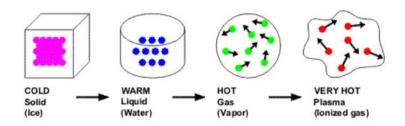
## Lecture 12

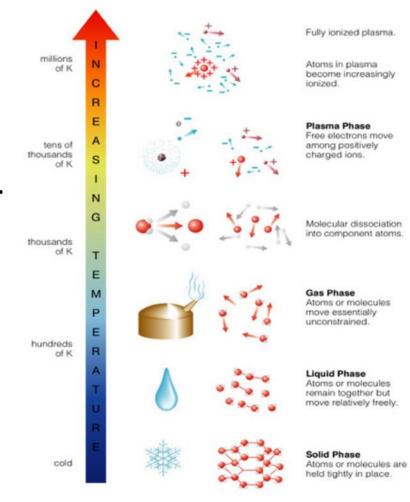
01.10.2018

Phase transformations

### A phase transformation

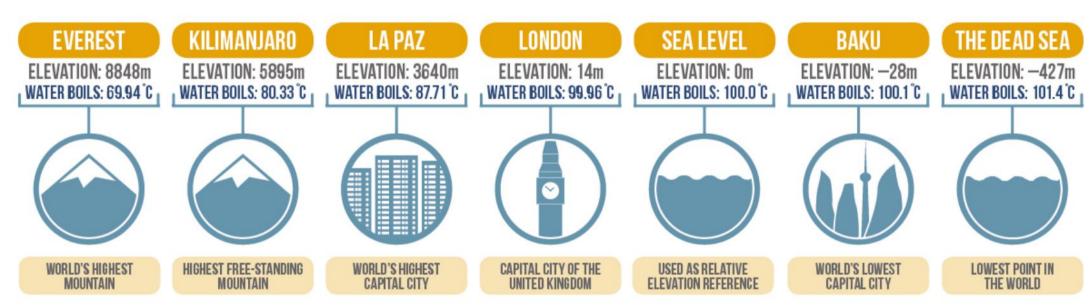
- **Discontinuous change** in the properties of substance when the environent is changed infinitesimaly.
- Change between phases different forms of the substance
- Example: water vs. temperature...





#### A phase transformation

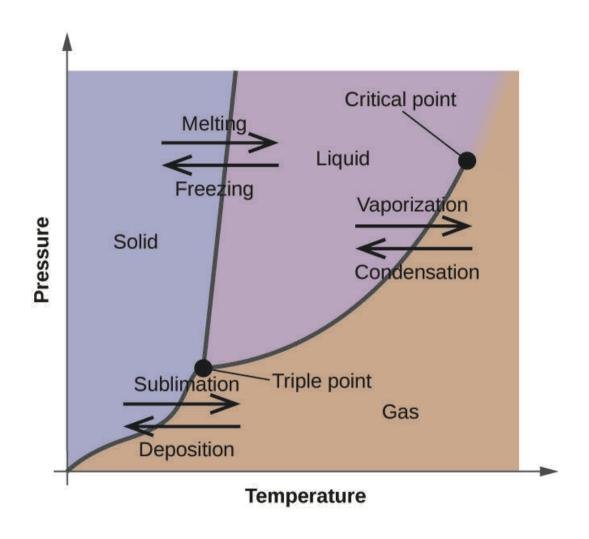
 Often more than one variable that affects the phase of substance (eg, temperature & pressure



