

MEK1100, utvalde formlar

Taylor polynom av andre orden

$$g(x, y) = g(x_0, y_0) + \frac{\partial g(x_0, y_0)}{\partial x}(x - x_0) + \frac{\partial g(x_0, y_0)}{\partial y}(y - y_0) \\ + \frac{1}{2} \frac{\partial^2 g(x_0, y_0)}{\partial x^2}(x - x_0)^2 + \frac{\partial^2 g(x_0, y_0)}{\partial x \partial y}(x - x_0)(y - y_0) + \frac{1}{2} \frac{\partial^2 g(x_0, y_0)}{\partial y^2}(y - y_0)^2$$

Derivasjonsformlar

$$\nabla \times \mathbf{A} = \left(\frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z} \right) \mathbf{i} + \left(\frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x} \right) \mathbf{j} + \left(\frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} \right) \mathbf{k}$$

$$\nabla(\kappa\beta) = \beta \nabla\kappa + \kappa \nabla\beta, \quad \nabla \cdot (\beta\mathbf{A}) = \nabla\beta \cdot \mathbf{A} + \beta \nabla \cdot \mathbf{A}, \quad \nabla \times (\beta\mathbf{A}) = \nabla\beta \times \mathbf{A} + \beta \nabla \times \mathbf{A}$$

$$\frac{D\beta}{dt} = \frac{\partial\beta}{\partial t} + \mathbf{v} \cdot \nabla\beta, \quad \frac{D\mathbf{v}}{dt} = \frac{\partial\mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla\mathbf{v}$$

$$\mathbf{v} \cdot \nabla\mathbf{v} = (\mathbf{v} \cdot \nabla v_x) \mathbf{i} + (\mathbf{v} \cdot \nabla v_y) \mathbf{j} + (\mathbf{v} \cdot \nabla v_z) \mathbf{k} = v_x \frac{\partial\mathbf{v}}{\partial x} + v_y \frac{\partial\mathbf{v}}{\partial y} + v_z \frac{\partial\mathbf{v}}{\partial z}$$

Sylinderkoordinatar

Transformasjon: $x = r \cos \theta$, $y = r \sin \theta$, $\mathbf{i}_r = \cos \theta \mathbf{i} + \sin \theta \mathbf{j}$, $\mathbf{i}_\theta = -\sin \theta \mathbf{i} + \cos \theta \mathbf{j}$.

Derivasjon: ($\mathbf{A} = A_r \mathbf{i}_r + A_\theta \mathbf{i}_\theta + A_z \mathbf{k}$)

$$\nabla\beta = \frac{\partial\beta}{\partial r} \mathbf{i}_r + \frac{1}{r} \frac{\partial\beta}{\partial\theta} \mathbf{i}_\theta + \frac{\partial\beta}{\partial z} \mathbf{k}, \quad \nabla^2\beta = \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial\beta}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2\beta}{\partial\theta^2} + \frac{\partial^2\beta}{\partial z^2}$$

$$\nabla \cdot \mathbf{A} = \frac{1}{r} \left(\frac{\partial}{\partial r} (r A_r) + \frac{\partial A_\theta}{\partial\theta} \right) + \frac{\partial A_z}{\partial z}$$

$$\nabla \times \mathbf{A} = \left(\frac{1}{r} \frac{\partial A_z}{\partial\theta} - \frac{\partial A_\theta}{\partial z} \right) \mathbf{i}_r + \left(\frac{\partial A_r}{\partial z} - \frac{\partial A_z}{\partial r} \right) \mathbf{i}_\theta + \frac{1}{r} \left(\frac{\partial}{\partial r} (r A_\theta) - \frac{\partial A_r}{\partial\theta} \right) \mathbf{k}$$

Kulekoordinatar

Transformasjon: $x = r \sin \theta \cos \phi$, $y = r \sin \theta \sin \phi$, $z = r \cos \theta$,

$\mathbf{i}_r = \sin \theta \cos \phi \mathbf{i} + \sin \theta \sin \phi \mathbf{j} + \cos \theta \mathbf{k}$,

$\mathbf{i}_\theta = \cos \theta \cos \phi \mathbf{i} + \cos \theta \sin \phi \mathbf{j} - \sin \theta \mathbf{k}$,

$\mathbf{i}_\phi = -\sin \phi \mathbf{i} + \cos \phi \mathbf{j}$.

