because, for example elastic waves in crystals often propagate in different directions, specifically can be longitudinal or transverse (shear) waves

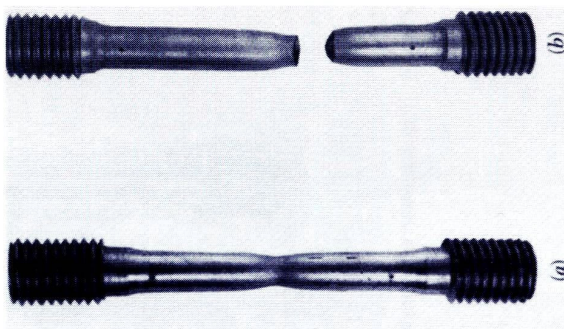
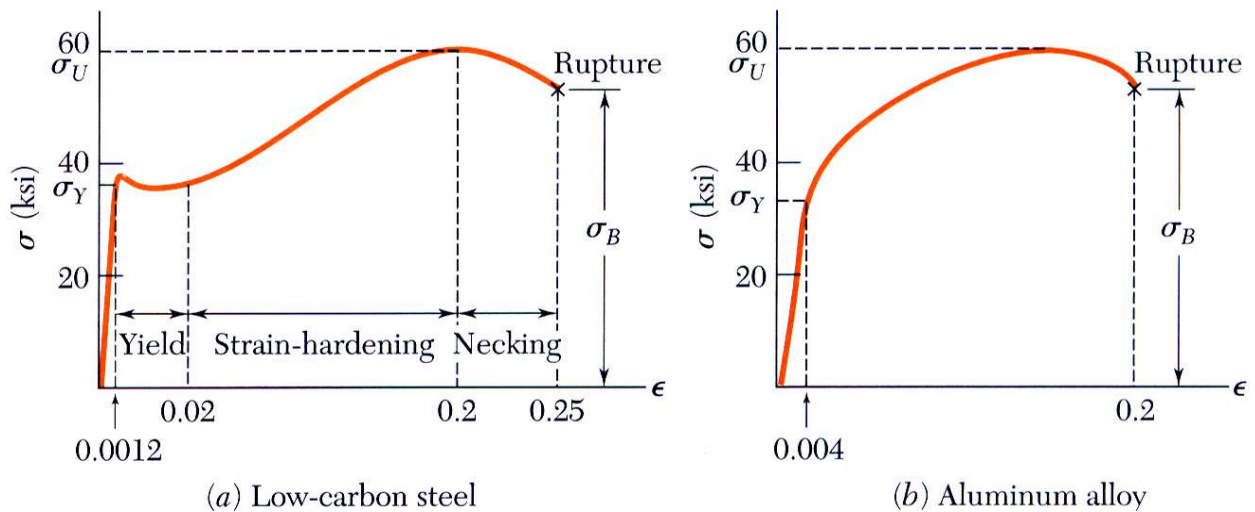
### Elastic strain (deformation)