After: www.compoundchem.com

#### Phase diagram

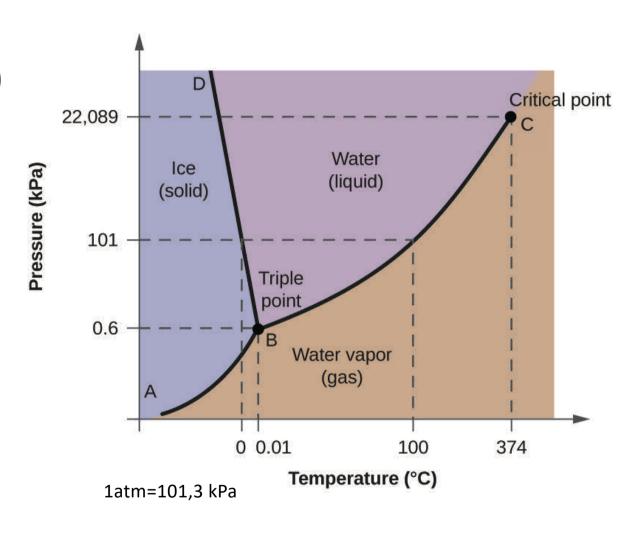
- A graph showing the equilibrium phases as a function of pressure and temperature.
- Critical point gas too dense -> fluid...



After: Chemistry, OpenStax

## Phase diagram, H<sub>2</sub>0

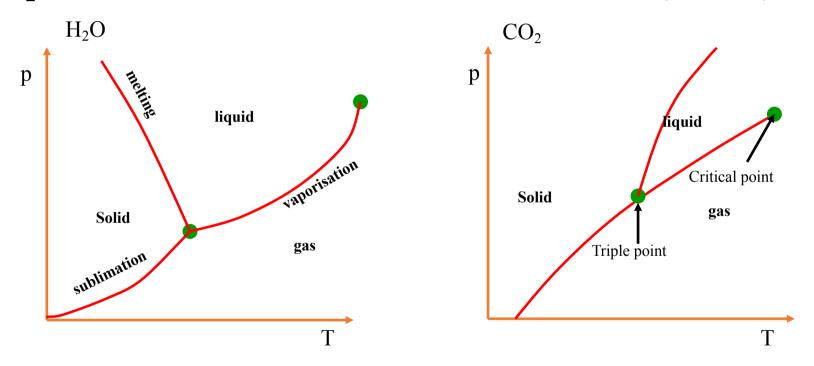
- Let us consider water, H<sub>2</sub>O
- Vapour pressure: coexisting gas and solid/liquid phase.
- Triple and critical points.
- Low pressures: no liquid phase
- We can still have metastable phases (superheated, supercooled)



After: www.compoundchem.com

## Phase diagram, H<sub>2</sub>0

• H<sub>2</sub>O vs. other most other substances – different slope of liquid/solid.



## Peculiarity of Ice



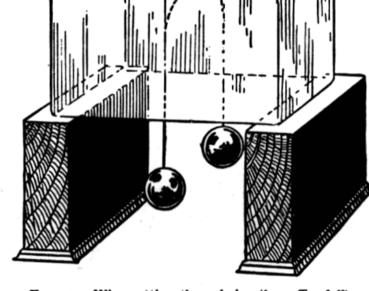


Fig. 47.—Wire cutting through ice (from Tyndall).

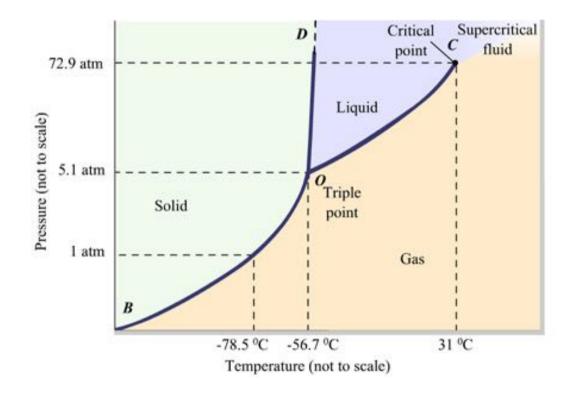
Ice less dense than water...

