

UNIVERSITETET I OSLO

Det matematisk-naturvitenskapelige fakultet

Exam in: FYS 4110 Non-relativistic quantum mechanics

Day of exam: Thursday, December 4, 2008

Exam hours: 3 hours, beginning at 14:30

This examination paper consists of 3 pages

Permitted materials: Calculator

Øgrim og Lian or Angell og Lian: Størrelser og enheter i fysikken

Rottmann: Matematisk formelsamling

Language: The solutions may be written in Norwegian or English depending on your own preference.

Make sure that your copy of this examination paper is complete before you begin.

PROBLEM 1

Spin half particle in a harmonic oscillator potential

A spin half particle is moving in a one-dimensional harmonic oscillator potential (in the x -direction) under the influence of a constant magnetic field (in the z -direction). The Hamiltonian is

$$\hat{H} = \hbar\omega_0(\hat{a}^\dagger\hat{a} + \frac{1}{2}) + \frac{1}{2}\hbar\omega_1\sigma_z + \lambda\hbar(\hat{a}^\dagger\sigma_- + \hat{a}\sigma_+) \quad (1)$$

where the first term is the harmonic oscillator part with ω_0 as the oscillator frequency, the second term is the spin energy due to the magnetic field, with ω_1 as the spin precession frequency, and the third term is a coupling term between the spin and the position coordinate of the particle. The spin flip operators are defined as $\sigma_\pm = \frac{1}{2}(\sigma_x \pm i\sigma_y)$, \hat{a} and \hat{a}^\dagger are the standard lowering and raising operators of the harmonic oscillator and $\sigma_x, \sigma_y, \sigma_z$ are the Pauli spin matrices.

When $\lambda = 0$, the spin and position of the particle are uncoupled and the energy eigenstates are $|n, m\rangle$ with $n = 0, 1, 2, \dots$ as the harmonic oscillator quantum number and $m = \pm 1$ as the spin quantum number, corresponding to spin up/down along the z -axis. When $\lambda \neq 0$, the unperturbed eigenstates will pairwise be coupled by the Hamiltonian, so that $|n, +1\rangle$ is coupled to $|n+1, -1\rangle$.

a) Consider the two-dimensional subspace spanned by basis vectors $|0, +1\rangle$ and $|1, -1\rangle$. Show that in this space the Hamiltonian takes the form of a 2x2 matrix which can be written as

$$H = \hbar\Delta \begin{pmatrix} \cos\theta & \sin\theta \\ \sin\theta & -\cos\theta \end{pmatrix} + \hbar\epsilon\mathbb{1} \quad (2)$$

with $\mathbb{1}$ as the 2x2 identity matrix. Determine Δ , $\cos\theta$, $\sin\theta$ and ϵ .

b) Find the energies and eigenstates of H in the two-dimensional subspace, expressed as functions of Δ , θ and ϵ .

c) The basis vectors $|n, m\rangle$ can be regarded as tensor products of position and spin vectors, $|n, m\rangle = |n\rangle \otimes |m\rangle$. The two eigenstates found under b) will be entangled with respect to the position and spin variables. Determine the degree of entanglement as function of θ . What value for θ gives the smallest what gives the largest degree of entanglement?

PROBLEM 2

Electric dipole transition

We consider the transition in hydrogen from the excited 2p level to the ground state 1s, where a single photon is emitted. The initial atomic state (A) we assume to have $m = 0$ for the z-component of the orbital angular momentum, so that the quantum numbers of this state are $(n, l, m) = (2, 1, 0)$, with n as the principle quantum number and l as the orbital angular momentum quantum number. Similarly the ground state (B) has quantum numbers $(n, l, m) = (1, 0, 0)$. When expressed in polar coordinates the wave functions of the two states (with intrinsic spin of the electron not included) are given by

$$\begin{aligned}\psi_A(r, \phi, \theta) &= \frac{1}{\sqrt{32\pi a_0^3}} \cos \theta \frac{r}{a_0} e^{-\frac{r}{2a_0}} \\ \psi_B(r, \phi, \theta) &= \frac{1}{\sqrt{\pi a_0^3}} e^{-\frac{r}{a_0}}\end{aligned}\quad (3)$$

where a_0 is the Bohr radius.

We remind you about the form of the interaction matrix element in the dipole approximation,

$$\langle B, 1_{\mathbf{k}a} | \hat{H}_{emis} | A, 0 \rangle = ie \sqrt{\frac{\hbar\omega}{2V\epsilon_0}} \boldsymbol{\epsilon}_{\mathbf{k}a}^* \cdot \mathbf{r}_{BA} \quad (4)$$

where e is the electron charge, \mathbf{k} is the wave vector of the photon, a is the polarization quantum number, ω is the photon frequency and $\boldsymbol{\epsilon}_{\mathbf{k}a}$ is a polarization vector. V is a normalization volume for the electromagnetic wave functions, ϵ_0 is the permittivity of vacuum and \mathbf{r}_{BA} is the matrix element of the electron position operator between the initial and final atomic states.

a) Explain why the x- and y-components of \mathbf{r}_{BA} vanish while the z-component has the form $z_{BA} = \nu a_0$, with ν as a numerical factor. Determine the value of ν . (A useful integration formula is $\int_0^\infty dx x^n e^{-x} = n!$.)

b) To first order in perturbation theory the interaction matrix element (4) determines the direction of the emitted photon, in the form of a probability distribution $p(\phi, \theta)$, where (ϕ, θ) are the polar angles of the wave vector \mathbf{k} . Determine $p(\phi, \theta)$ from the above expressions.

c) The life time of the 2p state is $\tau_{2p} = 1.6 \cdot 10^{-9} s$ while the excited 2s state (with angular momentum $l = 0$) has a much longer life time, $\tau_{2s} = 0.12 s$. Do you have a (qualitative) explanation for the large difference?

PROBLEM 3

Density operators and entanglement

Give a brief and concise discussion of the following points:

a) List the general properties of density operators (or density matrices) and specify the difference between a *pure* and a *mixed* state.

b) For a composite system consisting of two parts \mathcal{A} and \mathcal{B} use the density operator formulation to explain the difference between, uncorrelated states, states with classical correlations (separable states) and entangled states.

c) Assume the full system is in a pure state, described by the state vector $|\psi\rangle$. What is meant by the *Schmidt decomposition* of this state vector relative to the two subsystems \mathcal{A} and \mathcal{B} ? Use the decomposition to find expressions for the reduced density operators of the two subsystems, and show that the von Neumann entropy of the reduced density operators are equal.

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Det matematisk-naturvitenskapelige fakultet

Eksamens i: FYS 4110 Ikke-relativistisk kvantemekanikk

Eksamensdag: Torsdag 3. desember, 2009

Tid for eksamen: kl. 14.30 (3 timer)

Oppgavesettet er på 3 sider

Tilatte hjelpebidrifter: Kalkulator

Øgrim og Lian eller Angell og Lian: Størrelser og enheter i fysikken
Rottmann: Matematisk formelsamling

Kontroller at oppgavesettet er komplett før du begynner å besvare spørsmålene.

BOKMÅL

OPPGAVE 1

To-nivåsystemer

Vi studerer i det følgende et kvantemekanisk system som er sammensatt av to to-nivåsystemer \mathcal{A} og \mathcal{B} . Hilbertrommet til det fulle systemet $\mathcal{H} = \mathcal{H}_A \otimes \mathcal{H}_B$ er dermed firedimensjonalt. De to delsystemene er dynamisk koblet, og Hamiltonoperatoren for det fulle systemet har formen

$$\hat{H} = \frac{1}{2}\hbar\omega(\sigma_z \otimes \mathbb{1} + \mathbb{1} \otimes \sigma_z) - i\hbar\lambda(\sigma_+ \otimes \sigma_- - \sigma_- \otimes \sigma_+) \quad (1)$$

hvor σ_z og σ_{\pm} er Pauli-matriser, med $\sigma_{\pm} = \frac{1}{2}(\sigma_x \pm i\sigma_y)$, $\hbar\omega$ er (den like store) energi-splittingen i hvert av de to delsystemene, og λ er en koblingsparameter. I tensorproduktuttrykkene regner vi at første faktor virker på delsystem \mathcal{A} og andre faktor på delsystem \mathcal{B} .

a) Vis at den tidsavhengige Schrödinger-ligningen har en løsning på formen

$$|\psi(t)\rangle = \cos(\lambda t)|+-\rangle + \sin(\lambda t)|-+\rangle \quad (2)$$

hvor $|+-\rangle = |+\rangle \otimes |- \rangle$ og $|-\rangle = |- \rangle \otimes |+\rangle$ og hvor $\sigma_z|\pm\rangle = \pm|\pm\rangle$ for hvert av delsystemene. Hva blir uttrykket for den tilsvarende tetthetsoperator $\hat{\rho}(t)$ når den skrives på bra-ket form?

b) Den tidsavhengige tetthetsoperatoren kan også uttrykkes ved Pauli-matriser på en tilsvarende måte som i (1). Finn dette uttrykket, og finn også de reduserte tetthetsoperatorene $\hat{\rho}_A(t)$ og $\hat{\rho}_B(t)$, begge uttrykt ved Paulimatriser (og identitets-operatoren).

c) Angi det generelle uttrykket for graden av sammenfiltrering i et sammensatt systemet når det befinner seg i en kvantemekanisk *ren tilstand*. I den tidsavhengige tilstanden beskrevet ovenfor, hva blir da uttrykket?

