

# FYS3410 - Vår 2010 (Kondenserte fasers fysikk)

<http://www.uio.no/studier/emner/matnat/fys/FYS3410/index-eng.xml>

Based on Introduction to Solid State Physics by Kittel

## Course content

- Periodic structures, understanding of diffraction experiment and reciprocal lattice
- Crystal binding, elastic strain and waves
- Imperfections in crystals: diffusion, point defects, dislocations
- Crystal vibrations: phonon heat capacity and thermal conductivity
- Free electron Fermi gas: density of states, Fermi level, and electrical conductivity
- Electrons in periodic potential: energy bands theory classification of metals, semiconductors and insulators
- Semiconductors: band gap, effective masses, charge carrier distributions, doping, pn-junctions
- Metals: Fermi surfaces, temperature dependence of electrical conductivity

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## FYS3410 lecture schedule and exams: Spring 2010

M/18/1/2010:	Introduction and motivation. Periodicity and lattices	2h
W/20/1/2010:	Index system for crystal planes. Crystal structures	1h
M/25/1/2010:	Reciprocal space, Laue condition and Ewald construction	2h
W/27/1/2010:	Brillouin Zones. Interpretation of a diffraction experiment	1h
M/01/2/2010:	Crystal binding, elastic strain and waves	2h
W/03/2/2010:	Elastic waves in cubic crystals; defects in crystals	1h
M/08/2/2010:	Defects in crystals; case study - vacancies	2h
W/10/2/2010:	Atomic diffusion	1h
M/15/2/2010:	Crystal vibrations and phonons	2h
W/17/2/2010:	Crystal vibrations and phonons	1h
M/22/2/2010:	Planck distribution and density of states	2h
W/24/2/2010:	Debye model	1h
M/01/3/2010:	Einstein model and general result for density of states	2h
W/03/3/2010:	Thermal conductivity	1h
M/08/3/2010:	Free electron Fermi gas in 1D and 3D – ground state	2h
W/10/3/2010:	Density of states, effect of temperature – FD distribution	1h
M/15/3/2010:	Heat capacity of FEFG	2h
W/17/3/2010:	Repetition	1h
22/3/2010:	Mid-term exam	

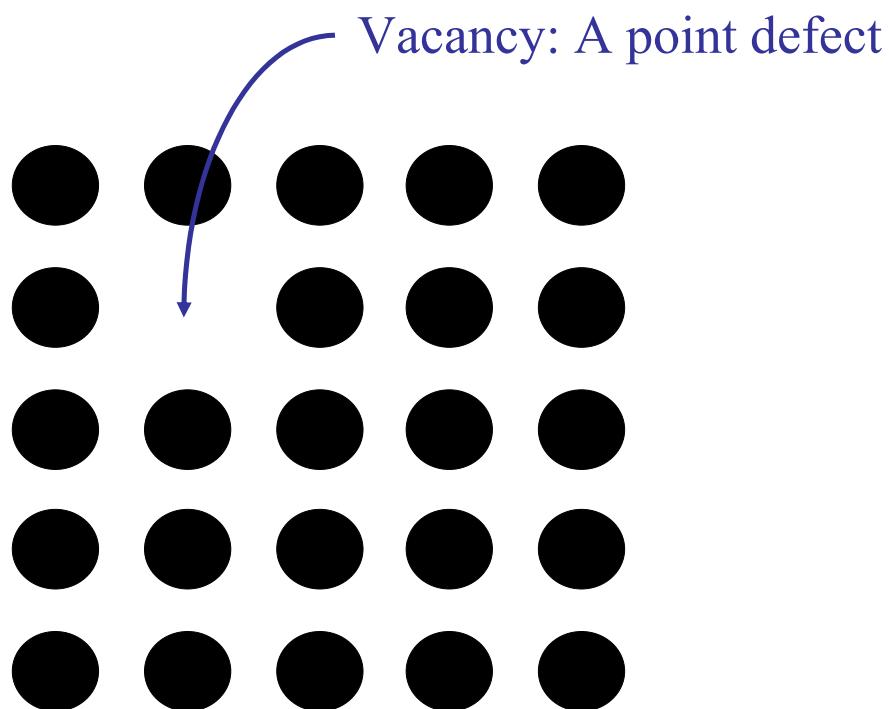
<b>M/14/4/2010:</b>	<b>Electrical and thermal conductivity in metals</b>	<b>2h</b>
<b>W/12/4/2010:</b>	<b>Bragg reflection of electron waves at the boundary of BZ</b>	<b>1h</b>
<b>M/19/4/2010:</b>	<b>Energy bands, Kronig - Penny model</b>	<b>2h</b>
<b>W/21/4/2010:</b>	<b>Empty lattice approximation; number of orbitals in a band</b>	<b>1h</b>
<b>M/26/4/2010:</b>	<b>Semiconductors, effective mass method, intrinsic carriers</b>	<b>2h</b>
<b>W/28/4/2010:</b>	<b>Impurity states in semiconductors and carrier statistics</b>	<b>1h</b>
<b>M/03/5/2010:</b>	<b>p-n junctions and heterojunctions</b>	<b>2h</b>
<b>W/05/5/2010:</b>	<b>surface structure, surface states, Schottky contacts</b>	<b>2h</b>
<b>M/10/5/2010:</b>	<b>no lectures</b>	
<b>W/12/5/2010:</b>	<b>no lectures</b>	
<b>W/19/5/2010:</b>	<b>Repetition</b>	<b>2h</b>
<b>W/26/5/2010:</b>	<b>Repetition</b>	<b>2h</b>
<b>28/5/2010:</b>	<b>Final Exam (sensor: Prof. Arne Nylandsted Larsen at the Århus University, Denmark, <a href="http://person.au.dk/en/anl@phys.au.dk">http://person.au.dk/en/anl@phys.au.dk</a>)</b>	