Bragg, Concering the Nature of Things, 1925

## Phase diagram, CO<sub>2</sub>

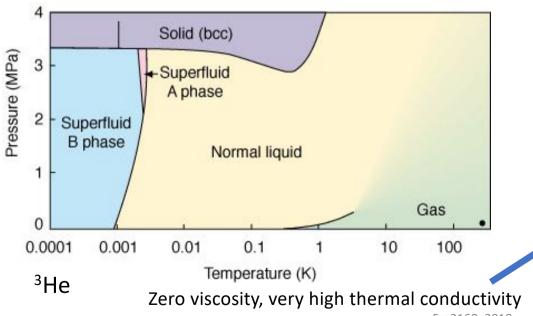
- Higher pressure raises melting temperature.
- Dry ice sublimation is easily observed for CO<sub>2</sub>

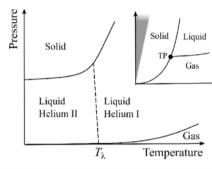


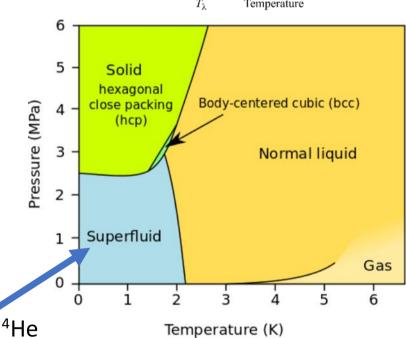


### Phase diagram, Helium

- The most «exotic» phase behavior
- Liquid at zero T.





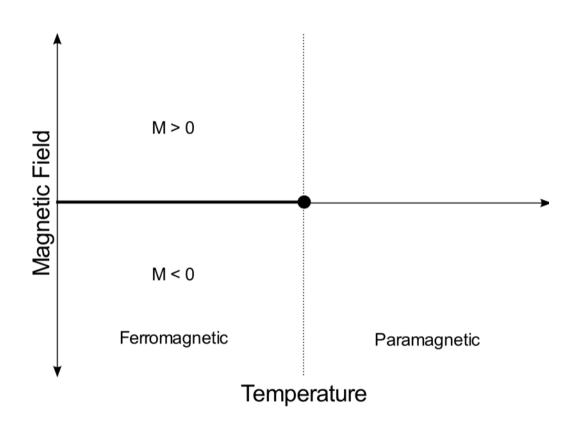


Fys2160, 2018

9

#### Ferromagnet

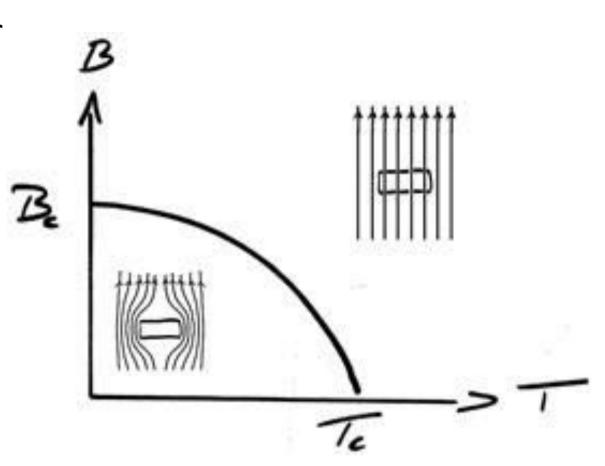
 Curie temperature (iron 1043 K) magnetization disappears – no more phase boundary above critical point



## Superconductor

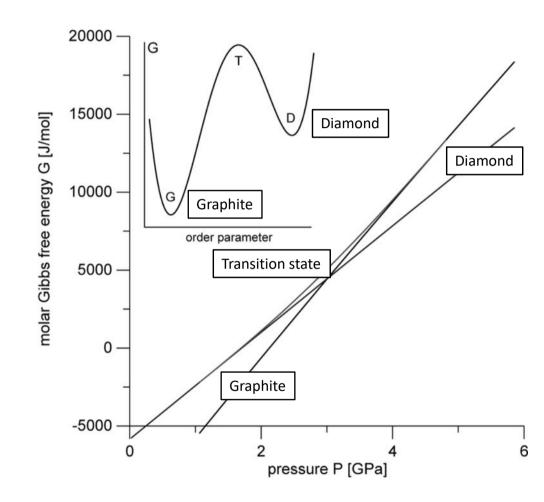
 Type I – tin, mercury, lead

• Low B-field and T gives zero electrical resistivity.