OPPGAVE 2

Koblede harmoniske oscillatorer

To harmoniske oscillatorer, kalt \mathcal{A} og \mathcal{B} , behandles som et sammensatt kvantemekanisk system. Hamiltonoperatoren til systemet har formen

$$\hat{H} = \hbar\omega(\hat{a}^\dagger\hat{a} + \hat{b}^\dagger\hat{b} + \mathbb{1}) + \hbar\lambda(\hat{a}^\dagger\hat{b} + \hat{b}^\dagger\hat{a}) \quad (3)$$

med $(\hat{a}, \hat{a}^\dagger)$ som senke- og heveoperatorer for \mathcal{A} og $(\hat{b}, \hat{b}^\dagger)$ som tilsvarende operatorer for \mathcal{B} . ω og λ er to reelle konstanter.

a) Vis at Hamiltonoperatoren kan skrives på diagonal form,

$$\hat{H} = \hbar\omega_c \hat{c}^\dagger \hat{c} + \hbar\omega_d \hat{d}^\dagger \hat{d} + \hbar\omega \mathbb{1} \quad (4)$$

hvor c og d er lineære kombinasjoner av a og b ,

$$c = \mu a + \nu b, \quad d = -\nu a + \mu b \quad (5)$$

og hvor μ og ν er reelle konstanter som tilfredsstiller $\mu^2 + \nu^2 = 1$. (Tilsvarende uttrykk gjelder for de hermitisk konjugerte operatorene \hat{c}^\dagger og \hat{d}^\dagger .) Bestem de nye parametrene ω_c , ω_d , μ og ν , uttrykt ved ω og λ . Sjekk at de nye operatorene \hat{c} og \hat{d} tilfredsstiller de samme kommutasjonsrelasjonene som \hat{a} og \hat{b} ved at $[\hat{c}, \hat{c}^\dagger] = [\hat{d}, \hat{d}^\dagger] = \mathbb{1}$ og $[\hat{c}, \hat{d}^\dagger] = 0$.

b) Anta at tilstanden $|\psi(0)\rangle$ til det sammensatte systemet ved $t = 0$ er en koherent tilstand for begge de nye variablene, slik at

$$\hat{c}|\psi(0)\rangle = z_{c0}|\psi(0)\rangle, \quad \hat{d}|\psi(0)\rangle = z_{d0}|\psi(0)\rangle \quad (6)$$

Tilstanden vil også på et senere tidspunkt være en koherent tilstand for \hat{c} og \hat{d} , med egenverdier

$$z_c(t) = e^{-i\omega_c t} z_{c0}, \quad z_d(t) = e^{-i\omega_d t} z_{d0} \quad (7)$$

Vis dette for $z_c(t)$. (Uttrykket for $z_d(t)$ følger på samme måte, og trengs derfor ikke vises.)

c) Vis at tilstanden $|\psi(t)\rangle$ også er en koherent tilstand for de opprinnelige harmonisk oscillatoroperatorene \hat{a} og \hat{b} , og bestem egenverdiene $z_a(t)$ og $z_b(t)$ uttrykt ved initialverdiene z_{a0} og z_{b0} .

OPPGAVE 3

Harmonisk oscillator i varmebad

En harmonisk oscillator med vinkelfrekvens ω er i termisk likevekt med et varmebad med temperatur T . Den befinner seg da i en blandet kvantemekanisk tilstand uttrykt ved tetthetsoperatoren

$$\hat{\rho} = N e^{-\beta \hat{H}} \quad (8)$$

med \hat{H} som Hamiltonoperatoren til den harmoniske oscillatoren, N som en normeringskonstant og $\beta = 1/kT$ hvor k er Boltzmanns konstant.

a) Sjekk at $\hat{\rho}$ tilfredsstiller kravene til en tetthetsmatrise og bestem normaliseringskonstanten N .

b) Vis at forventningsverdien for energien kan skrives som

$$E = \text{Tr}(\hat{H}\hat{\rho}) = \frac{1}{N} \frac{dN}{d\beta} \quad (9)$$

og finn energien som funksjon av β . Vis at for lav temperatur, $T \rightarrow 0$ eller $\beta \rightarrow \infty$, vil energien nærme seg grunntilstandsenergien til oscillatoren.

c) Tetthetsoperatoren kan skrives på diagonal form som

$$\hat{\rho} = \sum_{n=0}^{\infty} p_n |n\rangle \langle n| \quad (10)$$

hvor $p_n = Ne^{-\beta\hbar\omega(n+1/2)}$. Det vil si at vi kan se på tilstanden $\hat{\rho}$ som en statistisk blanding av energiegentilstander $|n\rangle$, vektet med sannsynlighetene p_n . Den samme tilstanden kan imidlertid også ses på som en statistisk blanding av koherente tilstander, på formen

$$\hat{\rho} = \int \frac{d^2 z}{\pi} p(|z|) |z\rangle \langle z| \quad (11)$$

hvor $p(|z|)$ er en sannsynlighetsfunksjon som bare avhenger av absoluttverdien $|z| = r$.

Vis at uttrykket (11) kan omformuleres til (10), og finn p_n uttrykt ved $p(r)$.

Vi minner om følgende uttrykk:

$$\langle n|z\rangle = \frac{z^n}{\sqrt{n!}} e^{-\frac{1}{2}|z|^2} \quad (12)$$

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Det matematisk-naturvitenskapelige fakultet

Eksamensdato: FYS 4110 / 9110 Ikke-relativistisk kvantemekanikk

Eksamensdag: Onsdag 8. desember, 2010

Tid for eksamen: kl. 14.30 (4 timer)

Oppgavesettet er på 4 sider

Tilatte hjelpebidrifter: Kalkulator

Øgrim og Lian eller Angell og Lian: Størrelser og enheter i fysikken
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OPPGAVE 1

Sammenfiltrering i et trepartikkelsystem

Tre partikler med halvtallig spinn, som vi referer til som \mathcal{A} , \mathcal{B} og \mathcal{C} , befinner seg i en sammensatt spinntilstand

$$|\psi\rangle = \frac{1}{\sqrt{2}}(|uuu\rangle - |ddd\rangle) \quad (1)$$

hvor $|uuu\rangle = |u\rangle_{\mathcal{A}} \otimes |u\rangle_{\mathcal{B}} \otimes |u\rangle_{\mathcal{C}}$, er tensorproduktet av *spinn-opp* (u) langs z-aksen for alle tre partiklene, mens $|ddd\rangle = |d\rangle_{\mathcal{A}} \otimes |d\rangle_{\mathcal{B}} \otimes |d\rangle_{\mathcal{C}}$ er tensorproduktilstanden svarende til *spinn-ned* (d) for alle tre partiklene. Vi antar at posisjonskoordinatene er helt frakoblet spinnkoordinatene og at spinnet derfor kan studeres separat.

a) Forklar hva vi mener med at de tre partiklene er i en *korrelert* spinntilstand, og hva vi mener med at spinnene er *sammenfiltret*.

Vi studerer i det følgende spinnsystemet som todelt, svarende til en oppsplittning $\mathcal{ABC} = \mathcal{A} + \mathcal{BC}$, slik at spinn \mathcal{A} definerer det ene undersystemet og de to andre spinnene, \mathcal{B} og \mathcal{C} , definerer det andre undersystemet.

b) Bestem de reduserte tetthetsoperatorene $\hat{\rho}_{\mathcal{A}}$ og $\hat{\rho}_{\mathcal{BC}}$ for de to delsystemene. Forklar hva vi mener med sammenfiltringsentropien til et todelt, sammensatt system og bestem verdien på denne for spinntilstanden (1). Hva menes med at spinntilstanden er en *maksimalt* sammenfiltret tilstand?

Tilstanden til delsystem \mathcal{BC} er beskrevet av tetthetsoperatoren $\hat{\rho}_{\mathcal{BC}}$. Hva sier denne om sammenfiltrering mellom de to spinnene \mathcal{B} og \mathcal{C} ?

Anta nå at de tre partiklene \mathcal{A} , \mathcal{B} og \mathcal{C} tas hånd om av tre fysikere (også identifisert som \mathcal{A} , \mathcal{B} og \mathcal{C}) som befinner seg på forskjellige steder, men som er i stand til å beskytte tilstanden til hver sin partikkel slik at den totale spinntilstanden (1) ikke forandrer seg før en av dem foretar en spinnmåling.

c) Ved et gitt tidspunkt måler \mathcal{A} spinnkomponenten langs x -aksen for sin partikkel og finner spinn-opp som måleresultatet. (Spinn-opp-tilstanden langs x -aksen blir betegnet $|f\rangle$ og spinn-ned-tilstanden $|b\rangle$.) Hun sender beskjed om dette til \mathcal{B} og \mathcal{C} , og disse beregner, med utgangspunkt i denne opplysningen, den nye tetthetsoperatoren $\hat{\rho}$ for det fulle systemet og bestemmer den nye reduserte tetthetsoperator $\hat{\rho}_{BC}$.

Hva er de nye uttrykkene for tetthetsoperatorene $\hat{\rho}$ og $\hat{\rho}_{BC}$? Er det noen endring i sammenfiltringen mellom spinnene \mathcal{B} og \mathcal{C} ?

OPPGAVE 2

Spinnflipp-stråling

Vi studerer i denne oppgaven overgang mellom spinntilstander for et elektron i et ytre magnetfelt som er rettet langs z -aksen, $\mathbf{B} = Be_z$. (Merk: vi benytter her \mathbf{e}_x , \mathbf{e}_y og \mathbf{e}_z som enhetsvektorer langs x -, y - og z -aksen, siden \mathbf{k} benyttes som bølgevektoren for det utsendte fotonet.) Hamiltonoperatoren skrives som $\hat{H} = \hat{H}_0 + \hat{H}_1$, hvor \hat{H}_0 svarer til den magnetiske dipolenergien i det ytre magnetfeltet, mens \hat{H}_1 beskriver koblingen mellom spinnet og strålingsfeltet. Vi har

$$\hat{H}_0 = \frac{1}{2}\omega_B\sigma_z, \quad \omega_B = -\frac{eB}{m} \quad (2)$$

med e som elektronladningen og m som elektronmassen. Frekvensen ω_B regnes som positiv.