### Lecture 7: Defects in crystals; case study - vacancies

- defects in crystals; vacancies
- Configurational entropy due to vacancy
- Equilibrium concentration of vacancies – temperature dependence
- Equilibrium concentration of vacancy – pressure dependence
- Excess concentration and clustering of vacancies
- Clustering of ion implantation induced vacancies in ZnO

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<i>Defects</i>	<i>Dimensionality</i>	<i>Examples</i>
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Point	0	Vacancy
Line	1	Dislocation
Surface	2	Free surface, Grain boundary

### **Vacancies**

#### Fact

There *may* be vacant sites in a crystal

#### *Surprising Fact*

There *must* be a certain fraction of vacant sites in a crystal in *equilibrium*.

## Vacancies

- Crystal in equilibrium
- Minimum Gibbs free energy  $G$  at constant  $T$  and  $P$
- A certain concentration of vacancy lowers the free energy of a crystal

## Vacancies

Gibbs free energy  $G$  involves two terms:

1. Enthalpy  $H = E + PV$        $E$  internal energy

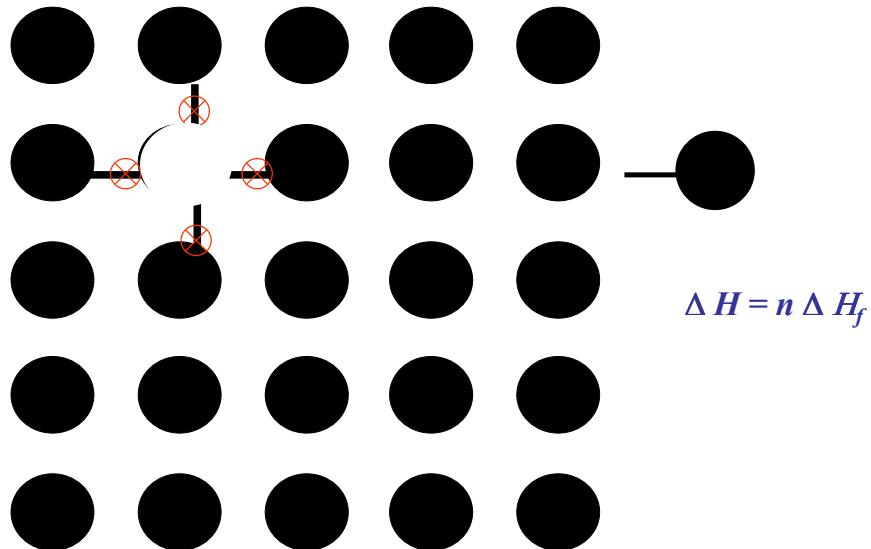
$P$  pressure

$V$  volume

2. Entropy  $S = k \ln W$        $k$  Boltzmann constant  
                                         $W$  number of microstates