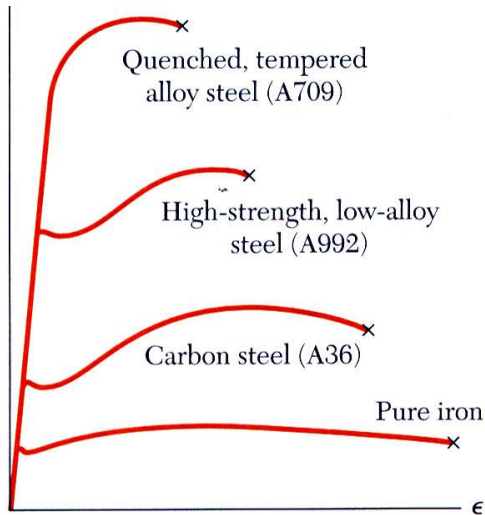
## Plastic deformation



## Elastic/plastic deformation



## Elastic/plastic deformation



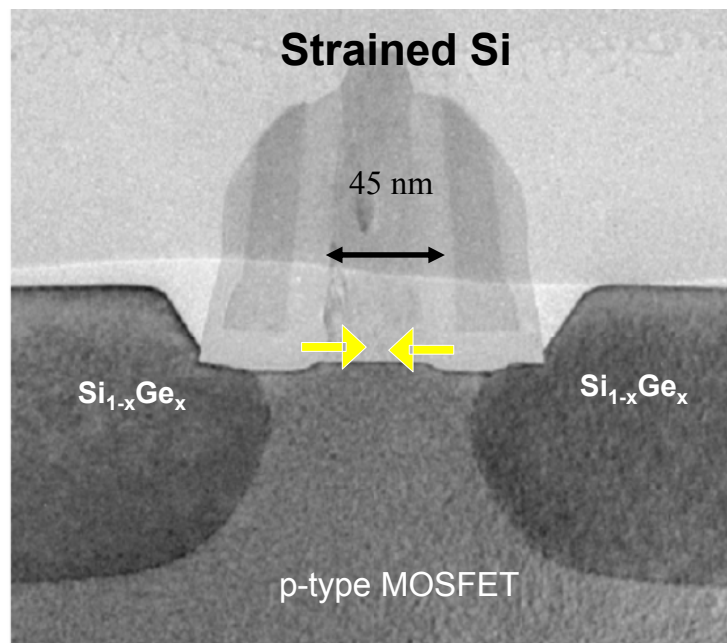
- Below the yield stress

$$\sigma = E\epsilon$$

$E$  = Youngs Modulus or  
Modulus of Elasticity

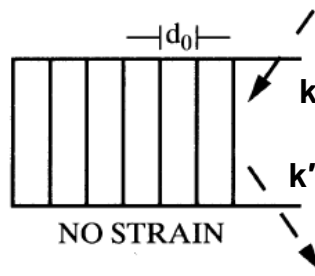
- Strength is affected by alloying, heat treating, and manufacturing process but stiffness (Modulus of Elasticity) is not.