- Graphite more stable than diamond
- Gibbs free energy (G) is larger by 2900 J from graphite per mole at 1atm and room temperature.
- G thermodynamic potential with thermal bath and fixed pressure.

$$G = U - TS + PV$$



 Pressure dependence of Gibbs free energy depends on volume

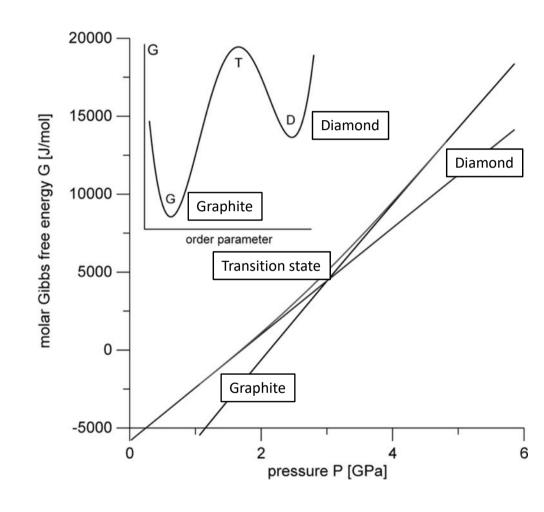
$$\left(\frac{\partial G}{\partial P}\right)_{T,N} = V$$

• This gives us slope:

 $V=5.31 \times 10^{-6} \text{ m}^3 \text{ (graphite)}$ 

 $V=3.42 \times 10^{-6} \text{ m}^3 \text{ (diamond)}$ 

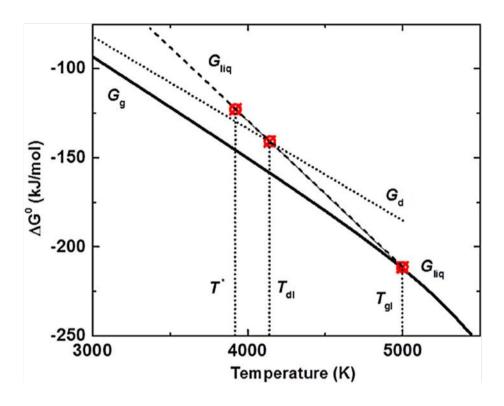
15 GPa => 50 km into the Earth.



 Temperature dependence of Gibbs free energy depends on entropy

$$\left(\frac{\partial G}{\partial T}\right)_{P,N} = -S$$

- Larger temperature => Lower G.
- Decrease in G is larger for graphite which has more entropy.
- High T => more pressure needed to keep diamonds stable



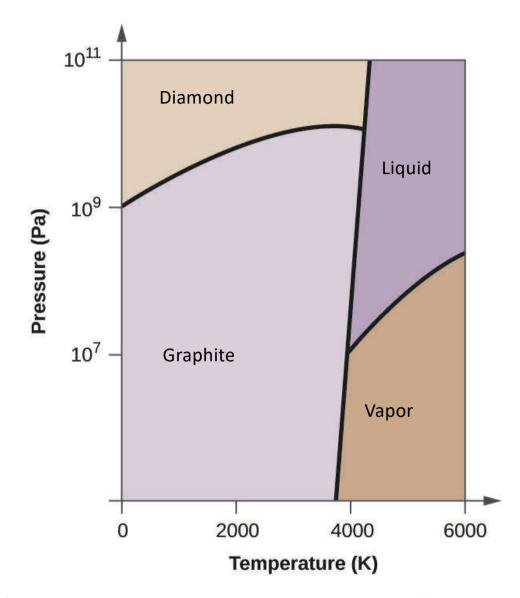
**Figure 1.** Gibbs free energy vs. temperature for graphite  $(G_g)$ , liquid carbon  $(G_{lig})$ , and diamond  $(G_d)$ .

From: *Mat. Res. Lett.* 7, 353, 2018

 Temperature dependence of Gibbs free energy depends on entropy

$$\left(\frac{\partial G}{\partial T}\right)_{P,N} = -S$$

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### Complex plasmas

Plasmas = ionised gas, electrically conductive

99% of visible matter in the universe

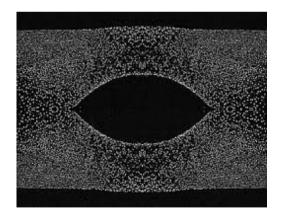
Complex plasma – small particles injected into plasma - get charged and interact with each other.

