# Lecture 12 <br> 01.10.2018 

Phase transformations

## A phase transformation

- Discontinuous change in the properties of substance when the environent is changed infinitesimaly.
- Change between phases - different forms of the substance
- Example: water vs. temperature...



## A phase transformation

- Often more than one variable that affects the phase of substance (eg, temperature \& pressure


WORLD'SHIGHEST MOUNTAIN


ELEVATION: -427m WATER BOILS: $101.4^{\circ} \mathrm{C}$,


## Phase diagram

- A graph showing the equilibrium phases as a function of pressure and temperature.
- Critical point - gas too dense -> fluid...


After: Chemistry, OpenStax

## Phase diagram, $\mathrm{H}_{2} \mathrm{O}$

- Let us consider water, $\mathrm{H}_{2} \mathrm{O}$
- Vapour pressure: coexisting gas and solid/liquid phase.
- Triple and critical points.
- Low pressures: no liquid phase
- We can still have metastable phases (superheated, supercooled)

1atm=101,3 kPa

Temperature ( ${ }^{\circ} \mathrm{C}$ )

## Phase diagram, $\mathrm{H}_{2} \mathrm{O}$

- $\mathrm{H}_{2} \mathrm{O}$ vs. other most other substances - different slope of liquid/solid.




## Peculiarity of Ice



Ice less dense than water...


Fig. 47.-Wire cutting through ice (from Tyndall).

Bragg, Concering the Nature of Things, 1925

## Phase diagram, $\mathrm{CO}_{2}$

- Higher pressure raises melting temperature.
- Dry ice sublimation is easily observed for $\mathrm{CO}_{2}$




## Phase diagram, Helium

- The most «exotic» phase behavior
- Liquid at zero T .




Zero viscosity, very high thermal conductivity
Fys2160, 2018
9

## Ferromagnet

- Curie temperature (iron 1043 K) magnetization disappears - no more phase boundary above critical point


## Superconductor

- Type I-tin, mercury, lead
- Low B-field and T gives zero electrical resistivity.



## Graphite vs. Diamonds

- Graphite more stable than diamond
- Gibbs free energy (G) is larger by 2900 J from graphite per mole at 1atm and room temperature.
- G - thermodynamic potential with thermal bath and fixed pressure.

$$
G=U-T S+P V
$$

## Graphite vs. Diamonds

- Pressure dependence of Gibbs free energy depends on volume

$$
\left(\frac{\partial G}{\partial P}\right)_{T, N}=V
$$

- This gives us slope:
$V=5.31 \times 10^{-6} \mathrm{~m}^{3}$ (graphite)
$V=3.42 \times 10^{-6} \mathrm{~m}^{3}$ (diamond)
$15 \mathrm{GPa}=>50 \mathrm{~km}$ into the Earth.



## Graphite vs. Diamonds

- Temperature dependence of Gibbs free energy depends on entropy

$$
\left(\frac{\partial G}{\partial T}\right)_{P, N}=-S
$$

- Larger temperature => Lower G.
- Decrease in G is larger for graphite which has more entropy.
- High T => more pressure needed to keep diamonds stable


Figure 1. Gibbs free energy vs. temperature for graphite ( $G_{g}$ ), liquid carbon ( $\mathrm{G}_{\text {liq }}$ ), and diamond $\left(\mathrm{G}_{\mathrm{d}}\right)$.

From : Mat. Res. Lett. 7, 353, 2018

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## Complex plasmas

Plasmas = ionised gas, electrically conductive
$99 \%$ of visible matter in the universe

Complex plasma - small particles injected into plasma get charged and interact with each other.


## Plasma crystal

To counteract the gravity we carry out experiments on parabolic flights or in on ISS

First experiment in physics on ISS!



## Complex plasmas

- Potential energy vs. Kinetic energy (coupling parameter)

$$
\Gamma=\frac{Z_{d}^{2} e^{2}}{4 \pi \epsilon_{0} \Delta k_{B} T_{d}}
$$

- Interparticle distance vs. sceening length

$$
\kappa=\frac{\Delta}{\lambda_{D}}
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## Phase diagram

- A summary



## Clausius-Clapeyron Relation

- Entropy determines the temperature dependence of the Gibbs free energy
- Volume determines the pressure dependence of the Gibbs free energy.
- The line on the phase transition diagram can be related to the entropies and volumes of the two phases!


