Lecture 3

Ideal gas model Equipartion of energy

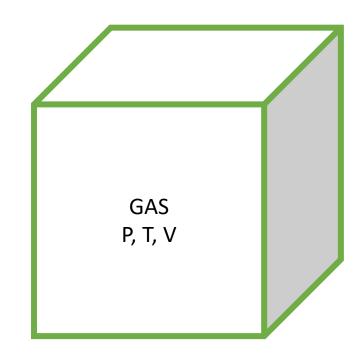
27.08.2018

Ideal gas in a box at equilibrium

Equation of state: PV = nRT

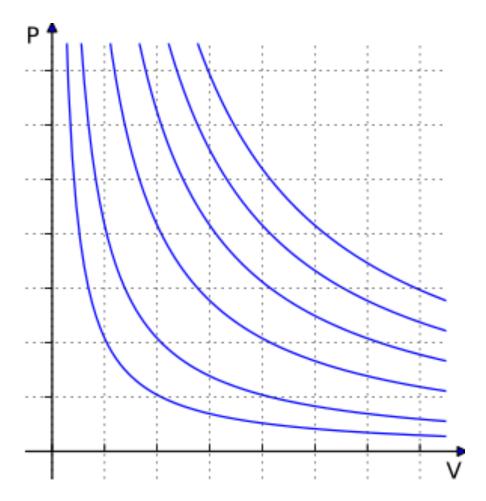
Experimental fact

- Pressure $[P] = 1 \ atm = 1.013 \times 10^5 \ Pa$
- Volume $[V] = 1 \ liter = 10^{-3} \ m^3$
- Temperature [T] = 1 K
- Gas constant $R = 8.31 \, J \, mol^{-1} K^{-1}$
- Number of moles $n = \frac{N}{N_A}$
- Avogadro's number $N_A = 6.02 \times 10^{23}$ of molecules in 1 mole



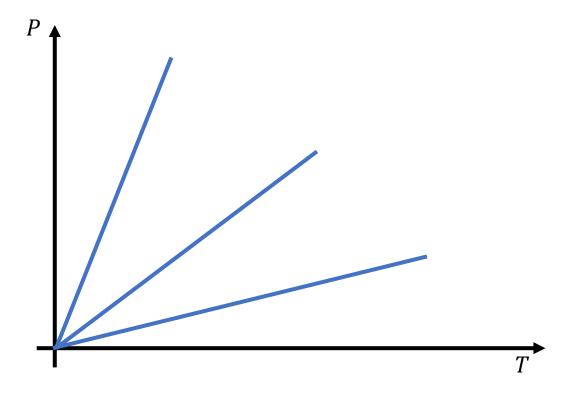
Isotherms

$$P = nRT \frac{1}{V}$$



P-T diagram: isocores

$$P = \frac{nR}{V} T$$



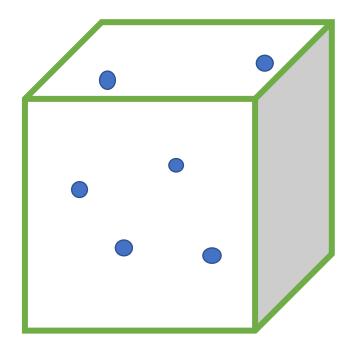
Ideal gas: Microscopic (statistical) model

Equation of state: PV = NkT

- Number of gas molecules N
- Boltzmann's constant

$$k = \frac{R}{N_A} = 1.381 \times 10^{-23} J K^{-1}$$

(conversion factor between energy and temperature?)

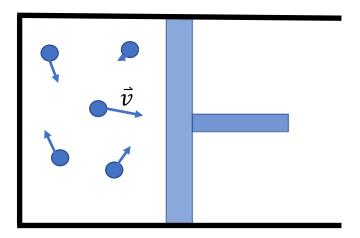


<u>Ideal gas model: Pressure</u>

Newtonian gas particle:

•
$$m\frac{d\vec{v}}{dt} = \vec{F}(=0)$$

Independent, identical particles

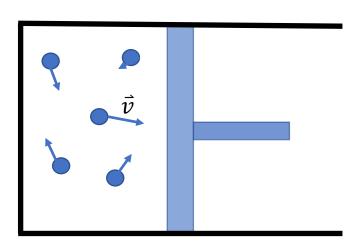


Q: What is the pressure in the ideal gas model?

Q: What is the pressure in the ideal gas model?



compute the pressure on the piston



What is the <u>long-time averaged</u> pressure exerted on the piston by one particle?