Matriseelementet til spinnvekselvirkningen \hat{H}_1 ved emisjon av ett foton er i dipoltilnærmelsen

$$\langle B, 1_{ka} | \hat{H}_1 | A, 0 \rangle = i \frac{e\hbar}{2m} \sqrt{\frac{\hbar}{2\omega V \epsilon_0}} (\mathbf{k} \times \boldsymbol{\epsilon}_{ka}) \cdot \boldsymbol{\sigma}_{BA} \quad (3)$$

hvor $|A\rangle$ er den eksitere spinntilstanden (spinn-opp) og $|B\rangle$ er grunntilstanden (spinn-ned). Videre er $\boldsymbol{\epsilon}_{ka}$ en polarisasjonsvektor og $\omega = ck$ er sirkelfrekvensen til det emitterte fotonet, V er et normeringsvolum for den elektromagnetiske strålingen og $\boldsymbol{\sigma}_{AB}$ er matriseelementet til Paulimatrissen $\boldsymbol{\sigma} = \sigma_x \mathbf{e}_x + \sigma_y \mathbf{e}_y + \sigma_z \mathbf{e}_z$ mellom de to spinntilstandene. Vi minner om formen på Paulimatrissene,

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \quad (4)$$

a) Til første orden i perturbasjonsteori vil vinkelavhengigheten til det kvadrerte matriselementet $|\langle B, 1_{ka} | \hat{H}_1 | A, 0 \rangle|^2$ bestemme sannsynlighetsfordelingen for retningen til fotonet, $p(\phi, \theta)$, hvor (ϕ, θ) er polarvinklene til bølgevektoren \mathbf{k} . Bestem $p(\phi, \theta)$ fra uttrykket ovenfor. Vi minner om at ved summasjon over polarisasjonsretningene har vi $\sum_a |\boldsymbol{\epsilon}_{ka} \cdot \mathbf{b}|^2 = |\mathbf{b}|^2 - |\mathbf{b} \cdot \frac{\mathbf{k}}{k}|^2$ for en vilkårlig vektor \mathbf{b} . Normeringen av sannsynlighetsfordelingen er $\int d\phi \int d\theta \sin \theta p(\phi, \theta) = 1$.

b) Det kvadrerte matriselementet bestemmer også, for gitt \mathbf{k} , sannsynlighetsfordelingen over polarisasjonsretningen til fotonet. Anta at en fotondetektor registrerer fotoner utsendt med retning langs x -aksen ($\mathbf{k} = k\mathbf{e}_x$) og med polarisasjon langs polarisasjonsvektoren $\boldsymbol{\epsilon}(\alpha) = \cos \alpha \mathbf{e}_y + \sin \alpha \mathbf{e}_x$. Hva er sannsynlighetsfordelingen $p(\alpha)$ for å detektere det emitterte fotonet, som funksjon av vinkelen α ? (Anta også her at fordelingen er normert til 1, dvs. den beskriver sannsynlighet for forskjellige polarisasjonstilstander, forutsatt at fotonet er emitert langs x -aksen.)

c) Til en god tilnærmelse vil besetningssannsynligheten for den eksitere spinntilstanden reduseres eksponensielt med tiden

$$P_A(t) = e^{-t/\tau_A} \quad (5)$$

hvor levetiden τ_A til første orden i vekselvirkningen er bestemt av (den tidsuavhengige) overgangsraten

$$w_{BA} = \frac{V}{(2\pi\hbar)^2} \int d^3k \sum_a |\langle B, 1_{ka} | \hat{H}_1 | A, 0 \rangle|^2 \delta(\omega - \omega_B) \quad (6)$$

Benytt dette til å finne et uttrykk for levetiden τ_A .

OPPGAVE 3

En tvungen harmonisk oscillator

En kvantemekanisk, tvungen harmonisk oscillator er beskrevet ved en Hamiltonoperator på formen

$$\hat{H} = \hbar\omega_0(\hat{a}^\dagger\hat{a} + \frac{1}{2}) + \hbar\lambda(\hat{a}^\dagger e^{-i\omega t} + \hat{a}e^{i\omega t}) \quad (7)$$

hvor \hat{a} og \hat{a}^\dagger oppfyller standard komutasjonsrelasjoner for heve og senke-operatorer, og hvor ω_0 , ω og λ er tre konstanter. Vi innfører dimensjonsløse posisjons og bevegelsesemengde-operatorer som

$$\hat{x} = \frac{1}{2}(\hat{a} + \hat{a}^\dagger), \quad \hat{p} = -\frac{i}{2}(\hat{a} - \hat{a}^\dagger) \quad (8)$$

a) Vi minner om den generelle form på Heisenbergs bevegelsesligning,

$$\frac{d}{dt}\hat{A} = \frac{i}{\hbar}[H, \hat{A}] + \frac{\partial}{\partial t}\hat{A} \quad (9)$$

for en observabel \hat{A} . Benytt denne på senkeoperatoren \hat{a} , utled en bevegelsesligning for \hat{x} på formen

$$\frac{d^2\hat{x}}{dt^2} + \omega_0^2\hat{x} = C \cos \omega t \quad (10)$$

og bestem konstanten C .

b) Ved å anvende følgende tidsavhengige, unitære transformasjon

$$\hat{T}(t) = e^{i\omega t \hat{a}^\dagger \hat{a}} \quad (11)$$

vil den nye Hamiltonoperatoren, $\hat{H}_T(t)$, som bestemmer tidsutviklingen til de transformerte tilstandsvektorene $|\psi_T(t)\rangle = \hat{T}(t)|\psi(t)\rangle$, bli tidsuavhengig. Finn utrykket for denne operatoren.

c) En koherent tilstand er definert som en egentilstand for senkeoperatoren \hat{a} ,

$$\hat{a}|z\rangle = z|z\rangle \quad (12)$$

Anta at ved tiden $t = 0$ er oscillatoren i grunntilstanden for den λ -uavhengige del av Hamilton-operatoren, dvs.

$$|\psi(0)\rangle = |0\rangle, \quad \hat{a}|0\rangle = 0 \quad (13)$$

Vis at den fortsetter å være i en koherent tilstand under tidsutviklingen altså slik at

$$|\psi(t)\rangle = e^{i\alpha(t)}|z(t)\rangle \quad (14)$$

med $\alpha(t)$ er en tidsavhengig fase og $z(t)$ som en kompleks tidsavhengig funksjon.

Bestem funksjonen $z(t)$ og gi en kvalitativ beskrivelse av bevegelsen i det komplekse z -planet. Vis at realdelen $x(t) = (z(t) + z(t)^*)/2$ oppfyller samme bevegelsesligning (10) som posisjonsoperatoren $\hat{x}(t)$.

Vi minner om operatorrelasjonen

$$e^{\hat{A}}\hat{B}e^{-\hat{A}} = \hat{B} + [\hat{A}, \hat{B}] + \frac{1}{2!}[\hat{A}, [\hat{A}, \hat{B}]] + \dots \quad (15)$$

som gjelder for to vilkårlig valgte operatorer \hat{A} og \hat{B} .

UNIVERSITETET I OSLO

Det matematisk-naturvitenskapelige fakultet

Exam in: FYS 4110/ 9110 Non-relativistic quantum mechanics

Day of exam: Wednesday, December 7, 2011

Exam hours: 4 hours, beginning at 14:30

This examination paper consists of 2 problems on 3 pages

Permitted materials: Calculator

Øgrim og Lian or Angell og Lian: Størrelser og enheter i fysikken

Rottmann: Matematisk formelsamling

Language: The solutions may be written in Norwegian or English depending on your own preference.

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PROBLEM 1

Dressed photon states

A photon is interacting with an atom within a small reflecting cavity. The photon energy is close to the excitation energy of the atom, which connects the ground state to the first excited state. Due to interaction between the photon and the atom the stationary states of the composite system are admixtures of the photon and atomic states. These are sometimes referred to as a "dressed" photon states. In this problem we examine some of the properties of the dressed states.

The Hamiltonian of the photon-atom system can be written as

$$\hat{H} = \frac{1}{2}\hbar\omega_0\sigma_z + \hbar\omega\hat{a}^\dagger\hat{a} + i\hbar\lambda(\hat{a}^\dagger\sigma_- - \hat{a}\sigma_+) \quad (1)$$

where $\hbar\omega_0$ is then the energy difference between the two atomic levels, $\hbar\omega$ is the photon energy, and $\lambda\hbar$ is an interaction energy. The Pauli matrices act between the two atomic levels, with $\sigma_z|\pm\rangle = \pm|\pm\rangle$, and with $\sigma_\pm = (1/2)(\sigma_x \pm i\sigma_y)$ as matrices that raise or lower the atomic energy. \hat{a} and \hat{a}^\dagger are the photon creation and destruction operators.

a) We introduce the notation $|+, 0\rangle = |+\rangle \otimes |0\rangle$ and $|-, 1\rangle = |-\rangle \otimes |1\rangle$ for the relevant product states of the composite system, with 0, 1 referring to the photon number. Show that in the two-dimensional subspace spanned by these vectors the Hamiltonian takes the form

$$H = \frac{1}{2}\hbar\Delta \begin{pmatrix} \cos\phi & -i\sin\phi \\ +i\sin\phi & -\cos\phi \end{pmatrix} + \epsilon\mathbb{1} \quad (2)$$

where we assume $|-, 1\rangle$ to correspond to the lower matrix position and $|+, 0\rangle$ to the upper one. $\mathbb{1}$ denotes the 2×2 identity matrix. Express the parameters Δ , $\cos\phi$, $\sin\phi$, and ϵ in terms of ω_0 , ω and λ .

b) Find the energy eigenvalues E_\pm . Find also the eigenstates $|\psi_\pm(\phi)\rangle$, expressed in terms of the product states $|+, 0\rangle$ and $|-, 1\rangle$, and show that they are related by $|\psi_-(\phi)\rangle = |\psi_+(\phi + \pi)\rangle$.