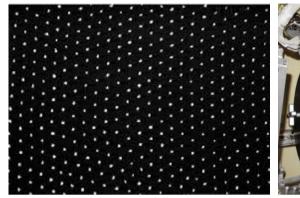


## Plasma crystal

To counteract the gravity we carry out experiments on parabolic flights or in on ISS

First experiment in physics on ISS!









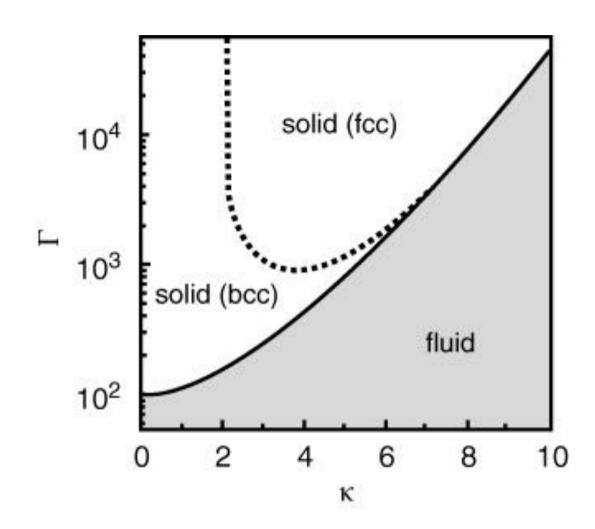
#### Complex plasmas

 Potential energy vs. Kinetic energy (coupling parameter)

$$\Gamma = \frac{Z_d^2 e^2}{4\pi\epsilon_0 \Delta k_B T_d}$$

 Interparticle distance vs. sceening length

$$\kappa = \frac{\Delta}{\lambda_D}$$



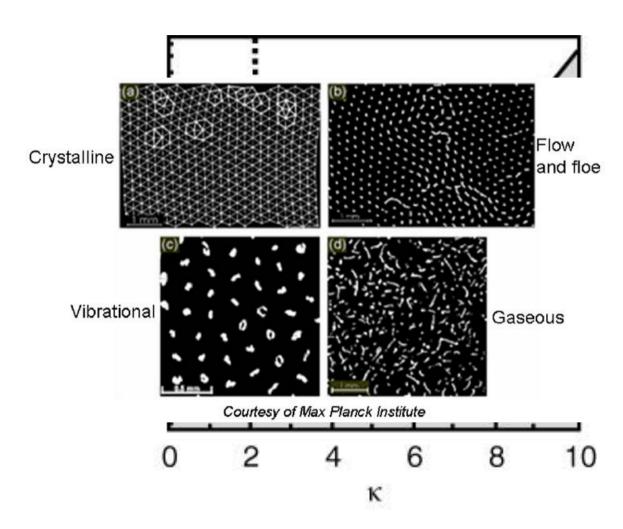
### Complex plasmas

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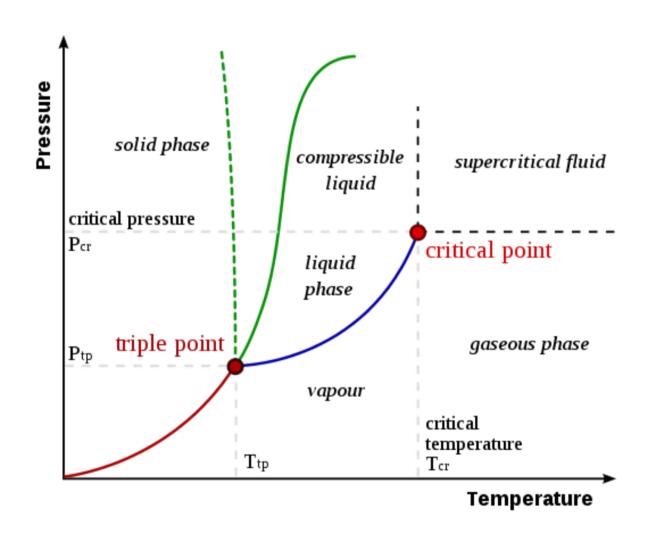
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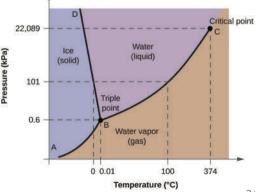
## Phase diagram

• A summary



- Entropy determines the temperature dependence of the Gibbs free energy
- Volume determines the pressure dependence of the Gibbs free energy.