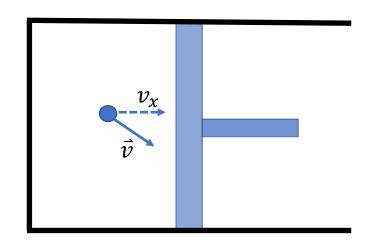
•
$$P = \frac{\overline{F}_{x,piston}}{A} = -\frac{\overline{F}_{x}}{A} = -\frac{m}{A} \overline{\left(\frac{\Delta v_{x}}{\Delta t}\right)}$$



$$\Delta t = \frac{2L}{v_x}$$

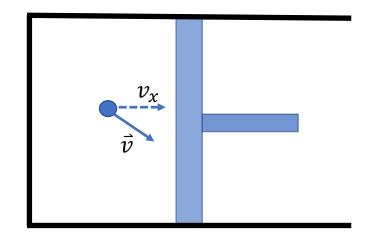
 Change in velocity after one perfect collision with the piston

$$\Delta v_x = v_{x,after} - v_{x,before} = -2v_x$$



What is the average pressure exerted on the piston by the gas particles?

•
$$P = -\frac{m}{A} \overline{\left(\frac{\Delta v_X}{\Delta t}\right)} = \frac{m\overline{v_X^2}}{AL} = \frac{m\overline{v_X^2}}{V}$$



using

$$\circ \Delta t = \frac{2L}{v_x}$$

$$0 \Delta v_x = v_{x,after} - v_{x,before} = -2v_x$$

What is the average pressure exerted by the gas particles?

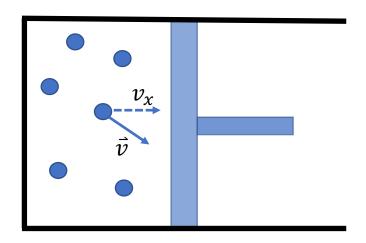
Pressure is the long-time average of v_x^2

$$P = \frac{m\overline{\mathbf{v}_{\mathbf{x}}^2}}{V}$$

Consider **ALL** N particles at a given time and average over their velocities (**ensemble average**)

$$PV = Nm\langle v_{x}^{2} \rangle = \frac{N}{3}m\langle v^{2} \rangle$$

$$PV = \frac{2N}{3} \langle K_{trans} \rangle, \qquad \langle K_{trans} \rangle = \frac{1}{2} m \langle v^2 \rangle$$

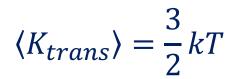


Ideal gas model: Temperature

$$PV = \frac{2N}{3} \langle K_{trans} \rangle$$

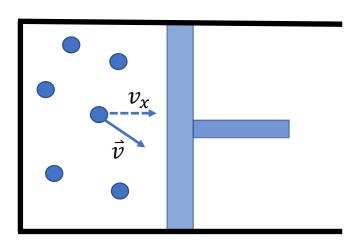
$$PV = NkT$$





In 1-D:
$$\langle K_{trans} \rangle = \frac{1}{2}kT$$

In d-D:
$$\langle K_{trans} \rangle = \frac{d}{2}kT$$



$$\langle |v| \rangle \approx \sqrt{\langle v^2 \rangle} = \sqrt{\frac{3kT}{m}}$$

Ideal gas model: Equipartition

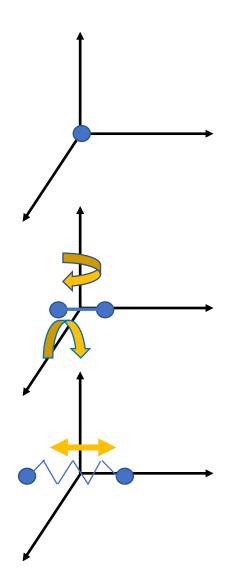
Equipartion of energy (theorem):

At equilibrium with temperature T, any quadratic form of the internal energy equal $\frac{1}{2}kT$ per degree of freedom

$$U = K + U_{potential} = \frac{f}{2}NkT$$

What is a degree of freedom?