In the following we focus on the state $|\psi_-(\phi)\rangle$, which we assume to be the one-photon state in the non-interacting case. This state (with a convenient choice of phase factor) can be written as $|\psi_-(\phi)\rangle = \cos \frac{\phi}{2} |-, 1\rangle + i \sin \frac{\phi}{2} |+, 0\rangle$.

c) Find expressions for the reduced density operators of the photon and of the atom for the state $|\psi_-(\phi)\rangle$. Discuss in what parameter interval the state is mostly a photon-like state and when it is mostly an atom-like state.

d) Determine the entanglement entropy as a function of ϕ , and find for what values the entropy is minimal and maximal. Relate this to the discussion in c).

e) At time $t = 0$ a single photon is sent into the cavity, where there is previously no photon and the atom is in its ground state. Determine the time dependent probability $p(t)$ for a photon later to be present in the cavity. Give a qualitative explanation of its oscillatory behaviour, and specify what determines the frequency and amplitude of the oscillations.

PROBLEM 2

A radiation problem

We consider a one-dimensional problem where a two-level system (A) interacts with a scalar radiation field (B). The notation we use is essentially the same as in Problem 1. The Hamiltonian of the system we consider is

$$\hat{H} = \frac{1}{2}\hbar\omega_A \sigma_z + \sum_k \hbar\omega_k \hat{a}_k^\dagger \hat{a}_k + \kappa \sum_k \sqrt{\frac{\hbar}{2L\omega_k}} (\hat{a}_k \sigma_+ + \hat{a}_k^\dagger \sigma_-) = \hat{H}_0 + \hat{H}_{int} \quad (3)$$

The first term is the two-level Hamiltonian, with energy splitting $\hbar\omega_A$, the second one is the free field contribution, with $k = 2\pi n/L$ (n - integer) as the wave number of the photon. L is a (large) normalization length. The third term is the interaction term \hat{H}_{int} , with κ as an interaction parameter. The frequency parameter is $\omega_k = ck$.

a) A general state of the two-level system is characterized by a vector \mathbf{r} , with $r \leq 1$, and with the corresponding density matrix as

$$\rho_A = \frac{1}{2}(\mathbb{1} + \mathbf{r} \cdot \boldsymbol{\sigma}) = \frac{1}{2} \begin{pmatrix} 1+z & x-iy \\ x+iy & 1-z \end{pmatrix} \quad (4)$$

Consider first that the interaction term \hat{H}_{int} is turned off, $\kappa = 0$, so that the time evolution operator of the two-level system is $\hat{\mathcal{U}}(t) = \exp(-\frac{i}{2}\omega_A t \sigma_z)$. Use this to determine the the density matrix $\rho_A(t)$ at time t , assuming that $\rho_A(0)$ is identical to the density matrix in (4), and show that the time evolution of \mathbf{r} is a precession around the z -axis with angular velocity ω_A .

b) Assume next that $\kappa \neq 0$ and that initially the two-level system is in the excited "spin up state", while the scalar field is in the vacuum state. Thus, the initial state is $|+, 0\rangle = |+\rangle \otimes |0\rangle$. It decays to the "spin down state" by emission of a field quantum. The final state we then write as $|-, 1_k\rangle = |-\rangle \otimes |1_k\rangle$.

The occupation probability of the excited state $|+\rangle$ decays exponentially, $P_+(t) = \exp(-\gamma t)$, with a decay rate γ that to first order in the interaction, and in the limit $L \rightarrow \infty$, is given by

$$\gamma = \frac{L}{(2\pi\hbar)^2} \int dk |\langle -, 1_k | \hat{H}_{int} | +, 0 \rangle|^2 \delta(\omega_k - \omega_A) \quad (5)$$

Determine the decay rate γ , expressed in terms of the parameters of the problem.

As discussed in the lectures an approximate way to handle the decay is to introduce an imaginary contribution to the energy of the decaying state. Assuming a more general initial state, of the form

$$|\psi(0)\rangle = (\alpha|+\rangle + \beta|-\rangle) \otimes |0\rangle = \alpha|+, 0\rangle + \beta|-, 0\rangle \quad (6)$$

with α and β as unspecified coefficients, with $|\alpha|^2 + |\beta|^2 = 1$, we make the corresponding *ansatz* for the time evolved state

$$|\psi(t)\rangle = (e^{-\frac{i}{2}\omega_A t - \gamma t/2} \alpha|+\rangle + e^{\frac{i}{2}\omega_A t} \beta|-\rangle) \otimes |0\rangle + \sum_k c_k(t)|-, 1_k\rangle \quad (7)$$

with $c_k(t)$ as decay parameters, which satify $c_k(0) = 0$.

c) Check what normalization of the state vector (7) means for the decay parameters, and determine the reduced density matrix matrix $\rho_A(t)$ of the two-level system.

d) Assume the same initial conditions as in b), $z(0) = 1$, $x(0) = y(0) = 0$ ($\alpha = 1$, $\beta = 0$). Determine the density matrix $\rho_A(t)$ and the corresponding time dependent vector $\mathbf{r}(t)$. Is the time evolution consistent with the expected exponential decay of the excited state of the two-level system? Give a brief description of the evolution of the entanglement between the two level system and the radiation field during the decay.

e) Choose another initial condition $x(0) = 1$, $y(0) = z(0) = 0$ ($\alpha = \beta = 1/\sqrt{2}$), and find also in this case the time evolution of the reduced density matrix and the components of the vector $\mathbf{r}(t)$. Sketch the time evolution of $\mathbf{r}(t)$ and compare qualitatively the motion with that in a) and d). Find $r(t)^2$ expressed as a function of γt , and sketch also this function. What does it show about the time evolution of the entanglement between the two subsystems A and B?

Assume in this paragraph $\gamma \ll \omega_A$.

UNIVERSITETET I OSLO

Det matematisk-naturvitenskapelige fakultet

Exam in: FYS 4110 Non-relativistic quantum mechanics

Day of exam: Friday, December 7, 2012

Exam hours: 4 hours, beginning at 14:30

This examination paper consists of 3 problems on 4 pages

Permitted materials: Calculator

Øgrim og Lian or Angell og Lian: Størrelser og enheter i fysikken

Rottmann: Matematisk formelsamling

Language: The solutions may be written in Norwegian or English depending on your own preference.

Make sure that your copy of this examination paper is complete before you begin.

PROBLEM 1

Two spin-half systems

A quantum system is composed of two interacting spin-half systems. The Hamiltonian has the form

$$\hat{H} = \frac{1}{2}\hbar\omega_1 \sigma_z \otimes \mathbb{1} + \frac{1}{2}\hbar\omega_2 \mathbb{1} \otimes \sigma_z + \frac{1}{2}\hbar\lambda(\sigma_+ \otimes \sigma_- + \sigma_- \otimes \sigma_+) \quad (1)$$

where σ_z og σ_{\pm} are Pauli matrices, with $\sigma_{\pm} = \frac{1}{2}(\sigma_x \pm i\sigma_y)$, $\hbar\omega_1$ and $\hbar\omega_2$ giving the splitting between the two energy levels of each of the spins, and with λ as a coupling parameter. The two factors of the tensor product refer to each of the two spin systems. We define the frequency difference as $\Delta = \omega_1 - \omega_2$ and introduce the following parametrization, $\Delta = \mu \cos \phi$ and $\lambda = \mu \sin \phi$. We further use $|\pm\rangle$ as notation for the eigenstates of σ_z . In the following we use the tensor products of these states as basis for the Hilbert space of the composite system.

a) Show that only the product states $|+-\rangle = |+\rangle \otimes |- \rangle$ and $|--\rangle = |+\rangle \otimes |- \rangle$ are mixed by the λ term in the Hamiltonian, and show that the mixing coefficients only depend on the angle ϕ , which we will assume to lie in the interval $0 \leq \phi \leq \pi/2$. Give the expression for the Hamiltonian as a 2x2 matrix, when restricted to the subspace spanned by $|+-\rangle$ and $|--\rangle$.

b) Find the corresponding two energy eigenvalues, and find the eigenstates expressed as functions of ϕ .

c) We now assume $\Delta = 0$. At time $t = 0$ the system is in the state $|+-\rangle$. Determine the time evolution of the state vector and the corresponding reduced density matrices for the two subsystems. Show that the entanglement entropy has a periodic behavior. What are the maximum and minimum values and what is the period of the oscillations.

PROBLEM 2

Atom-photon interaction in a cavity

An atom is trapped inside a small reflecting cavity. The energy difference between the ground state and the first excited state is $\Delta E = E_e - E_g \equiv \hbar\omega$, with ω matching the frequency of one of the electromagnetic cavity modes. This gives a strong coupling between the atomic states and this cavity mode, while the couplings to the other cavity modes are weak and can be neglected.

The composite system, atom plus cavity mode, is described by the following effective Hamiltonian

$$\hat{H} = \frac{1}{2}\hbar\omega\sigma_z + \hbar\omega\hat{a}^\dagger\hat{a} + \frac{1}{2}\hbar\lambda(\hat{a}^\dagger\sigma_- + \hat{a}\sigma_+) - i\gamma\hbar\hat{a}^\dagger\hat{a} \quad (2)$$

where the Pauli matrices act between the two atomic levels, with σ_z being diagonal in the energy basis, and $\sigma_\pm = (1/2)(\sigma_x \pm i\sigma_y)$ being matrices that raise or lower the atomic energy. \hat{a}^\dagger and \hat{a} are the photon creation and destruction operators. λ is an interaction parameter and γ is a decay parameter. The decay is due to the process where the photon escapes through the cavity walls. Both λ and γ are real-valued parameters, and we assume $\gamma \ll \lambda$ and $\gamma \ll \omega$.