 The line on the phase transition diagram can be related to the entropies and volumes of the two phases!



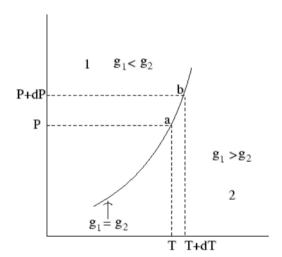
- Let us consider a mole of material, and for example liquid to gas transsition.
- At the phase boundary the Gibbs free energy must the the same:

$$G_l = G_q$$

(the same is valid for their chemical potentials i.e., G/molecule)

 Now let us change T and P by dT and dP in such a way that we stay on the line.

$$dG_l = dG_q$$



If we now take that

$$dG = -SdT + VdP + \mu dN$$

we can obtain (we do not change N in the system)

$$-S_l dT + V_l dP = -S_g dT + V_g dP$$

and we get the slope of the line:

$$\frac{dP}{dT} = \frac{S_g - S_l}{V_q - V_l}$$

Often it is more convenient to use difference in entropies:

$$S_g - S_l = L/T$$

where *L* is latent heat needed for converting the materials phase. An we get the **Clausius-Clapeyron Relation**:

$$\frac{dP}{dT} = \frac{L}{T\Delta V}$$

#### The van der Waals Model

Ideal gas low modified for liquid-gas system to account for interactions:

$$\left(P + \frac{aN^2}{V^2}\right)(V - Nb) = NkT$$

Volume modified by Nb – we can not compress a fluid to V=0! b – minimum volume occupied by a molecule x N

Pressure modified by  $aN^2/V^2$  reflects short range interactions between molecules.

#### The van der Waals Model

Pressure modified by  $aN^2/V^2$  reflects short range interactions between molecules.

Potential interactions of single molecule with all its neighbors is proportional to the density:  $\sim N/V$ 

Total potential energy (N molecules): 
$$E=-rac{aN^2}{V}$$

From thermodynamic identity and fixed entropy («frozen» particles):

$$dU = -PdV \rightarrow P = -(\partial U/\partial V)_S$$

Thus the pressure due to interactions between molecules:

$$P_e = -\frac{d}{dV} \left( -\frac{aN^2}{V} \right) = -\frac{aN^2}{V^2}$$

# van der Waals equation $P = \frac{NkT}{V - Nb} - \frac{aN^2}{V^2}$

Is qualitative exact for describing fluids.

• Does not account for that gas gets denser with P and clusterisation

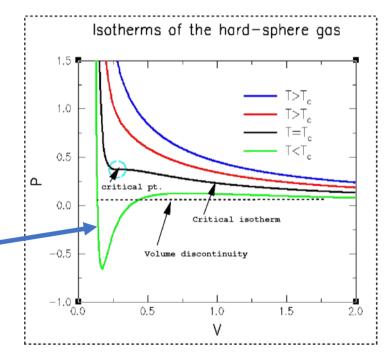
can occur.

• a and b depend on substances.

Eg. water:  $b^{\sim}(4 \text{ Å})^3$  a  $\sim 10 \text{ eV Å}^3$ 

Eg. nitrogen:  $b^{\sim}(4 \text{ Å})^3$ , a  $\sim 2.5 \text{ eV Å}^3$ 

Problem: Pressure decreases when we compress a fluid???