## How is strain applied to the electronic chips?

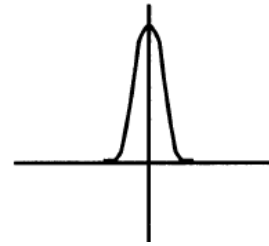


## Evolution of x-ray diffraction

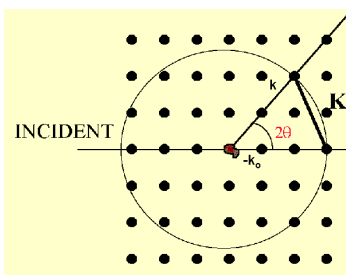
CRYSTAL LATTICE



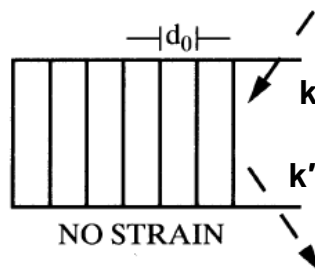
DIFFRACTION LINE



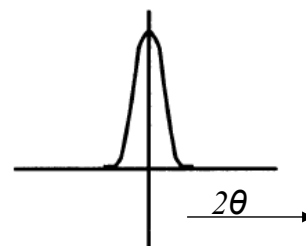
## Evolution of x-ray diffraction



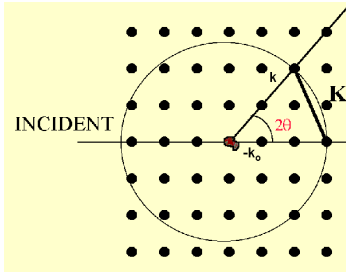
CRYSTAL LATTICE



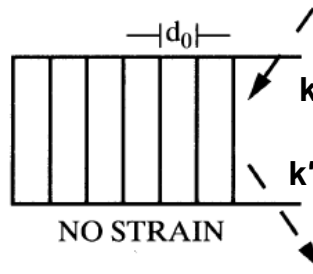
DIFFRACTION LINE



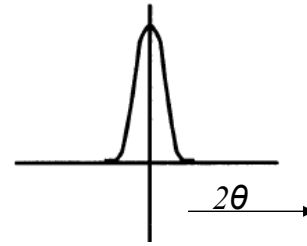
## Evolution of x-ray diffraction



CRYSTAL LATTICE

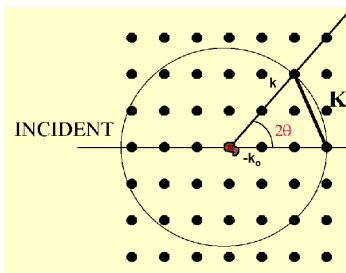


DIFFRACTION LINE

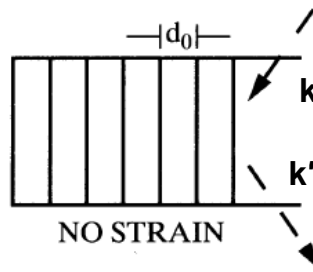


UNIFORM STRAIN

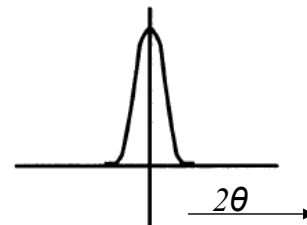
## Evolution of x-ray diffraction



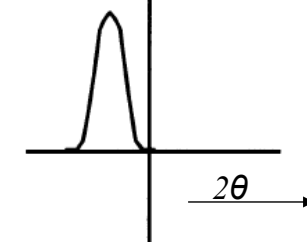
CRYSTAL LATTICE



DIFFRACTION LINE

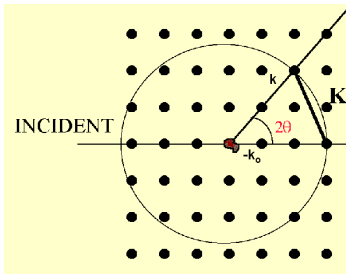


UNIFORM STRAIN



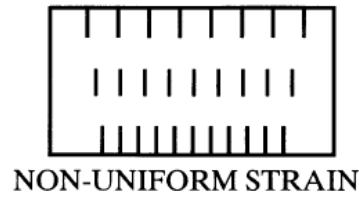
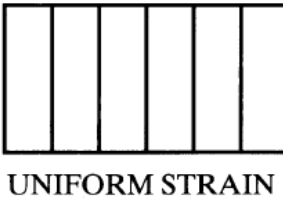
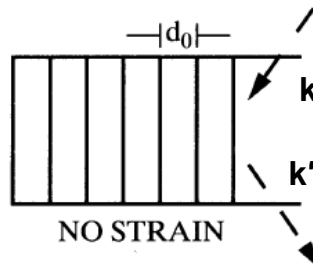
$$n\lambda = 2d \sin \theta$$

## Evolution of x-ray diffraction

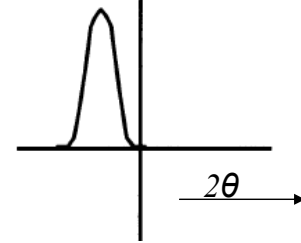
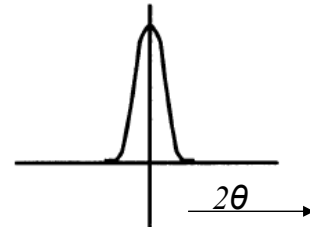


$$n\lambda = 2d \sin \theta$$

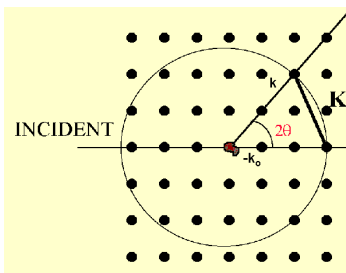
CRYSTAL LATTICE



DIFFRACTION LINE

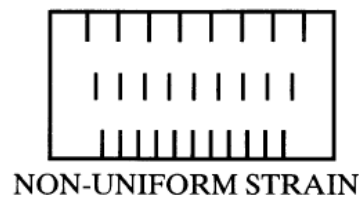
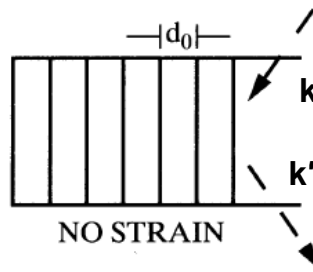


## Evolution of x-ray diffraction

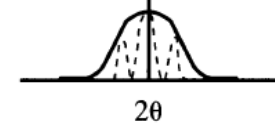
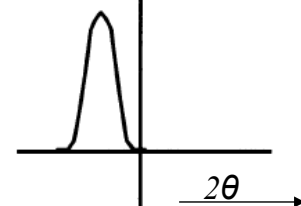
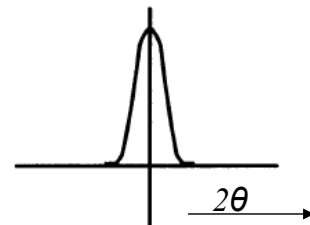


$$n\lambda = 2d \sin \theta$$

CRYSTAL LATTICE



DIFFRACTION LINE





## Lecture 5: Crystal binding, elastic strain and waves

- ionization and cohesive energies in the context of periodic table;
- interaction between two atoms in terms of attraction and repulsion forces;
- analysis of elastic strain
- elastic waves in cubic crystals