We characterize the relevant states of the composite system as $|g, 0\rangle$, $|g, 1\rangle$ and $|e, 0\rangle$, where g refers to the atomic ground state, e to the excited state, and 0 and 1 refers to the absence or presence of a photon in the cavity mode.

a) Show that in the two-dimensional subspace spanned by the vectors $|g, 1\rangle$ and $|e, 0\rangle$ the Hamiltonian takes the form

$$H = \frac{1}{2}\hbar(\omega - i\gamma)\mathbb{1} + \frac{1}{2}\hbar \begin{pmatrix} -i\gamma & \lambda \\ \lambda & i\gamma \end{pmatrix} \quad (3)$$

where $|g, 1\rangle$ corresponds to the upper row of the matrix and $|e, 0\rangle$ to the lower one, and $\mathbb{1}$ is the identity matrix.

b) Assume that initially the system is in the state $|\psi(0)\rangle = |e, 0\rangle$. Show that the time evolution of the state vector can be written as

$$|\psi(t)\rangle = e^{-\frac{i}{2}\omega t - \frac{1}{2}\gamma t} ((\cos(\Omega t) + a \sin(\Omega t))|e, 0\rangle + ib \sin(\Omega t)|g, 1\rangle) \quad (4)$$

and determine the constants Ω , a and b .

c) Denote the corresponding density operator as $\hat{\rho}(t)$. The norm of this operator is not conserved, but if we add a contribution

$$\hat{\rho}_{tot}(t) = \hat{\rho}(t) + f(t)|g, 0\rangle\langle g, 0| \quad (5)$$

then the norm is conserved, with value 1, for a particular function $f(t)$. Determine this function, and comment on in what sense the addition of the last term in (5) is reasonable, when considering the physical process described by the Hamiltonian (3). Give a short qualitative description of the process described by (5).

PROBLEM 3

Distributed information

A secret message is distributed to a party of three, denoted A, B, and C, in the form of an entangled three-spin state, coded into three spin-half particles. As the receiving party knows in advance, the quantum state is one out of a selection of three,

$$|\psi_n\rangle = \frac{1}{\sqrt{3}}(|+--\rangle + \eta^n|--+\rangle + (\eta^*)^n|---\rangle), \quad \eta = e^{2\pi i/3} \quad (6)$$

where $n = 0, 1, 2$. The message is identified by the value of n , which means by which of the three quantum states that is distributed.

We use the notation $|+--\rangle = |+\rangle \otimes |-\rangle \otimes |-\rangle$ etc., where the single spin states $|\pm\rangle$ are orthogonal states in a basis referred to as *basis I*. The three spinning particles are distributed to A, B and C, one particle to each of them, with the first state in the tensor product corresponding to the spin sent to A, the second one to B and the third one to C. We assume the three-spin state is preserved under this distribution.

Each person in the receiving party can make (spin) measurements on the spinning particle he/she receives. The three can also communicate over a classical channel, which means that they can correlate their measurements and also compare the results of the measurements. They have, however no quantum channel available for communication. This means that all the observables that are available for measurements by the receiving party are of product form.

a) Determine the reduced density operator of A, and explain why, for any measurement he/she performs on his particle, no information can be extracted about which of the three spin states $|\psi_n\rangle$ is distributed. Also show that if A, B and C all make their spin measurements in *basis I*, even if they communicate their measured results, these cannot make any distinction between the three values of n .

Next, consider the situation where A and B are not able to communicate with C. They decide to perform measurements on the two spins they have received, and to make a probabilistic evaluation for the different values of n , based on the measured results. In order to do so they decide both to make their spin measurements in a rotated basis, which we refer to as *basis II*. The vectors in this basis are

$$|0\rangle = \frac{1}{\sqrt{2}}(|+\rangle + |-\rangle), \quad |1\rangle = \frac{1}{\sqrt{2}}(|+\rangle - |-\rangle) \quad (7)$$

The possible outcomes of the measurements they list with numbers $k = 1, 2, 3, 4$, with the correspondence

$$k = 1 : (0, 0), \quad k = 2 : (0, 1), \quad k = 3 : (1, 0), \quad k = 4 : (1, 1) \quad (8)$$

We refer to the corresponding states as $|\phi_k\rangle$, with $|\phi_1\rangle = |00\rangle = |0\rangle \otimes |0\rangle$, etc.

Before they do the measurements they evaluate for each three-spin state $|\psi_n\rangle$ the probabilities for the different measurement results (labeled by k). These probabilities are referred to as $p(k|n)$.

b) Find the reduced density operator $\hat{\rho}_n^{AB}$ and determine the probabilities $p(k|n)$ for different values of k and n . It is sufficient, due to repetitions of results, to consider $n = 0, 1$ and $k = 1, 2$.

Do you, in particular, see a reason why the probabilities are the same for $n = 1$ and $n = 2$, for all k ?

c) Assume now that A and B perform their measurements, with the result labeled by k . The probability for the state to be $|\psi_n\rangle$, under the condition that the measured result is k , we denote by $\bar{p}(n|k)$. Under the assumption that all spin states $|\psi_n\rangle$ are equally probable until the result of the measurement is known, statistics theory gives us the following relation

$$\bar{p}(n|k) = \frac{p(k|n)}{p(k)} \quad (9)$$

with $p(k)$ as a normalization factor. Determine $p(k)$ and the probability $\bar{p}(n|k)$ for each n in the case $k = 1 : (0, 0)$. What is most probably the message that has been distributed?

UNIVERSITETET I OSLO

Det matematisk-naturvitenskapelige fakultet

Eksamensdato: FYS 4110 / 9110 Ikke-relativistisk kvantemekanikk

Eksamensdag: Mandag 9. desember, 2013

Tid for eksamen: 4 timer, fra kl. 14:30

Oppgavesettet består av 3 oppgaver på 3 sider

Tillatte hjelpeemidler: Godkjent kalkulator

Angell og Lian: Størrelser og enheter i fysikken

Rottmann: Matematisk formelsamling

Kontroller at oppgavesettet er komplett før du begynner å besvare spørsmålene.

OPPGAVE 1

Tidsutvikling i et to-nivåsystem

Hamiltonoperatoren for et isolert to-nivåsystem (betegnet A) har formen $\hat{H}_0 = (1/2)\hbar\omega \sigma_z$, med σ_z som den diagonale Paulimatrisen. Vi betegner den normerte grunntilstandsvektoren som $|g\rangle$ og den eksitere tilstanden som $|e\rangle$. Systemet er i realiteten koblet til et strålingsfelt (betegnet S), og den eksitere tilstanden vil derfor henfalle til grunntilstanden under utsendelse av et strålingskvant. Vi lar $\hat{\rho}$ betegne den reduserte tetthetsoperatoren til delsystem A . Med god tilnærming kan tidsutviklingen av denne beskrives av den såkalte Lindbladligningen, her på formen

$$\frac{d\hat{\rho}}{dt} = -\frac{i}{\hbar} [H_0, \hat{\rho}] - \frac{1}{2}\gamma [\hat{\alpha}^\dagger \hat{\alpha} \hat{\rho} + \hat{\rho} \hat{\alpha}^\dagger \hat{\alpha} - 2\hat{\alpha} \hat{\rho} \hat{\alpha}^\dagger] \quad (1)$$

med γ som henfallsraten for overgangen $|e\rangle \rightarrow |g\rangle$, $\hat{\alpha} = |g\rangle\langle e|$ og $\hat{\alpha}^\dagger = |e\rangle\langle g|$.

På matriseform, i basis $\{|e\rangle, |g\rangle\}$, skriver vi tetthetsoperatoren $\hat{\rho}$ som

$$\hat{\rho} = \begin{pmatrix} p_e & b \\ b^* & p_g \end{pmatrix} \quad (2)$$

med p_e som sannsynligheten for å finne systemet i tilstand $|e\rangle$ og p_g som sannsynligheten for å finne det i tilstand $|g\rangle$.

a) Anta først at to-nivåsystemet ved tiden $t = 0$ er i tilstanden $\hat{\rho} = |e\rangle\langle e|$. Vis ved bruk av ligning (1) at sannsynligheten p_e avtar eksponensielt, med γ som henfallsrate, mens total sannsynlighet $p_e + p_g$ er bevart.

b) Anta så en annen initialtilstand hvor to-nivåsystemet ved $t = 0$ er i den rene tilstanden $|\psi\rangle = \frac{1}{\sqrt{2}}(|e\rangle + |g\rangle)$. Bestem den tidsavhengige tetthetsmatrisen $\hat{\rho}(t)$ med denne initialtilstanden.

c) Tetthetsoperatoren for system A kan alternativt uttrykkes ved Paulimatrissene, som $\hat{\rho} = \frac{1}{2}(\mathbb{1} + \mathbf{r} \cdot \boldsymbol{\sigma})$. Bestem funksjonen $r^2(t)$ i de to tilfellene ovenfor og vis at den i begge tilfeller har

minimum for $t = (1/\gamma) \ln 2$. Hva blir minimalverdien til r i de to tilfellene? Gi en kommentar om hva dette sier om sammenfiltringen mellom systemene A og S . (Vi forutsetter at det fulle systemet $A+S$ hele tiden er i en ren tilstand.)

OPPGAVE 2

Tre partikler i en sammenfiltret tilstand

Tre partikler med halvtallig spinn, som vi referer til som \mathcal{A} , \mathcal{B} and \mathcal{C} , befinner seg i en sammensatt spinntilstand

$$|\psi\rangle = \frac{1}{\sqrt{2}}(|uuu\rangle + |ddd\rangle) \quad (3)$$

hvor $|uuu\rangle = |u\rangle_{\mathcal{A}} \otimes |u\rangle_{\mathcal{B}} \otimes |u\rangle_{\mathcal{C}}$, er tensorproduktet av *spinn-opp* (u) langs z -aksen for alle tre partiklene, mens $|ddd\rangle = |d\rangle_{\mathcal{A}} \otimes |d\rangle_{\mathcal{B}} \otimes |d\rangle_{\mathcal{C}}$ er tensorprodukttilstanden svarende til *spinn-ned* (d) for alle tre partiklene. Vi antar at posisjonskoordinatene er helt frakoblet spinnkoordinatene og at spinnet derfor kan studeres separat. Det er ingen vekselvirkning mellom partiklene, og tilstanden (3) er derfor uendret så lenge det ikke måles på noen av spinnene.