Derivasjon: ($\mathbf{A} = A_r \mathbf{i}_r + A_\theta \mathbf{i}_\theta + A_\phi \mathbf{i}_\phi$)

$$\nabla\beta = \frac{\partial\beta}{\partial r} \mathbf{i}_r + \frac{1}{r} \frac{\partial\beta}{\partial\theta} \mathbf{i}_\theta + \frac{1}{r \sin \theta} \frac{\partial\beta}{\partial\phi} \mathbf{i}_\phi$$

$$\nabla \cdot \mathbf{A} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 A_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial\theta} (\sin \theta A_\theta) + \frac{1}{r \sin \theta} \frac{\partial A_\phi}{\partial\phi}$$

Integral

Parametrisering av flate $\Omega \rightarrow \sigma: \mathbf{r} = \mathbf{r}(t, s)$

$$\int_{\sigma} \mathbf{v} \cdot \mathbf{n} d\sigma = \iint_{\Omega} \mathbf{v} \cdot \left(\pm \frac{\partial \mathbf{r}}{\partial t} \times \frac{\partial \mathbf{r}}{\partial s} \right) dt ds. \quad \int_{\sigma} p \mathbf{n} d\sigma = \iint_{\Omega} p \left(\pm \frac{\partial \mathbf{r}}{\partial t} \times \frac{\partial \mathbf{r}}{\partial s} \right) dt ds$$

Sylinder: $d\sigma = r d\theta dz$; Kule: $d\sigma = r^2 \sin \theta d\theta d\phi$

Integrert fluks gjennom kurve

$$\int_{\lambda} \mathbf{v} \cdot \mathbf{n} ds = \int_{\lambda} v_x dy - v_y dx$$

Mindre vanlege former av Gauss sats

$$\int_{\sigma} \mathbf{n} \times \mathbf{A} d\sigma = \int_{\tau} \nabla \times \mathbf{A} d\tau, \quad \int_{\sigma} \mathbf{n} \beta d\sigma = \int_{\tau} \nabla \beta d\tau$$

Divergens- og kvervelfrie felt i 2D

Straumfunksjon: $v_x = -\frac{\partial \psi}{\partial y}$, $v_y = \frac{\partial \psi}{\partial x}$ eller $\mathbf{v} = \nabla \times (-\mathbf{k}\psi)$.

Eksempel:	eksempel på ein stagnasjonsstrøm $\phi = \frac{1}{2}A(x^2 - y^2)$	kjelde/sluk $\phi = A \ln r$	punktkvervel $\phi = A\theta$	dipol $\phi = \frac{Ax}{x^2 + y^2}$
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Hydrodynamiske likningar

Kontinuitetslikninga

$$\frac{D\rho}{dt} = -\rho \nabla \cdot \mathbf{v}$$

Euler's bevegelseslikning i tyngdefeltet (\mathbf{g} er tyngdeakselerasjonen)

$$\frac{D\mathbf{v}}{dt} = -\frac{1}{\rho} \nabla p + \mathbf{g}$$

Bernoulli's likning i tyngdefeltet

$$\frac{1}{2} \mathbf{v}^2 + \frac{p}{\rho} + gz = \text{konstant}$$

Varme

Spesifikk varmeenergitettleik for inkompressible medium: $E(T)$.

Spesifikk varmekapasitet: $c = \frac{\partial E}{\partial T}$.

Varmeflukstettleik: $\mathbf{H} = \mathbf{H}_s + \mathbf{H}_l$, $\mathbf{H}_s = \rho E \mathbf{v}$ (adveksjon), $\mathbf{H}_l = -k \nabla T$ (varmeleiing).

Varmelikninga:

$$\frac{DT}{dt} = \kappa \nabla^2 T + \frac{q}{\rho c}, \quad \text{kor} \quad \kappa = \frac{k}{\rho c}$$