### Elastic Waves in Cubic Crystals

Newton's 2<sup>nd</sup> law: 
$$\rho \frac{\partial^2 u_i}{\partial t^2} = \frac{\partial \sigma_{ik}}{\partial x_k}$$

$$\sigma_{ik} = C_{ikjl} u_{jl} \quad \rightarrow \quad \rho \frac{\partial^2 u_i}{\partial t^2} = C_{ikjl} \frac{\partial u_{jl}}{\partial x_k} = \frac{1}{2} C_{ikjl} \left( \frac{\partial^2 u_l}{\partial x_k \partial x_j} + \frac{\partial^2 u_j}{\partial x_k \partial x_l} \right) = C_{ikjl} \frac{\partial^2 u_l}{\partial x_k \partial x_j}$$

$$\begin{aligned} \rho \frac{\partial^2 u_1}{\partial t^2} &= C_{1111} \frac{\partial^2 u_1}{\partial x_1^2} + C_{1122} \frac{\partial^2 u_2}{\partial x_1 \partial x_2} + C_{1133} \frac{\partial^2 u_3}{\partial x_1 \partial x_3} + C_{1212} \frac{\partial^2 u_2}{\partial x_2 \partial x_1} + C_{1221} \frac{\partial^2 u_1}{\partial x_2^2} + C_{1313} \frac{\partial^2 u_3}{\partial x_3 \partial x_1} + C_{1331} \frac{\partial^2 u_1}{\partial x_3^2} \\ &= C_{1111} \frac{\partial^2 u_1}{\partial x_1^2} + C_{1122} \left( \frac{\partial^2 u_2}{\partial x_1 \partial x_2} + \frac{\partial^2 u_3}{\partial x_1 \partial x_3} \right) + C_{1212} \left( \frac{\partial^2 u_2}{\partial x_2 \partial x_1} + \frac{\partial^2 u_1}{\partial x_2^2} + \frac{\partial^2 u_3}{\partial x_3 \partial x_1} + \frac{\partial^2 u_1}{\partial x_3^2} \right) \end{aligned}$$

$$\therefore \quad \rho \frac{\partial^2 u_1}{\partial t^2} = C_{11} \frac{\partial^2 u_1}{\partial x_1^2} + (C_{12} + C_{44}) \left( \frac{\partial^2 u_2}{\partial x_1 \partial x_2} + \frac{\partial^2 u_3}{\partial x_1 \partial x_3} \right) + C_{44} \left( \frac{\partial^2 u_1}{\partial x_2^2} + \frac{\partial^2 u_1}{\partial x_3^2} \right)$$

Similarly 
$$\rho \frac{\partial^2 u_2}{\partial t^2} = C_{11} \frac{\partial^2 u_2}{\partial x_2^2} + (C_{12} + C_{44}) \left( \frac{\partial^2 u_3}{\partial x_2 \partial x_3} + \frac{\partial^2 u_1}{\partial x_2 \partial x_1} \right) + C_{44} \left( \frac{\partial^2 u_2}{\partial x_1^2} + \frac{\partial^2 u_2}{\partial x_3^2} \right)$$

$$\rho \frac{\partial^2 u_3}{\partial t^2} = C_{11} \frac{\partial^2 u_3}{\partial x_3^2} + (C_{12} + C_{44}) \left( \frac{\partial^2 u_2}{\partial x_3 \partial x_2} + \frac{\partial^2 u_1}{\partial x_3 \partial x_1} \right) + C_{44} \left( \frac{\partial^2 u_3}{\partial x_2^2} + \frac{\partial^2 u_3}{\partial x_1^2} \right)$$

## Dispersion Equation

$$\rho \frac{\partial^2 u_i}{\partial t^2} = C_{ijkl} \frac{\partial^2 u_l}{\partial x_k \partial x_j}$$

$$u_i = u_{0i} e^{i(\mathbf{k}\cdot\mathbf{r} - \omega t)} \quad \rightarrow \quad \omega^2 \rho u_{0i} = C_{ijkl} k_k k_j u_{0l}$$

$$(\omega^2 \rho \delta_{il} - C_{ijkl} k_k k_j) u_{0l} = 0$$

$$\left| \omega^2 \rho \delta_{il} - C_{ijkl} k_k k_j \right| = 0 \quad \text{dispersion equation}$$

$$\left| \omega^2 \rho I - \mathbf{C}(\mathbf{k}) \right| = 0 \quad \mathbf{C}_{ij}(\mathbf{k}) = C_{imnj} k_m k_n$$