$$G = H - TS \quad T \text{ Absolute temperature}$$

## Vacancies

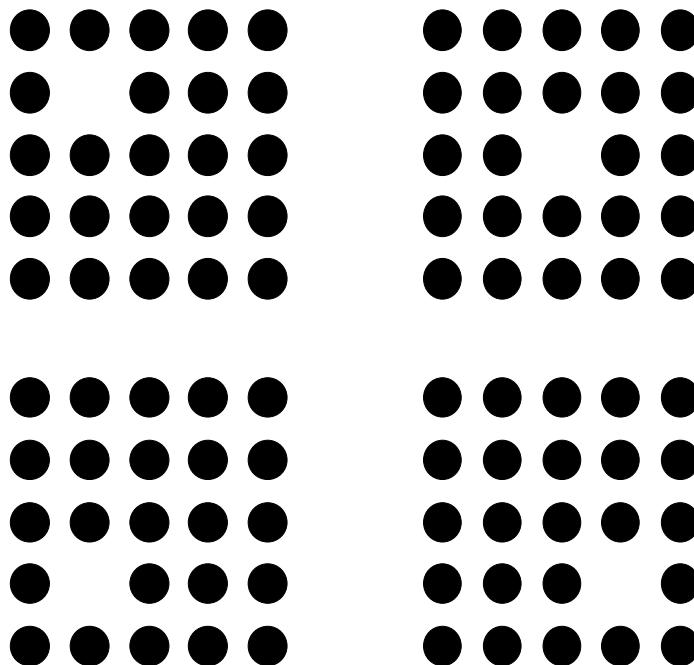


Vacancy increases  $H$  of the crystal due to energy required to break bonds

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## Configurational entropy due to vacancy



## Configurational entropy due to vacancy

Number of atoms:  $N$

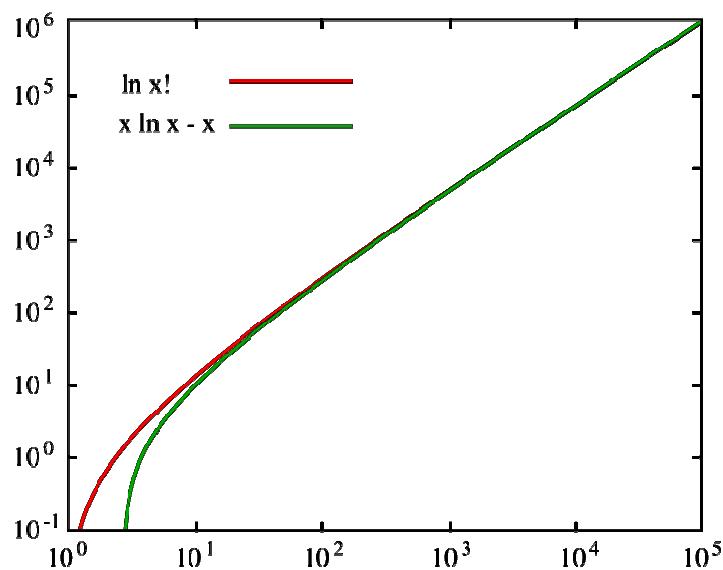
Number of vacacies:  $n$

Total number of sites:  $N+n$

How many distinguished configurations,  
so called microstates?

We calculate this explicitly

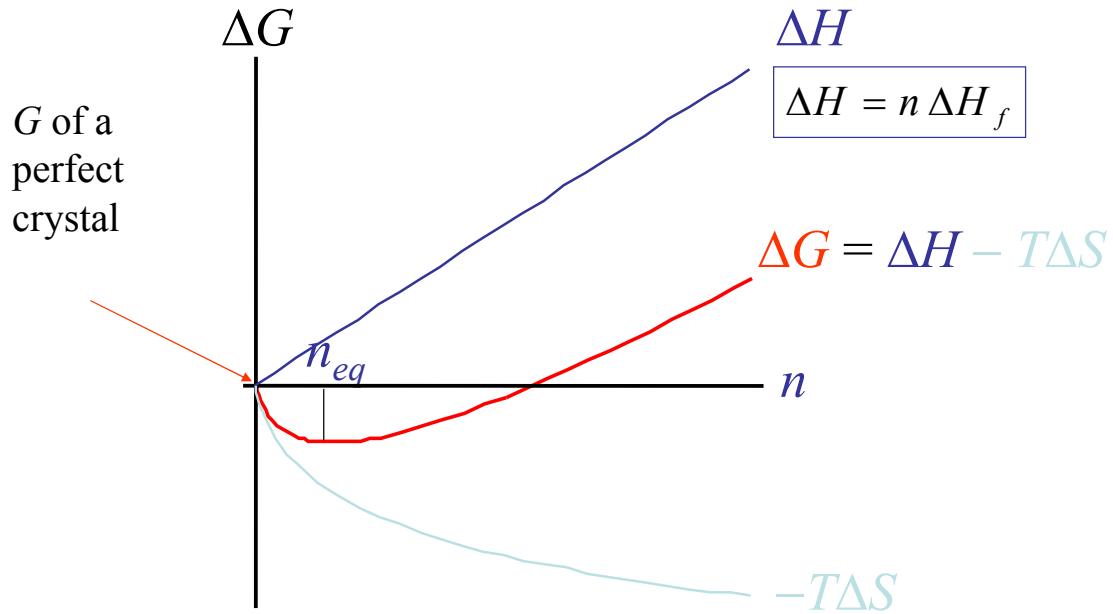
## Configurational entropy due to vacancy