Vi studerer i det følgende spinnsystemet som todelt, svarende til en oppsplitting $\mathcal{ABC} = \mathcal{A} + \mathcal{BC}$, slik at spinn \mathcal{A} definerer det ene undersystemet og de to andre spinnene, \mathcal{B} og \mathcal{C} , definerer det andre undersystemet.

a) Bestem de reduserte tetthetsoperatorene $\hat{\rho}_{\mathcal{A}}$ og $\hat{\rho}_{\mathcal{BC}}$ for de to delsystemene, og sammenfiltringsentropien til det sammensatte systemet. Hva menes med at de to delsystemene i denne tilstanden er *maksimalt* sammenfiltret? Delsystemet \mathcal{BC} kan videre tenkes sammensatt av undersystemene \mathcal{B} og \mathcal{C} . Hva sier tetthetsoperatoren $\hat{\rho}_{\mathcal{BC}}$ om sammenfiltrering mellom disse to.

b) Ved et gitt tidspunkt blir en spinnmåling utført på partikkel \mathcal{A} som bestemmer spinnkomponenten langs x -aksen som *spinn-opp* langs denne aksen. Denne informasjonen medfører at tetthetsoperatoren til systemet \mathcal{BC} blir endret. Hva blir den nye reduserte tetthetsoperator $\hat{\rho}'_{\mathcal{BC}}$? Har måling på spinnet til partikkel \mathcal{A} forandret sammenfiltringen mellom \mathcal{B} og \mathcal{C} ?

Vi minner om følgende: Med $|f\rangle$ som spinn-opp langs x -aksen og $|b\rangle$ som spinn-ned langs samme akse har vi relasjonene

$$|u\rangle = \frac{1}{\sqrt{2}}(|f\rangle - |b\rangle), \quad |d\rangle = \frac{1}{\sqrt{2}}(|f\rangle + |b\rangle) \quad (4)$$

c) Anta at tre-spinnsystemet igjen befinner seg i tilstanden (3). Denne gangen måles spinnet til \mathcal{A} langs en akse i xz -planet, som er rotert med vinkelen θ i forhold til z -aksen. Anta også at i dette tilfellet er måleresultatet *spinn-opp*. Finn hvordan måleresultatet nå påvirker tetthetsoperatoren for systemet \mathcal{BC} , og bestem sammenfiltringsentropien for sammensetningen $\mathcal{B} + \mathcal{C}$, som funksjon av vinkelen θ .

For de roterte spinntilstandene gjelder

$$\begin{aligned} |\theta, +\rangle &= \cos(\theta/2)|u\rangle + \sin(\theta/2)|d\rangle && (\text{spinn opp}) \\ |\theta, -\rangle &= -\sin(\theta/2)|u\rangle + \cos(\theta/2)|d\rangle && (\text{spinn ned}) \end{aligned} \quad (5)$$

der $\theta = 0$ svarer til kvantisert spinn langs z -aksen og $\theta = \pi/2$ til kvantisert spinn langs x -aksen.

OPPGAVE 3

Spinnflipp-stråling

Vi studerer i denne oppgaven overgang mellom to spinntilstander for et elektron i et ytre magnetfelt som er rettet langs z -aksen, $\mathbf{B} = B\mathbf{e}_z$. (Merk: vi benytter her \mathbf{e}_x , \mathbf{e}_y og \mathbf{e}_z som enhetsvektorer langs x -, y - og z -aksen.) Hamiltonoperatoren skrives som $\hat{H} = \hat{H}_0 + \hat{H}_1$, hvor \hat{H}_0 svarer til den magnetiske dipolenergien i det ytre magnetfeltet, mens \hat{H}_1 beskriver koblingen mellom spinnen og strålingsfeltet. Vi har

$$\hat{H}_0 = \frac{1}{2}\omega_B\sigma_z, \quad \omega_B = -\frac{eB}{m} \quad (6)$$

med e som elektronladningen og m som elektronmassen. Frekvensen ω_B regnes som positiv.

Matriseelementet til spinnvekselvirkningen \hat{H}_1 , ved emisjon av et foton, er i dipoltilnærmelsen

$$\langle B, 1_{\mathbf{k}a} | \hat{H}_1 | A, 0 \rangle = i \frac{e\hbar}{2m} \sqrt{\frac{\hbar}{2\omega V\epsilon_0}} (\mathbf{k} \times \boldsymbol{\epsilon}_{\mathbf{k}a}) \cdot \boldsymbol{\sigma}_{BA} \quad (7)$$

hvor $|A\rangle$ er den eksisterte spinntilstanden (spinn-opp) og $|B\rangle$ er grunntilstanden (spinn-ned). Videre er \mathbf{k} bølgetallsvektoren, $\boldsymbol{\epsilon}_{\mathbf{k}a}$ en polarisasjonsvektor og $\omega = c k$ er vinkelfrekvensen til det emitterte fotonet. V er et normeringsvolum for den elektromagnetiske strålingen og $\boldsymbol{\sigma}_{AB}$ er matriseelementet til Paulimatrisen $\boldsymbol{\sigma} = \sigma_x \mathbf{e}_x + \sigma_y \mathbf{e}_y + \sigma_z \mathbf{e}_z$ mellom de to spinntilstandene.

a) Til første orden i perturbasjonsteori vil vinkelavhengigheten til det kvadrerte matriselementet (summert over polarisasjonsindeksen) $\sum_a |\langle B, 1_{\mathbf{k}a} | \hat{H}_1 | A, 0 \rangle|^2$, bestemme sannsynlighetsfordelingen for retningen til fotonet, $p(\phi, \theta)$, hvor (ϕ, θ) er polarvinklene til bølgevektoren \mathbf{k} . Bestem $p(\phi, \theta)$ fra uttrykket ovenfor. Vi minner om at ved summasjon over polarisasjonsretningene har vi $\sum_a |\boldsymbol{\epsilon}_{\mathbf{k}a} \cdot \mathbf{b}|^2 = |\mathbf{b}|^2 - |\mathbf{b} \cdot \frac{\mathbf{k}}{k}|^2$ for en vilkårlig vektor \mathbf{b} . Normeringen av sannsynlighetsfordelingen er $\int d\phi \int d\theta \sin \theta p(\phi, \theta) = 1$.

b) Det kvadrerte matriselementet (uten sum over a) bestemmer også, for gitt \mathbf{k} , sannsynlighetsfordelingen over polarisasjonsretningen til fotonet. Anta at en fotondetektor registrerer fotoner utsendt langs x -aksen ($\mathbf{k} = k\mathbf{e}_x$), med polarisasjonsretning $\boldsymbol{\epsilon}(\alpha) = \cos \alpha \mathbf{e}_y + \sin \alpha \mathbf{e}_z$. Hva er sannsynligheten $p(\alpha)$ for å detektere det emitterte fotonet? Anta her at sannsynlighetsfordelingen er normert slik at summen over to ortogonale retninger er, $p(\alpha) + p(\alpha + \pi/2) = 1$. Hva sier resultatet om polarisasjonen til det emitterte fotonet?

Vi minner om standardformen på Paulimatrisesene,

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \quad (8)$$

UNIVERSITETET I OSLO

Det matematisk-naturvitenskapelige fakultet

Eksamensdato: FYS 4110 / 9110 Ikke-relativistisk kvantemekanikk

Eksamensdag: Mandag 8. desember, 2014

Tid for eksamen: 4 timer, fra kl. 14:30

Oppgavesettet består av x oppgaver på y sider

Tillatte hjelpeemidler: Godkjent kalkulator

Angell og Lian: Størrelser og enheter i fysikken

Rottmann: Matematisk formelsamling

Kontroller at oppgavesettet er komplett før du begynner å besvare spørsmålene.

1 Entanglement in a two-spin system

We consider a composite quantum system consisting of two spin-half systems, A and B . The relevant states are restricted to the two-dimensional subspace spanned by the two (orthogonal) Bell states

$$|1\rangle = \frac{1}{\sqrt{2}}(|+-\rangle + |-+\rangle), \quad |2\rangle = \frac{1}{\sqrt{2}}(|+-\rangle - |-+\rangle) \quad (1)$$

where we use the notation $|+-\rangle = |+\rangle \otimes |-\rangle$, with $|\pm\rangle$ referring to the two eigenstates of σ_z .

Consider first (Case I) a linear superposition of the two state vectors, of the form

$$|\psi(x)\rangle = \cos x |1\rangle + \sin x |2\rangle, \quad 0 \leq x \leq \frac{\pi}{2} \quad (2)$$

The corresponding density operator we denote by $\hat{\rho}_I(x) = |\psi(x)\rangle\langle\psi(x)|$.

a) Determine the reduced density operators $\hat{\rho}_{IA}(x)$ and $\hat{\rho}_{IB}(x)$ of the two spins and the corresponding entropies $S_{IA}(x)$ and $S_{IB}(x)$. Characterize the entanglement of the two spins for the special values $x = 0, \pi/4$, and $\pi/2$.

Consider next (Case II) the following linear combination of the density operators of the two Bell states,

$$\hat{\rho}_{II}(x) = \cos^2 x |1\rangle\langle 1| + \sin^2 x |2\rangle\langle 2|, \quad 0 \leq x \leq \frac{\pi}{2} \quad (3)$$

b) What is the von Neuman entropy of this state? Find the reduced density operators $\hat{\rho}_{IIA}(x)$ and $\hat{\rho}_{IIB}(x)$, and the corresponding entropies $S_{IIA}(x)$, and $S_{IIB}(x)$. Characterize also here the states of the full system for $x = 0, \pi/4$, and $\pi/2$.

For a composite quantum system in pure quantum state, the degree of entanglement is expressed by the von Neumann entropy of one of its subsystems. When the system is in a mixed

state we do not have a general, universally accepted, measure for the degree of entanglement. However, for a classical, statistical system we have the following inequality for the entropy of the full systems and its subsystem,

$$\Delta \equiv S - \max\{S_A, S_B\} \geq 0 \quad (4)$$

The breaking of this inequality in quantum system therefore indicates that the two subsystems are entangled.

c) Show that in the two cases *I* and *II* the functions $\Delta_I(x)$ and $\Delta_{II}(x)$ are negative for all x , except for one value of x .

