

UiO **Fysisk institutt**Det matematisk-naturvitenskapelige fakultet

Lecture 14



This week

- Wednesday: Tensors and vector fields. (Sections 6.5-6.6)
- Thursday: Problem set 7, some repetition of Lagrangians, tensors and relativistic problems including a (modified) exam question from 2006. (Please remember you need 6 problems sets before the exam!)
- Friday: Relativistic kinematics with four-vectors. (Section 7.1)

Recap

• We define the metric tensor g and inverse g-1 as

$$g = g^{-1} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix}$$

This is used to raise and lower indices

$$A^{\mu} = g^{\mu\nu} A_{\nu}, \ A_{\mu} = g_{\mu\nu} A^{\nu}$$

Lorentz transformations are given by

$$A^{\prime\mu}=L^{\mu}_{\nu}A^{\nu},~~A^{\prime}_{\mu}=L^{\nu}_{}A_{\nu}$$
 where L fulfils $g_{\mu\nu}L^{\mu}_{\rho}L^{\nu}_{\sigma}=g_{\rho\sigma}$

Today

- Tensors
 - Four-vectors with (possibly) more indices.
 - Rank (number of indices).
 - Tensor product.
- Vector and tensor fields
- Differentiation for four-vectors
 - The d'Alembertian.

Summary

- We can generalize four-vectors to rank-n tensors with n indices, e.g. the rank-2 F^{μν}.
- We also have tensor fields, e.g. $\phi(x)$, $A^{\mu}(x)$, $F^{\mu\nu}(x)$, that transform as relativistic tensors.
- We define a four-component derivative as

$$\partial_{\mu} \equiv \frac{\partial}{\partial x^{\mu}}; \quad \frac{\partial}{\partial x^{\mu}} = \left(\frac{1}{c}\frac{\partial}{\partial t}, -\frac{\partial}{\partial x}, -\frac{\partial}{\partial y}, -\frac{\partial}{\partial z}\right)$$
 that also transforms as a covariant four-vector.

This can be used to define the d'Alembertian

$$\partial_{\mu}\partial^{\mu} = \frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \vec{\nabla}^2$$