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## Equilibrium concentration of vacancies – temperature dependence

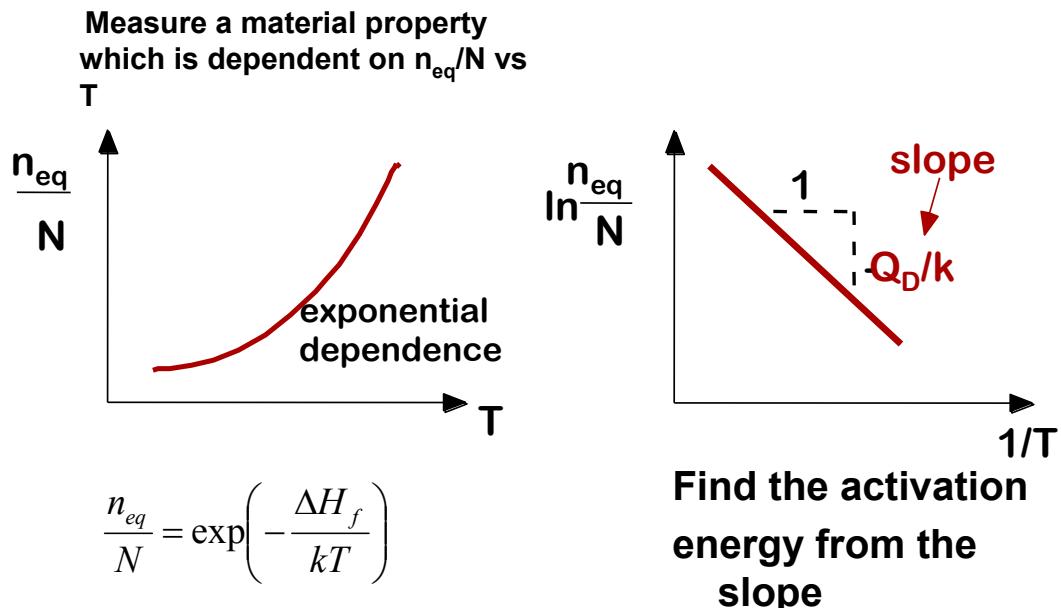


## Equilibrium concentration of vacancies – temperature dependence

$$\left. \frac{\partial \Delta G}{\partial n} \right|_{n=n_{eq}} = 0$$

$$\frac{n_{eq}}{N} = \exp\left(-\frac{\Delta H_f}{kT}\right)$$

## Equilibrium concentration of vacancies – temperature dependence



## Equilibrium concentration of vacancies – temperature dependence

– Copper at 1000 °C

$$H_f = 0.9 \text{ eV/at} \quad A_{Cu} = 63.5 \text{ g/mol} \quad \rho = 8400 \text{ kg/m}^{-3}$$

First find N in atoms/m<sup>-3</sup>

$$N = \frac{N_A \rho}{A_{Cu}} = \frac{(6.023 \times 10^{23})(8400)}{0.0635}$$

$$N =$$

units Check

$$N \rightarrow \frac{(at/mol)(kg/m^3)}{kg/mol} = \frac{at}{m^3}$$

## Equilibrium concentration of vacancies – temperature dependence

$$N = 7.97 \times 10^{28} \text{ at-sites / m}^3$$

- Now apply the Arrhenius relation @1000 °C

$$\begin{aligned} N_v &= N \exp\left(-\frac{H_f}{kT}\right) \\ &= 7.97 \times 10^{28} \exp\left[\frac{-0.9eV/at}{(8.62 \times 10^{-5} eV/at - K)1273K}\right] \\ N_v &= 2.18 \times 10^{25} \text{ vac / m}^3 \end{aligned}$$

## Equilibrium concentration of vacancies – temperature dependence

$$\frac{n_{eq}}{N} = \exp\left(-\frac{\Delta H_f}{kT}\right)$$

Al:  $\Delta H_f = 0.70 \text{ ev/vacancy}$

Ni:  $\Delta H_f = 1.74 \text{ ev/vacancy}$

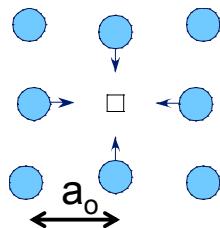
$n/N$	0 K	300 K	900 K
Al	0	$1.45 \times 10^{-12}$	$1.12 \times 10^{-4}$
Ni	0	$5.59 \times 10^{-30}$	$1.78 \times 10^{-10}$