OPPGAVE 2

Radiation damping

A charged particle is oscillating in a one-dimensional harmonic oscillator potential. It emits electric dipole radiation, with the rate for transition between an initial state i and a final state f given by the radiation formula

$$w_{fi} = \frac{4\alpha}{3c^2} \omega_{fi}^3 |x_{fi}|^2 \quad (5)$$

where α is the fine structure constant, $\hbar\omega_{fi}$ is the energy radiated in the transition, and c is the speed of light. x is the position coordinate of the particle, which is related to the raising and lowering operators of the harmonic oscillator by

$$x = \sqrt{\frac{\hbar}{2m\omega}} (\hat{a}^\dagger + \hat{a}) \quad (6)$$

with m as the mass of the particle.

a) Show that the non-vanishing transition rates are of the form

$$w_{n-1,n} = \gamma n \quad (7)$$

with $n = 0, 1, 2, \dots$ as referring to the energy levels of the harmonic oscillator, and γ as a constant decay parameter. Detemine γ .

The effect of the radiation on the state of the oscillating particle is described by the Lindblad equation in the following way

$$\frac{d\hat{\rho}}{dt} = -\frac{i}{\hbar} [H_0, \hat{\rho}] - \frac{1}{2}\gamma [\hat{a}^\dagger \hat{a} \hat{\rho} + \hat{\rho} \hat{a}^\dagger \hat{a} - 2\hat{a} \hat{\rho} \hat{a}^\dagger] \quad (8)$$

with $\hat{\rho}$ as the density operator of the particle and H_0 as the harmonic oscillator Hamiltonian, without decay.

b) In the following we focus on the diagonal terms of the density matrix, $p_n = \rho_{nn} = \langle n | \hat{\rho} | n \rangle$, which define the occupation probabilities of the energy eigenstates. Show that they satisfy the equation

$$\frac{dp_n}{dt} = -\gamma(np_n - (n+1)p_{n+1}) \quad (9)$$

Explain why this is consistent with the expression (7) for the transition rate $w_{n-1,n}$.

- c) Show that Eq. (9) implies that the expectation value of the excitation energy

$$E = \langle H_0 \rangle - \frac{1}{2}\hbar\omega \quad (10)$$

decays exponentially with time.

OPPGAVE 3

A state in thermal equilibrium

A quantum state in thermal equilibrium is described by the density operator

$$\hat{\rho}(\beta) = N(\beta)e^{-\beta\hat{H}} = N(\beta) \sum_n e^{-\beta E_n} |n\rangle\langle n| \quad (11)$$

with \hat{H} as the Hamiltonian, E_n as the corresponding energy eigenvalues, and $N(\beta)$ as a normalization factor. The parameter β is related to the temperature T by $\beta = 1/(k_B T)$, with k_B as Boltzmann's constant.

- a) Show that the expectation value for the energy can be expressed in terms of $N(\beta)$ as

$$E(\beta) = \frac{d}{d\beta} \ln N(\beta) \quad (12)$$

and find a similar expression for the von Neumann entropy $S(\beta) = \text{Tr}[\hat{\rho}(\beta) \ln \hat{\rho}(\beta)]$. (Use here the natural logarithm in the definition of S .)

b) For a two-level system, with Hamiltonian $\hat{H} = (\epsilon/2)\sigma_z$, determine the functions $N(\beta)$, $E(\beta)$ and $S(\beta)$, and make a sketch of the expectation value of the energy E as function of the temperature T .

c) Find the density operator expressed in the form $\hat{\rho} = (1/2)(\mathbb{1} + \mathbf{r} \cdot \boldsymbol{\sigma})$. Determine \mathbf{r} as a function of β and relate this to the results in b).

UNIVERSITY OF OSLO

Faculty of Mathematics and Natural Sciences

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Angell and Lian: Størrelser og enheter i fysikken

Rottmann: Matematisk formelsamling

Language: The solutions may be written in Norwegian or English depending on your own preference. Noen engelske ord er oversatt etter hver oppgave.

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PROBLEM 1

Two spin-half systems

A quantum system is composed of two interacting spin-half systems, referred to as system *A* and *B*. The Hamiltonian has the form

$$\hat{H} = \frac{1}{2}\hbar\omega(\sigma_z \otimes \mathbb{1} - \mathbb{1} \otimes \sigma_z) + \hbar\lambda(\sigma_+ \otimes \sigma_- + \sigma_- \otimes \sigma_+) \quad (1)$$

where σ_z og σ_{\pm} are Pauli matrices, with $\sigma_{\pm} = \frac{1}{2}(\sigma_x \pm i\sigma_y)$, $\hbar\omega$ giving the splitting between the two energy levels of each of the two spins, and with λ as a coupling parameter. The two factors of the tensor product refer to each of the two spin systems, with *A* corresponding to the first and *B* as the second factor. It is convenient to introduce new parameters by $\omega = a \cos \theta$ and $\lambda = a \sin \theta$, with $-\pi/2 < \theta \leq \pi/2$. We further use $|\pm\rangle$ as notation for the eigenstates of σ_z . In the following we use the tensor products of these states as basis for the Hilbert space of the composite system.

a) Show that only the product states $|+-\rangle = |+\rangle \otimes |- \rangle$ and $|-+\rangle = |+\rangle \otimes |- \rangle$ are mixed by the λ term in the Hamiltonian, and give the expression for the Hamiltonian as a 2x2 matrix, in the subspace spanned by $|+-\rangle$ and $|-+\rangle$.

b) Find the energy eigenvalues, and the energy eigenstates, expressed in terms of a and θ .

c) Determine, for the energy eigenstates, the density operator of the full system and the reduced density operators of the two subsystems, and determine the entanglement entropy of the eigenstates as functions of θ . What are the minimum and maximum values of the entropy functions? Make a comparison with the maximal possible value of the entanglement entropy of the two-spin system.

density operator = tetthetsoperator

entanglement = sammenfiltrering

PROBLEM 2

A driven harmonic oscillator

A quantum mechanical, driven harmonic oscillator is described by the following Hamiltonian

$$\hat{H} = \hbar\omega_0(\hat{a}^\dagger\hat{a} + \frac{1}{2}) + \hbar\lambda(\hat{a}^\dagger e^{-i\omega t} + \hat{a}e^{i\omega t}) \quad (2)$$

where \hat{a} og \hat{a}^\dagger satisfy the standard commutation relations for lowering and raising operators, and where ω_0 , ω og λ are three constants.

a) As a reminder, Heisenberg's equation of motion has the form

$$\frac{d}{dt}\hat{A} = \frac{i}{\hbar}[H, \hat{A}] + \frac{\partial}{\partial t}\hat{A} \quad (3)$$

for any given observable \hat{A} . Apply this to the operator \hat{a}_H , which is the operator \hat{a} transformed to the Heisenberg picture, and show that it satisfies an equation of the form

$$\frac{d^2\hat{a}_H}{dt^2} + \omega_0^2\hat{a}_H = Ce^{-i\omega t}\mathbb{1} \quad (4)$$

with C as a constant. Determine C .

b) Equation (4) can be solved as a linear differential equation, to give

$$\hat{a}_H(t) = \hat{a}e^{-i\omega_0 t} + D(e^{-i\omega t} - e^{-i\omega_0 t})\mathbb{1} \quad (5)$$

Show that (5) is a solution of (4) and determine the constant D .

c) A coherent state is defined as an eigenstate of the lowering operator \hat{a} ,

$$\hat{a}|z\rangle = z|z\rangle \quad (6)$$

Assume that the oscillator, at time $t = 0$, is in the ground state for the λ -independent part of the Hamiltonian, that is

$$|\psi(0)\rangle = |0\rangle, \quad \hat{a}|0\rangle = 0 \quad (7)$$

Show that, during the time evolution (in the Schrödinger picture), it will continue as a coherent state, so that

$$\hat{a}|\psi(t)\rangle = z(t)|\psi(t)\rangle \quad (8)$$

with $z(t)$ as a complex-valued function of time.

Find the function $z(t)$, and compare the time evolution of the real part $x(t) = (z(t) + z(t)^*)/2$ with the motion of the corresponding classical driven harmonic oscillator.

driven harmonic oscillator = tvungen harmonisk oscillator

coherent state = koherent tilstand

PROBLEM 3

Atom and photon in an optical microcavity

An atom is contained in an optical microcavity, with the energy difference between two of the atomic levels matching exactly the frequency of one of the electromagnetic cavity modes. A simplified description of the photon-atom system has the form of a two-level system coupled to a single electromagnetic mode. The Hamiltonian then takes the form

$$\hat{H} = \frac{1}{2}\hbar\omega\sigma_z + \hbar\omega\hat{a}^\dagger\hat{a} + \frac{1}{2}\hbar\lambda(\hat{a}^\dagger\sigma_- + \hat{a}\sigma_+) \quad (9)$$

where \hat{a}^\dagger and \hat{a} are photon creation and annihilation operators, and σ_z and σ_\pm are Pauli matrices with $\sigma_\pm = \frac{1}{2}(\sigma_x \pm i\sigma_y)$. These operators act between the two atomic levels with the upper and lower energy levels corresponding to the eigenvalues +1 and -1 respectively of σ_z . We refer in the following to $|\pm, n\rangle = |\pm\rangle \otimes |n\rangle$ as product states of the composite system, with $|\pm\rangle$ as the upper/lower atomic levels and $|n\rangle$ as the photon number states of the cavity mode.

a) Assume a single photon is introduced in the cavity at time $t = 0$ while the atom is in its ground state. Show that the atom-photon state will subsequently oscillate in the following way

$$|\psi(t)\rangle = e^{i\epsilon t}(\cos\Omega t|-, 1\rangle - i\sin(\Omega t)|+, 0\rangle) \quad (10)$$

and find Ω and ϵ expressed in terms of ω and λ .