# van der Waals equation $P = \frac{NkT}{V - Nb} - \frac{aN^2}{V^2}$

Equilibrium state at given T and P is determined by the Gibbs energy

$$dG = -SdT + VdP + \mu dN$$

we find, for a fixed amount N:

$$\left(\frac{\partial G}{\partial V}\right)_{N,T} = V \left(\frac{\partial P}{\partial V}\right)_{N,T}$$

we can now use the van der Waals equation and get:

$$\left(\frac{\partial G}{\partial V}\right)_{N,T} = \frac{NkTV}{(V-Nb)^2} + \frac{aN^2}{V^2}$$

## van der Waals equation $P = \frac{NkT}{V - Nb} - \frac{aN^2}{V^2}$

$$\left(\frac{\partial G}{\partial V}\right)_{N,T} = \frac{NkTV}{(V - Nb)^2} + \frac{aN^2}{V^2}$$

After taking V=(V-Nb)+Nb and integrating the last equation we get:

$$G = -NkT \ln(V - Nb) + \frac{(NkT)(Nb)}{(V - Nb)} - \frac{2aN^2}{V} + c(T)$$

which allows us to plot G for any given T.

## van der Waals equation

Thermodynamically stable are only the lowest Gibbs energies

We should go straight from point 2 to 6

What is the pressure at the phase transformation?

We can find it from the figure, but also calculate it from PV diagram...

Let us integrate around the closed loop... 2-3-4-5-6

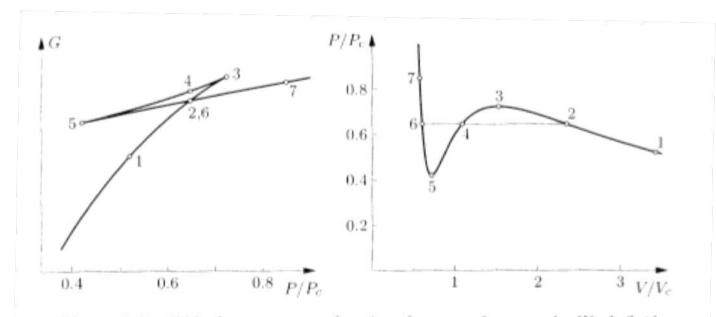
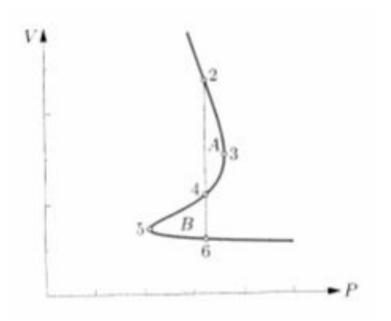


Figure 5.21. Gibbs free energy as a function of pressure for a van der Waals fluid at  $T = 0.9T_c$ . The corresponding isotherm is shown at right. States in the range 2-3-4-5-6 are unstable.

#### van der Waals equation

$$0 = \int_{\text{loop}} dG = \int_{\text{loop}} \left(\frac{\partial G}{\partial P}\right)_T dP = \int_{\text{loop}} V dP$$

This corresponds to the area below A – area below B. If A=B we found the pressure!



### Phase diagrams...

- Vapor pressure at phase transition
- We have critical temperature T<sub>c</sub> and critical volume and pressure: V and P at T<sub>c</sub>
- They all together define the critical point

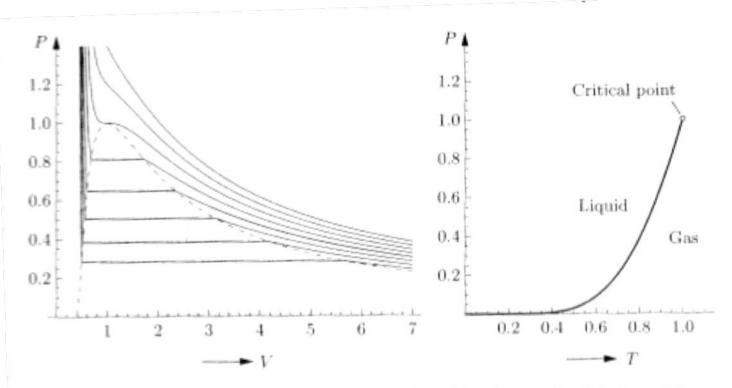


Figure 5.23. Complete phase diagrams predicted by the van der Waals model.