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### Equilibrium concentration of vacancy – pressure dependence

$$H_f = E_f + PV_f \quad \longrightarrow \quad \Delta G_f = E_f + PV_f - TS_f$$

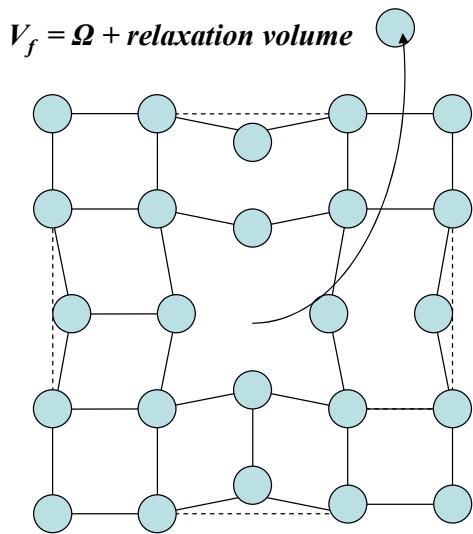


- Neighboring atoms tend to move into the vacancy, which creates a tensile stress field
- The stress/strain field is nearly spherical and short-range.

$$C_V^{eq} = e^{-\Delta G_f / kT} = e^{S_f / k} e^{-E_f / kT} e^{\sigma V_f / kT}$$

## Equilibrium concentration of vacancy – pressure dependence

$$C_V^{eq} = e^{-\Delta G_f/kT} = e^{S_f/k} e^{-E_f/kT} e^{\sigma V_f/kT}$$



How big the pressure should be to make a measurable effect on vacancy concentration?

Compare

$$\Delta H_f \quad \sigma V_f$$

## Equilibrium concentration of vacancy – pressure dependence

101.325 kPa is “one standard atmosphere” and **1 Pa = 1 N/m<sup>2</sup>**

$$1 \text{ J} = 1 \text{ N} \cdot \text{m} = \left( \frac{\text{kg} \cdot \text{m}}{\text{s}^2} \right) \cdot \text{m} = \frac{\text{kg} \cdot \text{m}^2}{\text{s}^2} = \text{Pa} \cdot \text{m}^3 = 1 \text{ W} \cdot \text{s}$$

$$1 \text{ eV} = 1.602176487 \times 10^{-19} \text{ Joule}$$

As we calculate the effect of pressure/stress on vacancy concentration starts to be significant at quite high values – in the range of 100 MPa.

Are these conditions available in real “life” or happens only in a laboratory experiment?

## Excess concentration and clustering of vacancies

### Radiation

### Chemical reactions

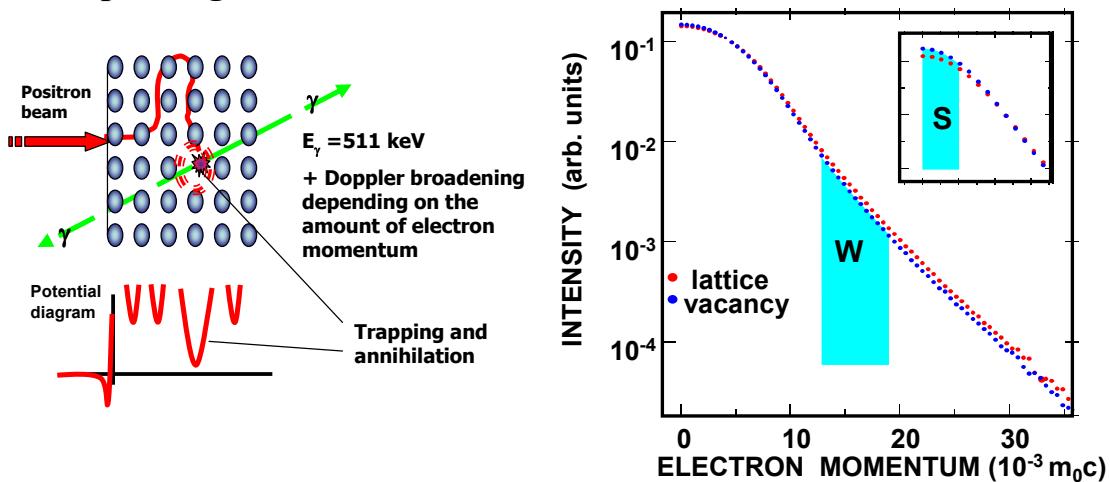
The amount of single vacancies exceed the "solubility" limit and vacancies may start cluster

$$\frac{\partial [V_n]}{\partial t} = \underbrace{4\pi R D_V [V_{n-1}] [V]}_{\text{generation}} - \underbrace{[V_n] \cdot C_0 \exp(-E_b/kT)}_{\text{dissociation}}$$

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## Positron probing of vacancies in semiconductors