## Waves in the [100] direction

$$\left| \omega^2 \rho I - \mathbf{C}(\mathbf{k}) \right| = 0 \quad \mathbf{C}_{ij}(\mathbf{k}) = C_{imnj} k_m k_n$$

$$\mathbf{k} = k(1, 0, 0) \quad \rightarrow \quad \mathbf{C}_{ij}(\mathbf{k}) = C_{i11j} k^2$$

$$\mathbf{C}(\mathbf{k}) = k^2 \begin{pmatrix} C_{1111} & 0 & 0 \\ 0 & C_{2112} & 0 \\ 0 & 0 & C_{3113} \end{pmatrix} = k^2 \begin{pmatrix} C_{11} & 0 & 0 \\ 0 & C_{44} & 0 \\ 0 & 0 & C_{44} \end{pmatrix}$$

$$\omega_L = \sqrt{\frac{C_{11}}{\rho}} k \quad \mathbf{u}_0 = (1, 0, 0) \quad \text{Longitudinal}$$

$$\omega_T = \sqrt{\frac{C_{44}}{\rho}} k \quad \mathbf{u}_0 = (0, 1, 0) \\ \mathbf{u}_0 = (0, 0, 1) \quad \text{Transverse, degenerate}$$

## Waves in the [110] direction

$$|\omega^2 \rho I - C(\mathbf{k})| = 0 \quad C_{ij}(\mathbf{k}) = C_{imnj} k_m k_n$$

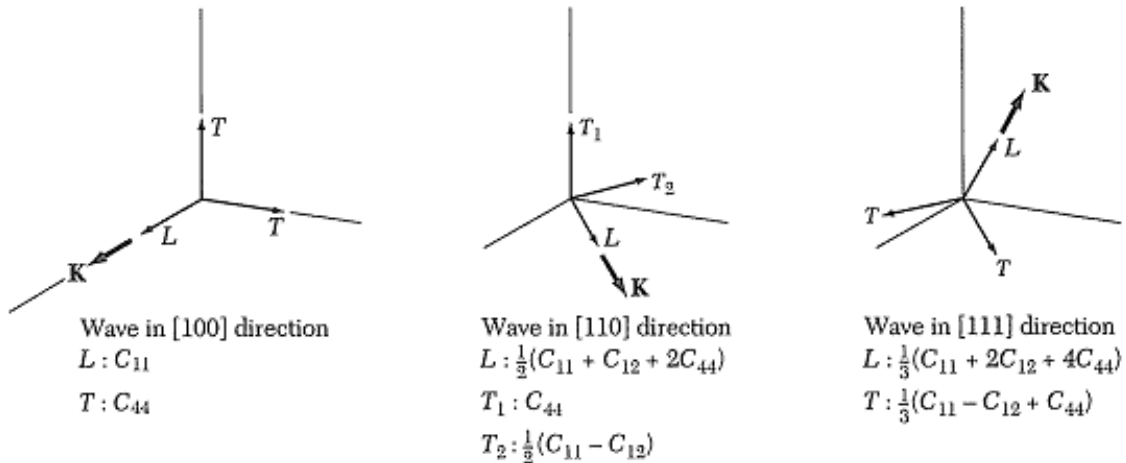
$$\mathbf{k} = \frac{k}{\sqrt{2}}(1,1,0) \quad \rightarrow \quad C_{ij}(\mathbf{k}) = (C_{i11j} + C_{i12j} + C_{i21j} + C_{i22j}) \frac{k^2}{2}$$

$$C(\mathbf{k}) = \frac{k^2}{2} \begin{pmatrix} C_{1111} + C_{1221} & C_{1122} + C_{1212} & 0 \\ C_{2121} + C_{2211} & C_{2112} + C_{2222} & 0 \\ 0 & 0 & C_{3113} + C_{3223} \end{pmatrix} = \frac{k^2}{2} \begin{pmatrix} C_{11} + C_{44} & C_{12} + C_{44} & 0 \\ C_{12} + C_{44} & C_{11} + C_{44} & 0 \\ 0 & 0 & 2C_{44} \end{pmatrix}$$

$$\omega_L = \sqrt{\frac{1}{2\rho}(C_{11} + C_{12} + 2C_{44})} k \quad \mathbf{u}_0 = (1,1,0) \quad \text{Longitudinal}$$

$$\omega_{T1} = \sqrt{\frac{1}{2\rho}(C_{11} - C_{12})} k \quad \mathbf{u}_0 = (1,-1,0) \quad \text{Transverse}$$

$$\omega_{T2} = \sqrt{\frac{C_{44}}{\rho}} k \quad \mathbf{u}_0 = (0,0,1) \quad \text{Transverse}$$



**Figure 20** Effective elastic constants for the three modes of elastic waves in the principal propagation directions in cubic crystals. The two transverse modes are degenerate for propagation in the [100] and [111] directions.