

UiO **Fysisk institutt**Det matematisk-naturvitenskapelige fakultet

Lecture 7



This week

- Wednesday: Hamilton's equations, Hamiltonian for charged particles in an electromagnetic field. (Sections 3.1 and 3.2)
- Thursday: Problem set 3 (main topic: even more e.o.m. from Lagrange's equation, cyclic coordinates and how to remove them)
- Friday: Application of Hamilton's equations, phase space. (Sections 3.2.1, 3.3, and 3.4)

Today

- Hamilton's equations
 - Derivation from Lagrange's equations
 - Simple example with 1D harmonic oscillator
 - Once more (with feeling) the charged particle in magnetic field example
 - Introducing (?) the Levi-Civita symbol (The Horror! The Horror!)

Recap

The Hamiltonian H is given by

$$H \equiv \sum_{i} \frac{\partial L}{\partial \dot{q}_{i}} \dot{q}_{i} - L = \sum_{i} p_{i} \dot{q}_{i} - L$$

 $p_i = \frac{\partial L}{\partial \dot{q}_i}$

 By solving the velocity in terms of p_i we can write H as a function of q_i and p_i (canonical position and momenta).

Recap

The e.m. potentials give a potential energy

$$U = e \phi - e \vec{v} \cdot \vec{A}$$

The conjugate momentum is

$$\vec{p} = m\vec{v} + e\vec{A}$$

and the Hamiltonian

$$H = \frac{1}{2m} (\vec{p} - e\vec{A})^2 + e\phi$$

The resulting e.o.m. are invariant under the gauge transformations

$$\phi \rightarrow \phi' = \phi - \frac{\partial \chi}{\partial t}, \quad \vec{A} \rightarrow \vec{A}' = \vec{A} + \vec{\nabla} \chi$$

Summary

 Expressing the Hamiltonian H in terms of the canonical position and momentum q_i and p_i the e.o.m can be written as Hamilton's equations

$$\dot{q}_i = \frac{\partial H}{\partial p_i}, \quad \dot{p}_i = -\frac{\partial H}{\partial q_i}, \quad i=1,\ldots,d$$

 To simplify calculations it is useful to introduce the totally antisymmetric Levi-Civita symbol

$$\epsilon_{ijk} \equiv \begin{cases} 1 & \text{for } i \neq j \neq k \text{ and cyclic} \\ -1 & \text{for } i \neq j \neq k \text{ and not cyclic} \\ 0 & \text{otherwise} \end{cases}$$