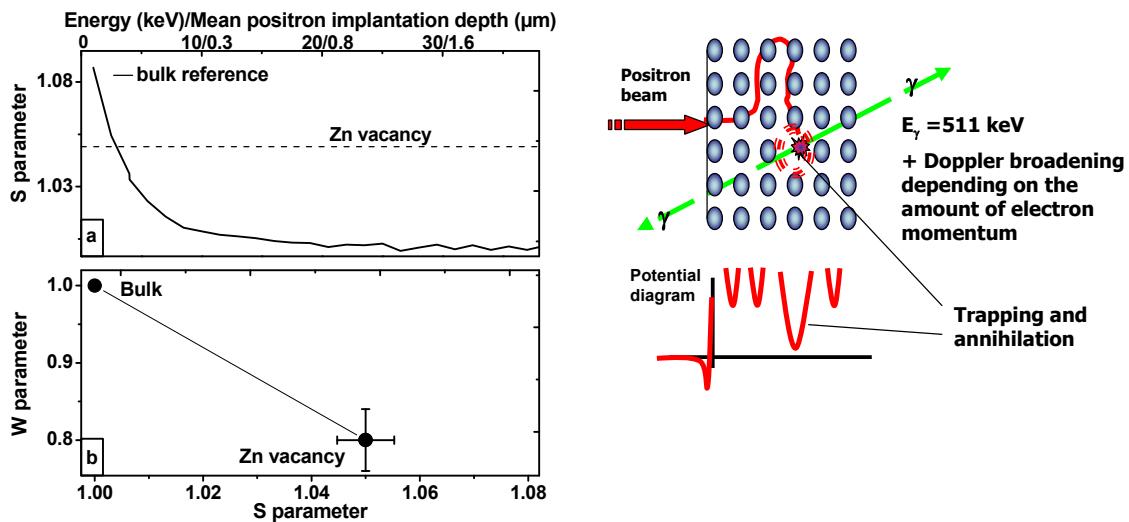
S-parameter characterizes annihilation with low momentum valence electrons. Increase in S-parameter is naturally interpreted as an increase in vacancy concentration

W-parameter characterizes annihilation with high momentum core electrons and increase in vacancy concentration results in decrease of W-parameter

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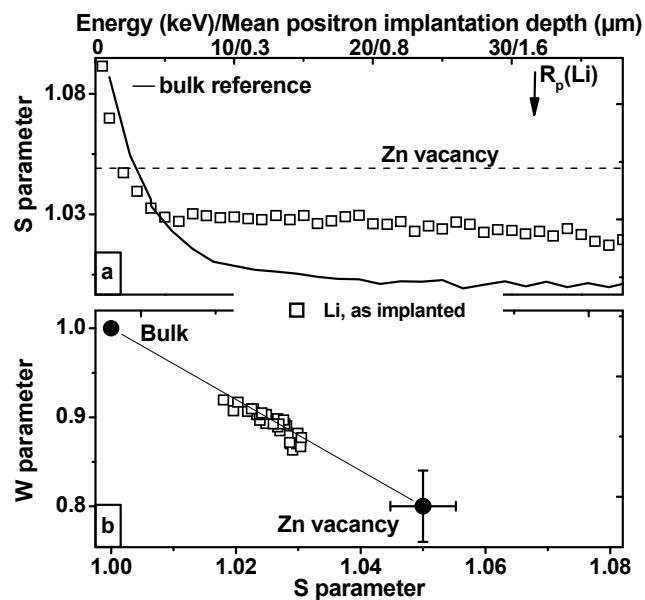
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## Clustering of ion implantation induced vacancies in ZnO

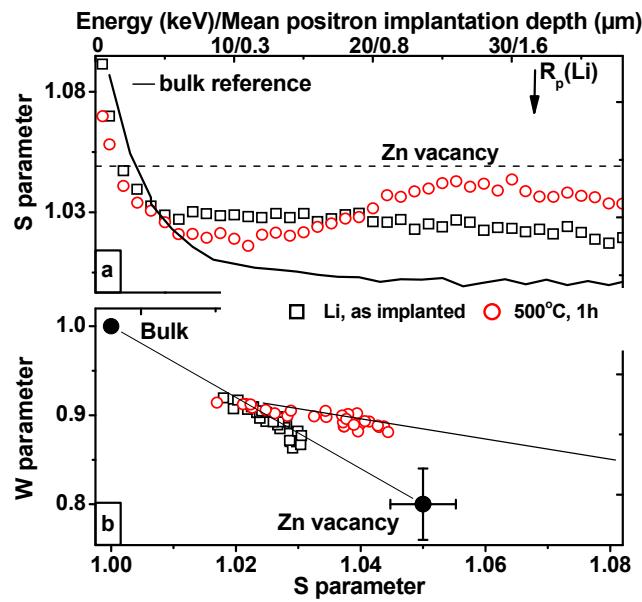


**Experimental points group around a line in the W-S plane if there are only two annihilation states available in the sample**