To take into account leakage of photons from the cavity, we turn to a description of the time evolution in terms of the density operator. It is assumed to satisfy the Lindblad equation,

$$\frac{d\hat{\rho}}{dt} = -\frac{i}{\hbar}[H, \hat{\rho}] - \frac{1}{2}\gamma[\hat{a}^\dagger\hat{a}\hat{\rho} + \hat{\rho}\hat{a}^\dagger\hat{a} - 2\hat{a}\hat{\rho}\hat{a}^\dagger] \quad (11)$$

where γ is the escape rate for photons from the cavity.

b) The probability for finding the atom in the ground state with no photon in the cavity is $p_g = \langle -, 0|\hat{\rho}|-, 0\rangle$. Assume that there is initially a non-vanishing probability for a photon being present in the cavity. Show that that this will result in an increase in p_g with time, which is consistent with the expectation that the photon will escape from the cavity.

c) Assuming there is no contribution to $\hat{\rho}$ from higher excited states than $|-, 1\rangle$ and $|+, 0\rangle$, show that a closed set of coupled differential equations for the three variables $p_1 = \langle -, 1|\hat{\rho}|-, 1\rangle$, $p_0 = \langle +, 0|\hat{\rho}|+, 0\rangle$ and $b = \text{Im}\langle -, 1|\hat{\rho}|+, 0\rangle$ can be derived from the Lindblad equation.

Without solving the equations, give a qualitative description of what the expected time evolution will be with the same initial condition as in a).

cavity mode = kavitetsmode, hulromsmode

UNIVERSITY OF OSLO

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PROBLEM 1

Spin-half particle in a harmonic oscillator potential

A spin-half particle is moving in a one-dimensional harmonic oscillator potential (in the x -direction) under the influence of a constant magnetic field (in the z -direction). The Hamiltonian is

$$\hat{H} = \hbar\omega_0(\hat{a}^\dagger\hat{a} + \frac{1}{2}) + \frac{1}{2}\hbar\omega_1\sigma_z + \lambda\hbar(\hat{a}^\dagger\sigma_- + \hat{a}\sigma_+) \quad (1)$$

where the first term is the harmonic oscillator part, with ω_0 as the oscillator angular frequency, and the second term is the spin energy due to the magnetic field, with ω_1 as the angular spin precession frequency. The third term is a coupling term between the spin and the position coordinates of the particle, with λ as a coupling parameter. The spin flip operators are defined as $\sigma_\pm = \frac{1}{2}(\sigma_x \pm i\sigma_y)$, \hat{a} , \hat{a}^\dagger are the standard lowering and raising operators of the harmonic oscillator, and $\sigma_x, \sigma_y, \sigma_z$ are the Pauli spin matrices.

When $\lambda = 0$, the spin and position of the particle are uncoupled and the energy eigenstates are $|n, m\rangle$, with $n = 0, 1, 2, \dots$ as the harmonic oscillator quantum number and $m = \pm 1$ as the spin quantum number, corresponding to spin up/down along the z -axis. When $\lambda \neq 0$, the unperturbed eigenstates will pairwise be coupled by the Hamiltonian, so that $|n, +1\rangle$ is coupled to $|n + 1, -1\rangle$. (The state $|0, -\rangle$ is an exception; it is not affected by the coupling term and remains the non-degenerate ground state also for $\lambda \neq 0$.)

a) Consider the two-dimensional subspace spanned by the basis vectors $|0, +1\rangle$ and $|1, -1\rangle$. Show that in this space, and in the given basis, the Hamiltonian takes the form of a 2x2 matrix

$$H = \frac{1}{2}\hbar\Delta \begin{pmatrix} \cos\theta & \sin\theta \\ \sin\theta & -\cos\theta \end{pmatrix} + \hbar\epsilon\mathbb{1} \quad (2)$$

with $\mathbb{1}$ as the 2x2 identity matrix. Determine Δ , $\cos\theta$, $\sin\theta$ and ϵ .

b) Find the energies and eigenstates of H in the two-dimensional subspace, expressed as functions of Δ , θ and ϵ .

c) The basis vectors $|n, m\rangle$ can be regarded as tensor products of position and spin vectors, $|n, m\rangle = |n\rangle \otimes |m\rangle$. The two eigenstates found under b) will be entangled with respect to the position and spin variables. Determine the entanglement entropy as function of θ . What value for θ gives the least and what gives the greatest entanglement?

PROBLEM 2

Coupled harmonic oscillators

Two harmonic oscillators, referred to as \mathcal{A} and \mathcal{B} , form a composite quantum mechanical system. The Hamiltonian of the system has the form

$$\hat{H} = \hbar\omega(\hat{a}^\dagger\hat{a} + \hat{b}^\dagger\hat{b} + \mathbb{1}) + \hbar\lambda(\hat{a}^\dagger\hat{b} + \hat{b}^\dagger\hat{a}) \quad (3)$$

with $(\hat{a}, \hat{a}^\dagger)$ as lowering and raising operators for \mathcal{A} and $(\hat{b}, \hat{b}^\dagger)$ as corresponding operators for \mathcal{B} , while ω and λ are real valued constants.

a) Show that the Hamiltonoperator can be expressed in diagonal form as

$$\hat{H} = \hbar\omega_c\hat{c}^\dagger\hat{c} + \hbar\omega_d\hat{d}^\dagger\hat{d} + \hbar\omega\mathbb{1} \quad (4)$$

where c and d are linear combinations of a and b ,

$$\hat{c} = \mu\hat{a} + \nu\hat{b}, \quad \hat{d} = -\nu\hat{a} + \mu\hat{b} \quad (5)$$

with μ and ν as real constants satisfying $\mu^2 + \nu^2 = 1$, and determine the new parameters μ , ν , ω_c , and ω_d , expressed in terms of ω og λ . (The same type of expressions are found for the hermitian conjugate operators \hat{c}^\dagger og \hat{d}^\dagger .) Check that the new operators \hat{c} and \hat{d} satisfy the same set of harmonic oscillator commutation relations as \hat{a} and \hat{b} . It is sufficient to show

$$[\hat{c}, \hat{c}^\dagger] = [\hat{d}, \hat{d}^\dagger] = \mathbb{1}, \quad [\hat{c}, \hat{d}^\dagger] = 0 \quad (6)$$

b) Assume that the state $|\psi(0)\rangle$ of the composite system, at time $t = 0$, is a coherent state when expressed in terms of the new variables,

$$\hat{c}|\psi(0)\rangle = z_{c0}|\psi(0)\rangle, \quad \hat{d}|\psi(0)\rangle = z_{d0}|\psi(0)\rangle \quad (7)$$

Also at a later time the state $|\psi(t)\rangle$ will be a coherent state for both \hat{c} og \hat{d} , with eigenvalues

$$z_c(t) = e^{-i\omega_c t}z_{c0}, \quad z_d(t) = e^{-i\omega_d t}z_{d0} \quad (8)$$

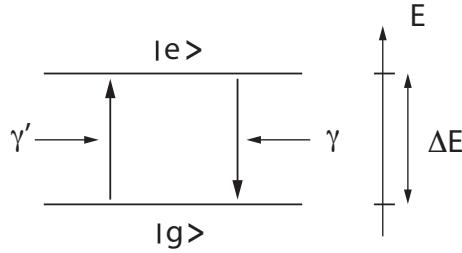
Show this for $z_c(t)$. (The expression for $z_d(t)$ follows in the same way, and is therefore not needed to be shown.)

c) Show that the state $|\psi(t)\rangle$ is a coherent state also for the original harmonic oscillator operators \hat{a} og \hat{b} , and find the eigenvalues $z_a(t)$ and $z_b(t)$ expressed in terms of z_{a0} and z_{b0} .

PROBLEM 3

Two-level system in a heat bath

We consider a two-level system, with $|g\rangle$ as the ground state and $|e\rangle$ as the excited state of the Hamiltonian \hat{H}_0 . The energy difference between the corresponding two energy levels is ΔE .



The system interacts weakly with a heat bath with temperature T . Energy can flow both ways, with γ as the rate for emission of energy to the heat bath in the transition $|e\rangle \rightarrow |g\rangle$ and γ' as the rate for absorption of energy in the transition $|g\rangle \rightarrow |e\rangle$. The situation is illustrated in the figure.

The temperature T of the heat bath and the energy gap ΔE determine the ratio between γ' and γ ,

$$\gamma' = \gamma e^{-\Delta E/kT} \quad (9)$$

where k is the Boltzman constant. Both transitions, corresponding to γ and γ' , contribute to the time evolution of the density operator of the two-level system. This is expressed by the Lindblad equation in the following way,

$$\begin{aligned} \frac{d}{dt}\hat{\rho} = & -\frac{i}{\hbar} [\hat{H}_0, \hat{\rho}] - \frac{1}{2}\gamma(|e\rangle\langle e|\hat{\rho} + \hat{\rho}|e\rangle\langle e| - 2|g\rangle\langle e|\hat{\rho}|e\rangle\langle g|) \\ & - \frac{1}{2}\gamma'(|g\rangle\langle g|\hat{\rho} + \hat{\rho}|g\rangle\langle g| - 2|e\rangle\langle g|\hat{\rho}|g\rangle\langle e|) \end{aligned} \quad (10)$$

The 2×2 matrix form of $\hat{\rho}$, in the basis $\{|g\rangle, |e\rangle\}$, we write as

$$\rho = \begin{pmatrix} p_e & b \\ b^* & p_g \end{pmatrix} \quad (11)$$

with p_e interpreted as the probability of occupation of the excited level and p_g as the probability of occupation of the ground state.

- a) Find from equation (10) expressions for the time derivatives of p_e , p_g and b , and check that they are consistent with preservation of total probability, $p_e + p_g$.
- b) The conditions for $\hat{\rho}$ to be a density operator give restrictions on the matrix elements in (11). What are these?
- c) Assume first that the two-level system and the heat bath are in thermal equilibrium, and the density matrix (11) therefore is time independent. Determine the values of variables p_e , p_g and b in this case.
- d) Consider next the situation with initial values $p_g = 1$, $p_e = 0$. Determine the time evolution of the occupation probabilities towards thermal equilibrium. What happens in the limits $T \rightarrow 0$ and $T \rightarrow \infty$?