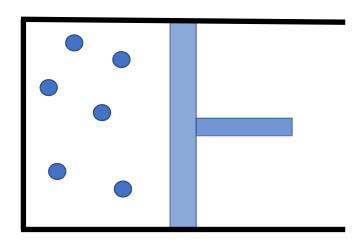
- Translation $K_{trans} = \frac{1}{2}mv^2$
- Rotation $K_{rot} = \frac{1}{2}I\omega^2$
- Vibration/oscillation $U_{harm} = -\frac{1}{2} k(\Delta x)^2$



Ideal gas model: Equipartition

Equipartion of energy (theorem):

At equilibrium with temperature T, any quadratic form of the internal energy is $\frac{1}{2}kT$ per degree of freedom

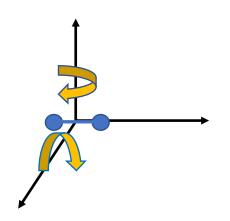


Example:

What is the internal energy of a diatomic gas:

$$U = K + U_{potential}$$

$$U = \frac{5}{2}NkT$$



Ideal gas model: Equipartition

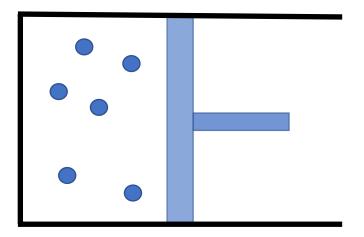
Equipartion of energy (theorem):

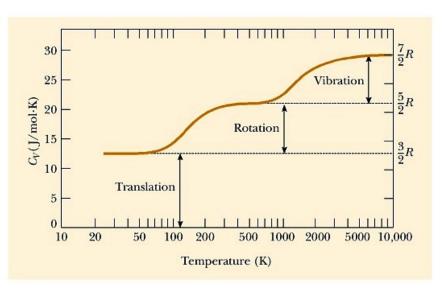
$$U = \frac{f}{2}NkT$$

What are the heat capacities C_V , C_P ?

$$C_V = \frac{dU}{dT} = \frac{f}{2}Nk$$

$$C_P = C_V + P\left(\frac{\partial V}{\partial T}\right)_P = \frac{f+2}{2}Nk$$





The molar specific heat of hydrogen as a function of temperature. The horizontal scale is logarithmic. Note that hydrogen liquefies at $20~\rm K$. Fys 2160~2018

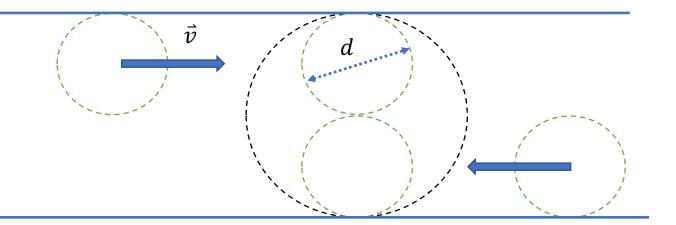
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Mean free-path

Free path between collisions

Number of molecules per unit volume (density):

$$n_V = \frac{N}{V}$$



Effective collision area: $\sigma = \pi d^2$

Effective collision volume:

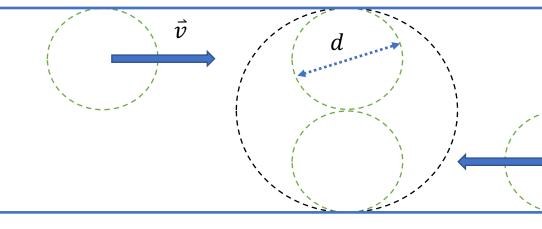
$$V_c = \overline{v}t\sigma = \pi d^2\overline{v}t$$

Mean free path:
$$\lambda = \frac{\overline{v}t}{n_V V_c}$$

$$\lambda = \frac{1}{\pi d^2 n_V}$$

Mean free path

Gas particles spend most of their time between collisions



Mean free path:
$$\lambda = \frac{1}{\pi d^2 n_V}$$

Average time between collisions:

$$au = rac{\lambda}{\overline{v}}, \qquad \overline{v} pprox rac{\sqrt{kT}}{\sqrt{m}}$$

$$\tau = \frac{\lambda}{\overline{v}} \approx \frac{1}{\pi d^2 n_V} \frac{\sqrt{m}}{\sqrt{kT}}$$

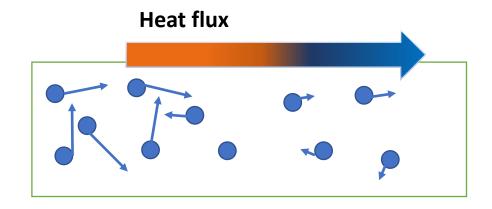
The mean-free path λ and mean lag time τ dictate the *kinetics* of the gas (*diffusion and heat conduction properties*)

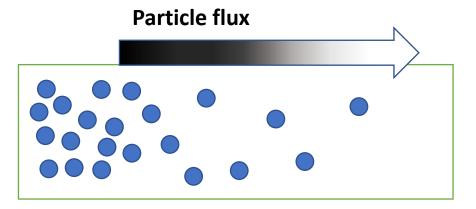
Relaxation to equilibrium by diffusion

Diffusion:

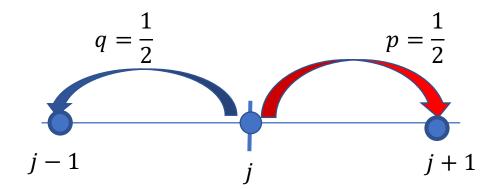
Net transport of *energy, momentum* or *particles* through random thermal motion and particle collisions until thermodynamic equilibrium is reached

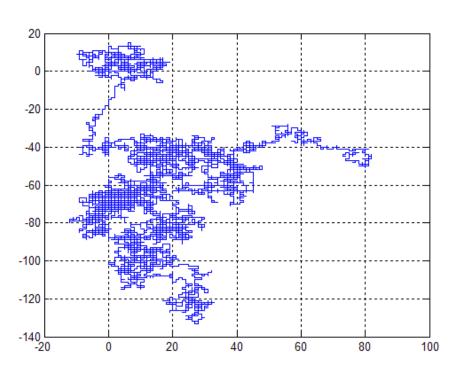
- At any T > 0K, particles are in *thermal motion*
- Collisions between particles -> particle trajectory is a zigzag -- random (diffusive particle)

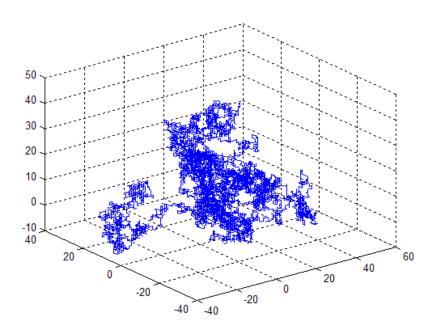




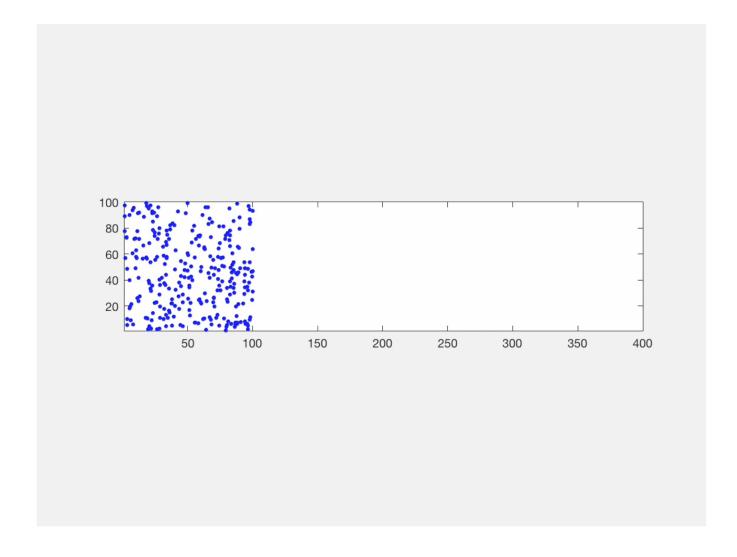
Particle diffusion: Random walk (RW)

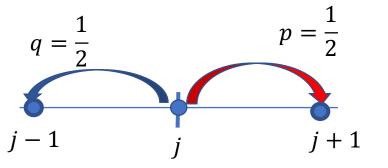






Random walk and diffusion





Fick's first law:

• Diffusive particle drift from high to low concentration $C = \frac{N}{V}$ of particles

$$J = -D\nabla C$$

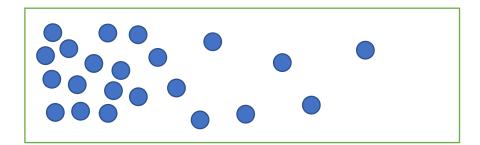
• **Diffusivity** *D***:** mobility of the diffusing particles

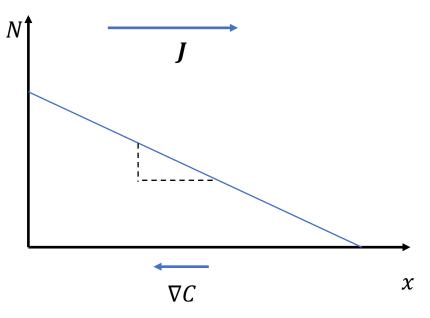
$$\frac{N}{A\Delta t} \approx D \frac{\Delta (N V^{-1})}{\Delta x}$$

$$\frac{moles}{m^2s} = [D] \frac{moles \cdot m^{-3}}{m}$$

$$[D] = \frac{m^2}{s}$$

- $D_{C0} = 0.2 \frac{cm^2}{s} \text{ in } air$
- $D_{C0} = 2 \times 10^{-5} \frac{cm^2}{s}$ in water