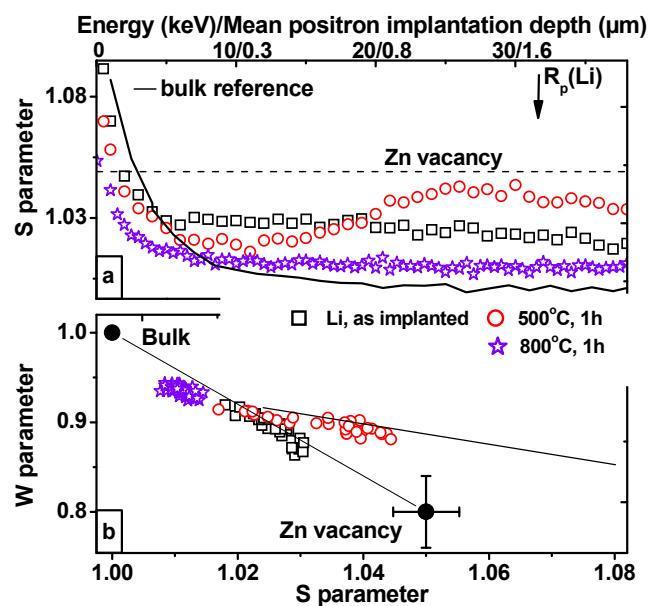
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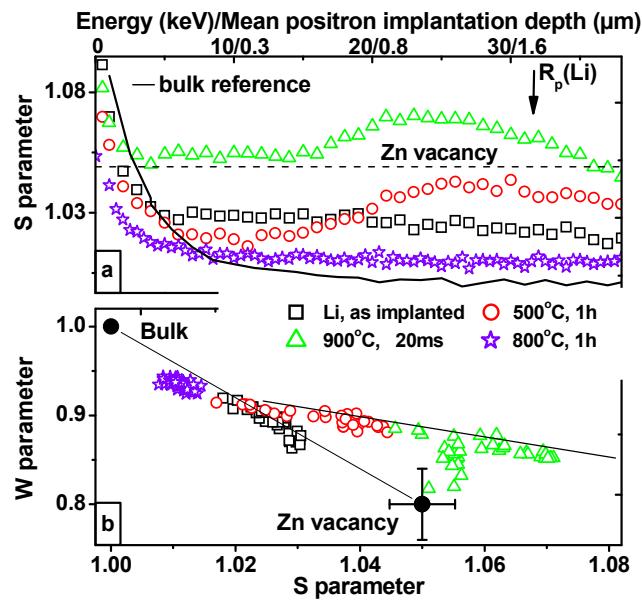
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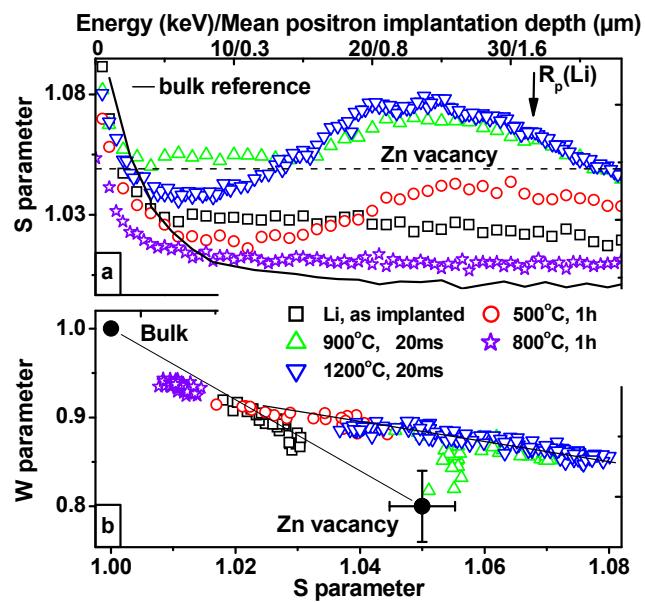
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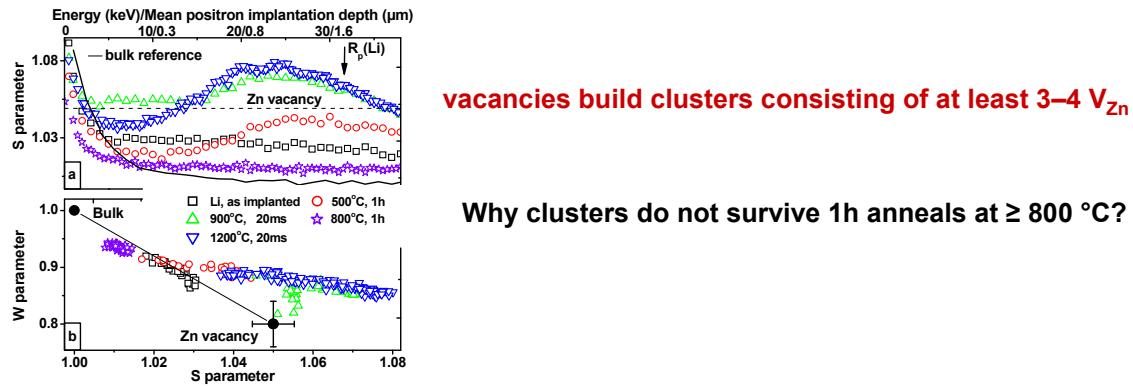
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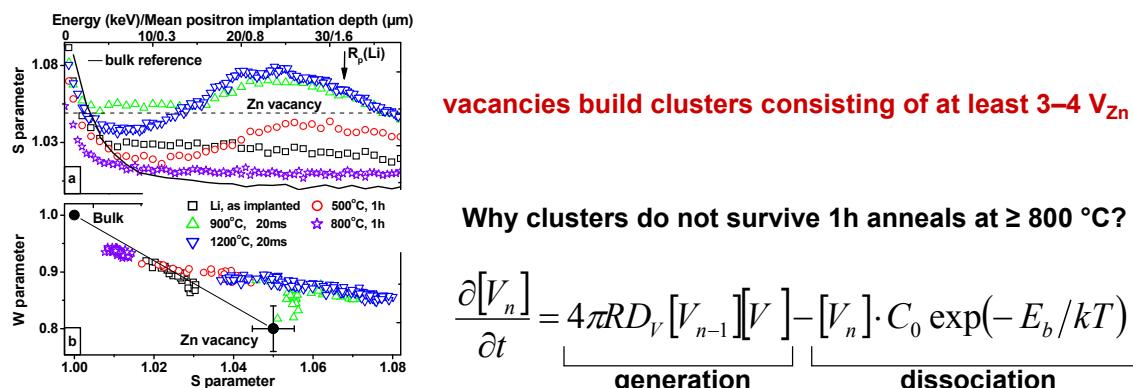
## Clustering of ion implantation induced vacancies in ZnO