## Clausius-Clapeyron Relation

- Let us consider a mole of material, and for example liquid to gas tranisition.
- At the phase boundary the Gibbs free energy must the the same:

$$
G_{l}=G_{g}
$$

(the same is valid for their chemical potentials i.e., $\mathrm{G} /$ molecule)

- Now let us change $T$ and $P$ by dT and dP in such a way that we stay on the line.

$$
d G_{l}=d G_{g}
$$



## Clausius-Clapeyron Relation

If we now take that

$$
d G=-S d T+V d P+\mu d N
$$

we can obtain (we do not change $N$ in the system)

$$
-S_{l} d T+V_{l} d P=-S_{g} d T+V_{g} d P
$$

and we get the slope of the line:

$$
\frac{d P}{d T}=\frac{S_{g}-S_{l}}{V_{g}-V_{l}}
$$

## Clausius-Clapeyron Relation

Often it is more convenient to use difference in entropies:

$$
S_{g}-S_{l}=L / T
$$

where $L$ is latent heat needed for converting the materials phase. An we get the Clausius-Clapeyron Relation:

$$
\frac{d P}{d T}=\frac{L}{T \Delta V}
$$

## The van der Waals Model

Ideal gas low modified for liquid-gas system to account for interactions:

$$
\left(P+\frac{a N^{2}}{V^{2}}\right)(V-N b)=N k T
$$

Volume modified by Nb - we can not compress a fluid to $\mathrm{V}=0$ ! $b$ - minimum volume occupied by a molecule $\times N$

Pressure modified by $a N^{2} / V^{2}$ reflects short range interactions between molecules.

## The van der Waals Model

Pressure modified by $a N^{2} / V^{2}$ reflects short range interactions between molecules.
Potential interactions of single molecule with all its neighbors is proportional to the density: $\sim N / V$
Total potential energy ( $N$ molecules): $E=-\frac{a N^{2}}{V}$
From thermodynamic identitty and fixed entropy («frozen» particles):

$$
d U=-P d V \rightarrow P=-(\partial U / \partial V)_{S}
$$

Thus the pressure due to interactions between molecules:

$$
P_{e}=-\frac{d}{d V}\left(-\frac{a N^{2}}{V}\right)=-\frac{a N^{2}}{V^{2}}
$$

## van der Waals equation $P=\frac{N k T}{V-N b}-\frac{a N^{2}}{V^{2}}$

- Is qualitative exact for describing fluids.
- Does not account for that gas gets denser with $P$ and clusterisation can occur.
- $a$ and $b$ depend on substances. Eg. water: $b^{\sim}(4 \AA)^{3}$, a $\sim 10 \mathrm{eV} \AA^{3}$
Eg. nitrogen: $b^{\sim}(4 \AA)^{3}, a \sim 2.5 \mathrm{eV} \AA^{3}$

Problem: Pressure decreases when we compress a fluid???

Isotherms of the hard-sphere gas

van der Waals equation $P=\frac{N k T}{V-N b}-\frac{a N^{2}}{V^{2}}$
Equilibrium state at given T and P is determined by the Gibbs energy

$$
d G=-S d T+V d P+\mu d N
$$

we find, for a fixed amount $N$ :

$$
\left(\frac{\partial G}{\partial V}\right)_{N, T}=V\left(\frac{\partial P}{\partial V}\right)_{N, T}
$$

we can now use the van der Waals equation and get:

$$
\left(\frac{\partial G}{\partial V}\right)_{N, T}=\frac{N k T V}{(V-N b)^{2}}+\frac{a N^{2}}{V^{2}}
$$

van der Waals equation $P=\frac{N k T}{V-N b}-\frac{a N^{2}}{V^{2}}$

$$
\left(\frac{\partial G}{\partial V}\right)_{N, T}=\frac{N k T V}{(V-N b)^{2}}+\frac{a N^{2}}{V^{2}}
$$

After taking $V=(V-N b)+N b$ and integrating the last equation we get:

$$
G=-N k T \ln (V-N b)+\frac{(N k T)(N b)}{(V-N b)}-\frac{2 a N^{2}}{V}+c(T)
$$

which allows us to plot G for any given T .

## van der Waals equation

Thermodynamically stable are only the lowest Gibbs energies

We should go straight from point 2 to 6

What is the pressure at the phase transformation?

We can find it from the figure, but also calculate it from PV diagram...
Let us integrate around the closed loop... 2-3-4-5-6


Figure 5.21. Gibbe free energy as a function of pressure for a van der Waals fluid at $T=0.9 T_{c}$. The corresponding isotherm is shown at right. States in the range 2-3-4-5-6 are unstable.

## van der Waals equation

$$
0=\int_{\text {loop }} d G=\int_{\text {loop }}\left(\frac{\partial G}{\partial P}\right)_{T} d P=\int_{\text {loop }} V d P
$$

This corresponds to the area below A - area below B. If $A=B$ we found the pressure!


## Phase diagrams...

- Vapor pressure at phas $\epsilon$ transition
- We have critical temperature $\mathrm{T}_{\mathrm{c}}$ and critical volume and pressure: V and P at $\mathrm{T}_{\mathrm{c}}$
- They all together define the critical point



Figure 5.23. Complete phase diagrams predicted by the van der Waals model.