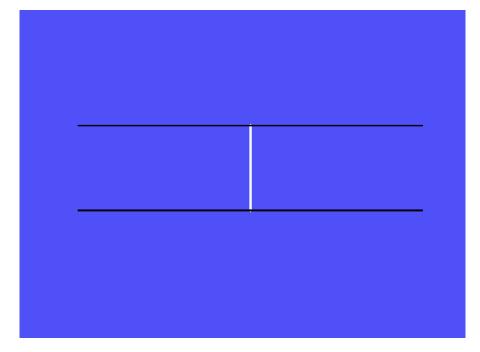


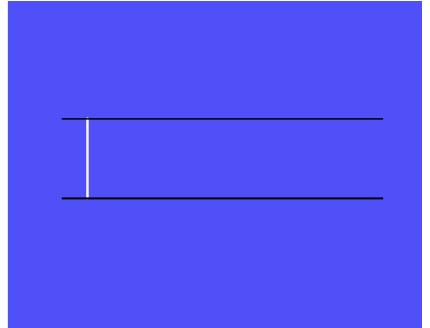


Flow and dispersion

Molecular dispersion (Diffusion)

Dispersion by flow





Molecular Diffusivity: gas kinetics

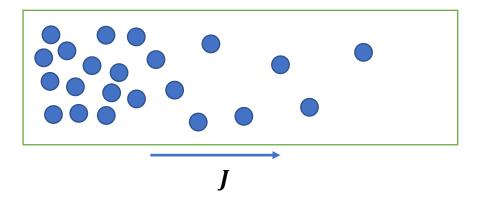
• Flux across a surface in a average time interval between collisions au

$$J \approx \frac{N}{A\tau} = \frac{N}{A\lambda} \overline{v} \sim c\overline{v}$$

$$J = -D\frac{dc}{dx} \approx D\frac{c}{\lambda}$$

$$D \approx \lambda \overline{v}$$

$$D \approx \frac{1}{\pi d^2} \frac{kT}{P} \times \frac{\sqrt{kT}}{\sqrt{m}} \sim \frac{T^{\frac{3}{2}}}{P}$$



Heat conduction

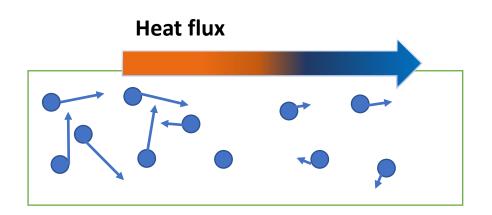
Fourier's law:

Rate of heat conduction is proportional to the temperature difference

$$J_Q = \frac{Q}{A\Delta t} = -k_t \nabla T$$

• Isocoric heat conduction $Q = \Delta U = C_V \Delta T$

$$rac{C_V \Delta T}{A au} pprox k_t rac{\Delta T}{\lambda}
ightarrow k_t pprox rac{C_V}{V} \lambda \overline{v}, \qquad C_V = rac{f}{2} rac{PV}{T}$$
 $k_t pprox rac{P}{T} D
ightarrow k_t \sim \sqrt{T}$



Fick's second law:

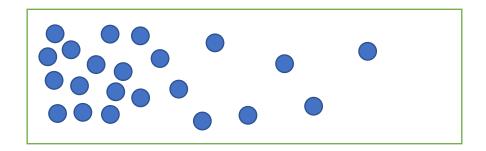
• Rate of change in the number of particles = net flux times the cross-section area

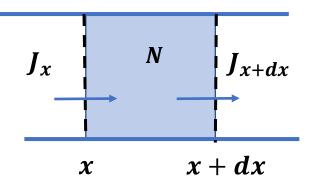
$$\frac{dN}{dt} = J_{net} \times A$$

$$\frac{dN}{dt} = (J_x - J_{x+dx}) \times A$$

$$\frac{d(N/A)}{dt} = -\frac{dJ_x}{dx}dx$$

$$\frac{\partial c}{\partial t} = D \frac{\partial^2 c}{\partial x^2}$$









Summary

Equipartition of energy

$$U = \frac{f}{2}kT$$
, f quadratic degrees of freedom

 Kinetic properties of gas depend on the mean-free path and mean velocity of the gas particles

$$D pprox \lambda \overline{v} \sim \frac{T^{rac{3}{2}}}{P}, \qquad \qquad k_t pprox \frac{C_V}{V} \lambda \overline{v} \sim \frac{\sqrt{T}}{P}$$