vacancies build clusters consisting of at least 3–4  $V_{\text{Zn}}$

Why clusters do not survive 1h anneals at  $\geq 800 \text{ }^{\circ}\text{C}$ ?

## Clustering of ion implantation induced vacancies in ZnO

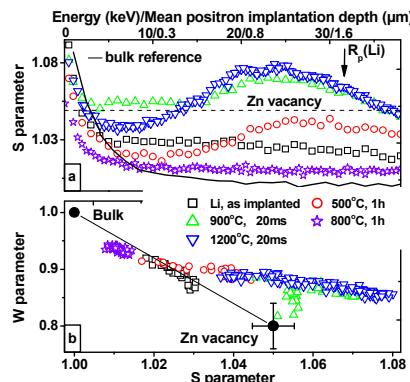


vacancies build clusters consisting of at least 3–4  $V_{\text{Zn}}$

Why clusters do not survive 1h anneals at  $\geq 800 \text{ }^{\circ}\text{C}$ ?

- clustering take place as long as the vacancy diffusivity ( $D_V$ ), supersaturation level ( $[V]$ ), and clustering reaction radii ( $R$ ) are high enough
- $E_b$  – dissociation energy – determines the dissociation rate at a given temperature

# Clustering of ion implantation induced vacancies in ZnO



**vacancies build clusters consisting of at least 3–4 V<sub>Zn</sub>**

## Why clusters do not survive 1h anneals at $\geq 800$ °C?

$$\frac{\partial [V_n]}{\partial t} = \underbrace{4\pi RD_V[V_{n-1}][V]}_{\text{generation}} - \underbrace{[V_n] \cdot C_0 \exp(-E_b/kT)}_{\text{dissociation}}$$

- clustering take place as long as the vacancy diffusivity ( $D_v$ ), supersaturation level ([V]), and clustering reaction radii (R) are high enough
  - $E_b$  – dissociation energy – determines the dissociation rate at a given temperature

temperature (°C)	500	800	900	1200	1400
clusters survive	1h		20 ms	20 ms	20 ms
clusters dissolve		1h			

$E_b = 2.6 \pm 0.3$  eV

T.M. Børseth, et al Phys. Rev. B 89, (2006)